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**Faculty of Sciences**

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**TITLE**

**MECHANICS**  
**OF THE MATERIAL POINT**  
**(COURSES AND CORRECTED EXERCISES)**

**Addressed to students level : License (L1)**

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# PREFACE

This physics text book is intended for a broad range of students; it covers the fields of Sciences and Technology (ST), Material sciences (SM), and in particular is suitable for first-year LMD (License, Master, and Doctorate) students in Mathematics and Computer Sciences.

This document follows the official curriculum of the Physics 1 module taught in the first year of the LMD programs mentioned earlier. This manuscript is essential for first-year students given the importance of concentration and a thorough understanding of scientific concepts. The combination of these elements, along with the requirement to take notes during the course, can be crucial for effective learning in disciplines such as physics and the sciences. Hence, the necessity to provide such documents to our students.

The proposed work is the result of my experience in courses and supervised work for several years at the University of Tlemcen. This manual compiles all the courses on the mechanics of the material point, with varying levels of detail, along with exercises and problems with solutions and a set of exams from previous years.

This handout comprises six chapters on mechanics. The first one focuses on dimensional analysis. Mathematical reminders and concepts about vectors are presented in the second chapter. The third chapter deals with the kinematics of material point in various coordinate systems. Chapter 4 details relative motion or changes in reference frames. The dynamics of material points is covered in the fifth chapter. Finally, the sixth chapter is dedicated to work and energy. I must stress that this document in no way replaces face-to-face tutorials.

I hope that this collection of exercises and solved test problems in point mechanics will be of great help to the majority of students.

Any comment, proposal or constructive criticism allowing the improvement of the texts will be collected with great interest.

Dr. Hadjou Bélaïd Zakia

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## Glossary

<b>In English</b>	<b>In French</b>	<b>In Arabic</b>
Mechanics	La mécanique	ميكانيك
Physical quantities	Une grandeur physique	مقدار فيزيائي
Dimensional analysis	Analyse dimensionnelle	التحليل البعدي
Uncertainty Calculation	Calcul d'incertitudes	حساب الارتيابات
Unit	Unité	الوحدة
Force	La force	القوة
Motion	Le Mouvement	الحركة
Velocity (speed)	La vitesse	السرعة
Acceleration	L'accélération	التسارع
Length	Lalongeur	المسافة
Time	Le temps	الزمن
Mass	La masse	الكتلة
Weight	Le poids	الوزن
Momentum	La quantité du mouvement	كمية الحركة
Work	Le travail	العمل
Energy	L'énergie	الطاقة
Power	La puissance	الاستطاعة
Equilibrium	Etat d'équilibre	التوازن
Surface	La surface	المساحة
Volume	Le volume	الحجم
Density	La masse volumique	الكتلة الحجمية
Frequency	La fréquence	التواتر
Linear velocity	La vitesse linéaire	السرعة الخطية
Angular velocity	La vitesse angulaire	السرعة الزاوية
Linear Acceleration	L'accélération linéaire	التسارع الخطي
Angular Acceleration	L'accélération angulaire	التسارع الزاوي
Pressure	La prèssion	الضغط
Acceleration of gravity	Accélération de pesanteur	تسارع الجاذبية



## Chapter I: Dimensional Analysis and Uncertainty Calculation

Current intensity	L'intensité du courant	شدة التيار الكهربائي
Light intensity	L'intensité de la lumière	شدة الضوء
Quantity of material	La quantité de la matière	كمية المادة
Height	La hauteur	الارتفاع
Dimensionless	Sans dimension	بدون بعد
The period of a pendulum	La période d'une pendule simple	دور نواس بسيط
The sound	Le son	الصوت
Radius	Le rayon	نصف قطر
Relative uncertainty	L'incertitude relative	الارتياح النسبي
Absolute uncertainty	L'incertitude absolue	الارتياح المطلق
Total Differential method	Méthode différentielle totale	طريقة التفاضلية الكلية
Logarithmic Method	Méthode logarithmique	الطريقة اللوغارتمية
Absolute Error	Erreur absolue	الخطا المطلق
Relative Error	Erreur relative	الخطا النسبي
The International System (SI)	Système international SI	النظام الدولي
The CGS system	Système CGS	النظام CGS
Average speed	La vitesse moyenne	السرعة المتوسطة
Instantaneous speed	La vitesse instantanée	السرعة اللحظية

## Part 1: Dimensional analysis

### التحليل البعدي

#### 1. Introduction

The observation of physical phenomena is incomplete if it does not lead to quantitative information, which is the measurement of physical quantities. To study a physical phenomenon, one must examine the important variables; the mathematical relationship between these variables constitutes a physical law.

This is possible in certain cases, but for other cases, it is necessary to use a modeling method such as dimensional analysis. (التحليل البعدي)

#### 2. Definition of Dimensional Analysis تعريف التحليل البعدي

It is a theoretical tool for interpreting problems based on the dimensions of the involved physical quantities: length, time, mass, and so on.

Dimensional analysis allows for:

- Verifying the validity of dimensioned equations.
- Investigating the nature of physical quantities.
- Exploring the homogeneity of physical laws.
- Determining the unit of a physical quantity based on fundamental units (meter, second, kilogram, etc.).

#### 3. Physical Quantity مقدار فيزيائي

A physical quantity is an observable and measurable property through a specifically designed instrument. Mechanics acknowledges seven fundamental physical quantities: length, time, mass, electric current, temperature, quantity of material, and luminous intensity. Other physical quantities, known as derived quantities, are expressed in terms of these three fundamental quantities, such as velocity, acceleration, force, and more.....

# Chapter I: Dimensional Analysis and Uncertainty Calculation

## Note :

In general, for first-year students in Mathematics (M), and Computer Science (I), the focus is primarily on the first three fundamental quantities: length, time, and mass.

## **4. International System of Units** الوحدة في النظام العالمي

The value of a physical quantity is given in relation to a standard known as a "unit". The first four fundamental units constitute the MKSA International System (Meter, Kilogram, Second, Ampere). Using these fundamental units, derived units can be constructed: area ( $m^2$ ), velocity ( $m/s$ ), force ( $kg\ m/s^2$ )...

<b>Fundamental quantities</b> المقادير الأساسية	<b>Units الوحدة (in the international system MKSA)</b>	<b>Symbols الرمز</b>
<b>Length</b>	<b>Meter</b>	<b>(m)</b>
<b>Mass</b>	<b>Kilogram (kg)</b>	<b>(kg)</b>
<b>Time</b>	<b>Second</b>	<b>(s)</b>
<b>Current intensity</b>	<b>Ampere</b>	<b>(A)</b>
<b>Temperature</b>	<b>Kelvin</b>	<b>(K)</b>
<b>Light intensity</b>	<b>Candela</b>	<b>(Cd)</b>
<b>Quantity of material</b>	<b>Mole</b>	<b>(mol)</b>

There are specific units such as N (Newton) for force, Hz (Hertz) for frequency, Watt for power, Pascal (Pa) for pressure...

**Note:** There are two systems of units:

- The International System (SI) known as MKSA (Meter, Kilogram, Second, Ampere), which is the most widely used system.

- The CGS system (Centimeter, Gram, Second), which is less commonly used.

# Chapter I: Dimensional Analysis and Uncertainty Calculation

## 5. Dimensional Equations معادلة ابعاد

Dimension represents the nature of a physical quantity. A physical quantity has only one possible dimension.

The dimension of a quantity  $G$  is denoted by:  $[G] = L$

By denoting  $M$ ,  $L$ , and  $T$  as the dimensions of the fundamental quantities mass, length, and time, we can express the dimensions of other derived quantities in terms of these three fundamental dimensions. The resulting equations are the dimension equations for these physical quantities.

The Fundamental Quantities المقادير الأساسية	Symbols الرمز	Dimensions الأبعاد	Units الوحدة (International System (SI))
<i>Length</i>	<i>L</i>	$[l] = L$	<i>Meter (m)</i>
<i>Mass</i>	<i>M</i>	$[m] = M$	<i>Kilogram (kg)</i>
<i>Time</i>	<i>T</i>	$[t] = T$	<i>Second (s)</i>
Current intensity	I	$[I] = I$	Ampere (A)
Temperature	T	$[T] = \theta$	Kelvin (K)
Light intensity	J	$[j] = J$	Candela (Cd)
Quantity of material	N	$[n] = N$	Mole (mol)

### Example :

- $[\text{speed}] = [v] = \frac{[\text{length}]}{[\text{time}]} = \frac{[l]}{[t]} = \frac{L}{T} = LT^{-1}$  and the unit of speed is (m/s).
- $[\text{acceleration}] = [a] = \frac{[\text{speed}]}{[\text{time}]} = \frac{[v]}{[t]} = \frac{LT^{-1}}{T} = LT^{-2}$  and the unit of acceleration is (m/s<sup>2</sup>).
- $[\text{Force}] = [F] = [\text{mass}][\text{acceleration}] = [m][a] = MLT^{-2}$  and the unit of force is Newton or (kg.m/s<sup>2</sup>).

# Chapter I: Dimensional Analysis and Uncertainty Calculation

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## Notes :

- The dimension of constants is always equal to 1; we say they are dimensionless.
- Angles and functions like sin, cos, tan, exp, ln, and log are dimensionless functions.

$$[\text{Numeric value}] = 1, [\text{angle}] = 1, [\cos \alpha] = [\sin \alpha] = [\tan \alpha] = [\cot \alpha] = [\ln x] = [e^x] = 1$$

## **6. Homogeneity of Dimensional Equations** *تجانس معادلة ابعاد*

The two sides of a dimension equation must have the same dimensions since they represent quantities of the same nature.

G is a physical quantity:

- $G = A \pm B \Rightarrow [G] = [A] = [B]$
- $G = A * B \Rightarrow [G] = [A] * [B]$
- $G = A/B \Rightarrow [G] = [A]/[B]$
- $G = A^n \Rightarrow [G] = [A]^n$

## Note:

- A heterogeneous (non-homogeneous *غير متجانسة*) equation is necessarily False.
- A homogeneous equation is not necessarily true.
- Dimensions cannot be added (or subtracted).

## Example 1:

$y = \frac{1}{2}at^2 + v_0t + y_0$  is the equation of a physical law.

- Check that this equation is homogeneous?

This equation is homogeneous if:  $[y] = \left[ \frac{1}{2} at^2 \right] = [v_0t] = [y_0]$

We have :

$$[y] = [y_0] = L,$$

$$\left[ \frac{1}{2} at^2 \right] = \left[ \frac{1}{2} \right] [a][t]^2 = 1 LT^{-2}T^2 = L,$$

## Chapter I: Dimensional Analysis and Uncertainty Calculation

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And  $[v_0t] = [v_0][t] = LT^{-1}T = L$

So :

$[y] = \left[\frac{1}{2}at^2\right] = [v_0t] = [y_0]$  is checked.

Hence the equation  $y = \frac{1}{2}at^2 + v_0t + y_0$  is homogeneous.

### Notes:

We can use this property of dimension equations to discover physical laws by knowing the variables involved in the given physical phenomenon and the relationship among them.

### Example 2:

The period is given in terms of length and severity by the following relationship:

$$T = k \cdot l^x \cdot g^y$$

- Give the physical law of period T *الطور*?

For this it is necessary to determine the exponents x and y.

It is assumed that the equation is homogeneous so:  $[T] = [k][l]^x[g]^y$

The dimensions of all physical quantities in the study relationship are written.

$$[l] = L, [k] = 1 \text{ and } T \text{ is a time so } [T] = T$$

We have weight force *قوة النقل*  $p = mg$  with:

$$[p] = [m][g] \Rightarrow [g] = \frac{[p]}{[m]}$$

P : the weight, is a force so it has the dimension of a force:  $[p] = [F] = MLT^{-2}$ .

$$\text{So } [g] = \frac{[F]}{[m]} = \frac{MLT^{-2}}{M} = LT^{-2}$$

g is the acceleration  $[g] = LT^{-2}$

$$\text{Hence : } T = 1 \cdot L^x \cdot (LT^{-2})^y \Rightarrow M^0 L^0 T^1 = L^{x+y} \cdot T^{-2y}$$

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By identification we will have:

$$\begin{cases} x + y = 0 \\ -2y = 1 \end{cases} \Rightarrow \begin{cases} y = -\frac{1}{2} \\ x = -y = \frac{1}{2} \end{cases} \quad \text{so} \quad T = kl^{\frac{1}{2}}g^{-\frac{1}{2}}$$

$\Rightarrow T = k\sqrt{\frac{l}{g}}$  it's the law of the period.

### Example 3:

The average speed of the particles is expressed as a function of the mass **m**, the volume **V**, and the pressure **p** by the following expression:

$$v=f(m, V, p) \Rightarrow v = k \cdot m^\alpha \cdot V^\beta \cdot p^\gamma$$

It is assumed that the equation is homogeneous, therefore:  $[v] = k[m]^\alpha[V]^\beta[p]^\gamma \dots\dots\dots(1)$

$$\text{with } [m] = M, [v] = LT^{-1}, [V] = L^3, [p] = \frac{[F]}{[s]} = \frac{[m][a]}{[s]} = \frac{MLT^{-2}}{L^2} = ML^{-1}T^{-2}$$

$$(1) \Rightarrow M^0LT^{-1} = M^\alpha L^{3\beta} (ML^{-1}T^{-2})^\gamma \Rightarrow M^0LT^{-1} = M^{\alpha+\gamma} L^{3\beta-\gamma} T^{-2\gamma}$$

By identification we will have:

$$\begin{cases} \alpha + \gamma = 0 \\ 3\beta - \gamma = 1 \\ -2\gamma = -1 \end{cases} \Rightarrow \begin{cases} \gamma = \frac{1}{2} \\ \alpha = -\gamma = -\frac{1}{2} \\ \beta = \frac{1 + \gamma}{3} = \frac{1 + 1/2}{3} = \frac{1}{2} \end{cases}$$

$$\text{So: } v = km^{-\frac{1}{2}}V^{\frac{1}{2}}p^{\frac{1}{2}} \Rightarrow v = k\sqrt{\frac{pV}{m}}$$

It's a law of the average speed of the particles.

### Conclusions:

Dimensional analysis serves the following purposes:

- Verification of the homogeneity of physical formulas.
- Determination of the nature and the unit of a physical quantity.
- Exploration of the general form of physical laws.

## 2<sup>nd</sup> part: Uncertainty Calculation

### حساب الارتيابات

#### 1. Introduction

In an experiment, exact measurements do not exist. These are always accompanied by more or less significant errors depending on the measurement method used, the quality of the instruments used, and the role of the operator. The measuring instrument, even if built upon a standard, also has a certain precision as provided by the manufacturer. Therefore, measurements are carried out with approximations. Estimating the errors made in measurements and their consequences is essential.

#### 2. Absolute and relative uncertainty الارتياب المطلق و الارتياب النسبي

##### 2.1. Absolute error الخطا المطلق

The absolute error of a measured quantity  $G$  is the difference  $\delta G$  between the experimental value  $G_m$  and a reference value that can be considered as exact,  $G_e$ . In reality, since the exact value is inaccessible, it is approximated by taking the average of a series of measurements of the quantity  $G$ .

$$\delta G = |G_{measured} - G_{exact}|$$

##### 2.2. Relative error الخطا النسبي

The relative error is the ratio of the absolute error to the reference value. The relative error is dimensionless; it indicates the quality (precision) of the obtained result. It is expressed in terms of a percentage.

$$\frac{\delta G}{G} = \frac{|G_{measured} - G_{exact}|}{G_{measured}}$$

##### 2.3. Absolute uncertainty الارتياب المطلق

This is the maximum error that can be committed in the evaluation.

$$\Delta G \geq |\delta G| \Rightarrow \Delta G = |G_{ex} - G_m| \Rightarrow G_{ex} = G_m \pm \Delta G$$



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$$\Rightarrow G_m - \Delta G \leq G_{ex} \leq G_m + \Delta G$$

The general form is:

$$G_{ex} = G_m \pm \Delta G$$

The absolute uncertainty has the same unit as the measured quantity and is always positive.

**Example:**  $m=12,121\text{g}$  and  $\Delta m=0,02\text{g}$

The correct expression for the measurement (the condensed writing) of  $m$  is:

$$m = (12,121 \pm 0,020) \text{ g}$$

## 2.4. Relative uncertainty الارتياب النسبي

The relative uncertainty is the ratio between the absolute uncertainty and the measured value of  $G$ . It is also expressed in terms of a percentage and is a convenient way to quantify the precision of a measurement. It is denoted as:  $\Delta G/G$

It is given in percentage and it is always smaller than 1.

$$\frac{\Delta m}{m} = \frac{0,02}{12,121} \approx 0.16\%$$

## 3. Uncertainty Calculation حساب الارتيابات

Generally, there are two mathematical methods for uncertainty calculation: the total differential method, which is a general approach, and the logarithmic method, which is limited to physical laws expressed as a product or a ratio.

### 3.1. The total differential method الطريقة التفاضلية الكلية

Let  $f(x, y, z)$  be a function that depends on three variables  $x, y, z$ :

The total differential of the function ( $f$ ) is expressed by the following equation:

$$df = \left(\frac{\partial f}{\partial x}\right)_{y,z=cst} dx + \left(\frac{\partial f}{\partial y}\right)_{x,z=cst} dy + \left(\frac{\partial f}{\partial z}\right)_{x,y=cst} dz$$

$\left(\frac{\partial f}{\partial x}\right)$  is the partial differential of the function  $f$  with respect to  $x$ , considering  $y$  and  $z$  as constants.

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$\left(\frac{\partial f}{\partial y}\right)$  is the partial differential of the function  $f$  with respect to  $y$ , considering  $x$  and  $z$  as constants.

$\left(\frac{\partial f}{\partial x}\right)$  is the partial differential of the function  $f$  with respect to  $x$ , considering  $y$  and  $z$  as constants.

The absolute uncertainty on  $(f)$  is generally expressed in the form:

$$\Delta f = \left|\frac{\partial f}{\partial x}\right| \Delta x + \left|\frac{\partial f}{\partial y}\right| \Delta y + \left|\frac{\partial f}{\partial z}\right| \Delta z$$

### **Example:**

Let  $f(x, y)$  be a physical quantity that depends on two variables  $x$  and  $y$ .

$f$  is expressed as  $f(x, y) = 2xy + x^2y$

The total differential of “ $f$ ” will be given by :  $df = \left(\frac{\partial f}{\partial x}\right) dx + \left(\frac{\partial f}{\partial y}\right) dy$

With:  $\left(\frac{\partial f}{\partial x}\right)_{y=cst} = 2y + 2xy$  and  $\left(\frac{\partial f}{\partial y}\right)_{x=cst} = 2x + x^2$

So:  $df = (2y + 2xy)dx + (2x + x^2)dy$

Hence the absolute uncertainty on the quantity «  $f$  » is given by:

$$\Delta f = |2y + 2xy|\Delta x + |2x + x^2|\Delta y$$

### **3.2. Logarithmic method الطريقة اللوغاريتمية**

This method is based on the logarithm and its derivative.

Consider a three-variable function,  $G = f(x, y, z)$ . To calculate the relative uncertainty on the function  $G$  using the logarithmic differential method, the following steps should be followed:

1. Introduce the logarithmic function to the function  $G$ .
2. Calculate  $d(\log G) = dG / (G \ln 10)$  or  $d(\ln G) = dG / G$ .
3.  $\frac{dG}{G} \leq \frac{\Delta G}{G}$ , and deduce the relative uncertainty on  $G$ .

## Chapter I: Dimensional Analysis and Uncertainty Calculation

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### Example

Let the function  $f(x,y,z)=x^a y^b z^c$

x, y et z are a variables and a, b et c are constants in the exponents.

First, we find the logarithm of “f”:

$$\log f = \log x^a y^b z^c = \log x^a + \log y^b + \log z^c = a \log x + b \log y + c \log z$$

Then, we calculate the derivative of log f:

$$d \log f = a d \log x + b d \log y + c d \log z$$

$$\frac{df}{f} = a \frac{dx}{x} + b \frac{dy}{y} + c \frac{dz}{z}$$

$$\frac{\Delta f}{f} = \left| \frac{a}{x} \right| \Delta x + \left| \frac{b}{y} \right| \Delta y + \left| \frac{c}{z} \right| \Delta z$$

If x, y, z, a, b and c are positive, we can write  $\Delta f/f$  by :

$$\frac{\Delta f}{f} = a \frac{\Delta x}{x} + b \frac{\Delta y}{y} + c \frac{\Delta z}{z}$$

$\frac{\Delta f}{f}$  Represents relative uncertainty on the quantity « f » الارتياب النسبي.

## Proposed exercises about chapter I

### Exercise 1

Find the dimension of the following physical quantities:

Surface, Volume, Density, Frequency, Linear Velocity, Angular Velocity, Linear Acceleration, Angular Acceleration, Force, Work, Energy, Power, and Pressure.

### Exercise 2

The characteristic equation of a constant temperature fluid is as follows:

$$\left(p + \frac{a}{V^2}\right)(V - b) = c$$

Or  $p$  is the pressure and  $V$  is the volume.

Determine the dimensions of quantities  $a$ ,  $b$  and  $c$ .

### Exercise 3

Check the homogeneity of this formula:

$$p = \rho g h_1 + h_2 F$$

Such as:  $P$  pression,  $\rho$  density,  $g$  an acceleration of gravity,  $h_1$  and  $h_2$  are heights and  $F$  a force.

### Exercise 4

The trajectory  $y=f(x)$  of a projectile with an initial velocity( $v_0$ ) from a point (O) located at heigh ( $h$ ) above the impact plane is given by the following formula

$$y = \frac{g}{2v_0^2} x^2 + h$$

Show that this formula is homogeneous.

### Exercise 5

Are the following formulas dimensionally valid?

1.  $F = \frac{Gm}{r}$ , such as:  $F$  is a force,  $G$  is a constant expressed in  $\frac{m^3}{kg s^2}$ ,  $m$  is a mass, and  $r$  is a length.
2.  $p = \rho g h_1 + h_2 F$  such as:  $P$  is a pressure,  $g$  is the acceleration due to gravity,  $h_1$  and  $h_2$  are heights, and  $F$  is a force.

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3.  $\theta = \frac{b \sin(a)}{t \cos(c)}$ , such as: b and t have dimensions of length.

### Exercise 6

1. In a fluid, a ray ball (نصف القطر)  $r$  animated by a velocity  $v$ , is subjected to a friction force given by  $F = -6 \cdot \pi \cdot \eta \cdot r \cdot v$ , where  $\eta$  is the viscosity of the fluid.

What is the dimension of  $\eta$ ?

2. When the ball is dropped without initial speed at the moment  $t = 0$ , its speed is written to

$$t > 0: \quad v = a \left( 1 - \exp\left(-\frac{t}{b}\right) \right)$$

Where  $a$  and  $b$  are two quantities that depend on the characteristics of the fluid. What are the dimensions of  $a$  and  $b$ ?

### Exercise 7

Experiments have shown that the speed  $v$  of sound in a gas is only dependent on the volumetric mass density  $\rho$  and the coefficient of compressibility  $\chi$ . What is the law that provides the speed  $v$  as a function of the gas's characteristics? It is noted that  $\chi$  has unit equivalent to the inverse of pressure.

### Exercise 8

The sound emitted by the wire of a guitar is characterized by its frequency  $f$ . This frequency is a function of the force  $F$  of the wire tension, the length  $L$  and the density  $\rho$  of the wire.

Find the expression of frequency  $f$  assuming the form:

$$f = K F^a L^b \rho^c$$

With  $K$  a dimensionless constant and the frequency dimension  $[f] = T^{-1}$ .

### Exercise 9

Let the simple pendulum formed of a ball (sphere) of radius  $R$  and mass  $m$ . The study of the effect of the air on this pendulum shows that its period  $T$  depends on a constant  $k$ , the coefficient of the air  $\eta$ , the radius of the ball  $R$  and its density  $\rho$ .

- 1- Find the expression of the period assuming the form:

$$T = K \eta^x R^y \rho^z \text{ with } [\eta] = ML^{-1}T^{-1}$$

- 2- Determine relative uncertainty on  $T$  based on  $\Delta\eta$ ,  $\Delta R$  and  $\Delta m$ .

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### Exercise 10

The expression for a physical quantity  $G$  is:

$$G = \frac{T^2 g a}{4\pi^2} - a^2.$$

Where  $T$  represents time,  $a$  is a length, and  $g$  is the acceleration due to gravity.

1. Determine the dimension of  $G$  and deduce its unit.
2. Calculate the absolute uncertainty  $\Delta G$  in terms of  $\Delta T$  and  $\Delta a$ .

### Exercise 11

The refractive index  $n$  of a substance is given by the relation:

$$n = \sqrt{N^2 - \sin^2 \alpha}$$

Where  $N$  is the prism index and  $\alpha$  is an angle.

1. Calculate the absolute uncertainty  $\Delta n$  by considering that  $\Delta n = f(N, \alpha)$ .
2. Deduce the relative uncertainty  $\Delta n/n$ .

### Exercise 12

The speed limit reached by a weighted parachute is a function of its weight  $P$  and its surface

$S$ , is given by:  $v = \sqrt{\frac{P}{k.S}}$

- 1) Give the dimension of the constant  $k$ .
- 2) Calculate the speed limit of a parachute having the following characteristics:
- 3) The weight being known to the nearest 2 % and the surface to 3 %, calculate the relative uncertainty  $\Delta v/v$  on the velocity  $v$ , thus the absolute uncertainty  $\Delta v$  and deduce the condensed writing of this velocity.

We give:  $M=90$  kg,  $S=80$  m<sup>2</sup>,  $g=9,81$  m/s<sup>2</sup>, and  $k=1,15$  MKS.

### Exercise 13

The height  $H$  of a liquid of mass  $M$  contained in a cylinder of radius  $R$  is given by the relation:

$$H = \frac{(2. \sigma. \cos \alpha)}{(R. g. \rho)}$$

Where  $\alpha$  is the liquid-cylinder contact angle,  $\rho$  the density of the liquid and  $g$  the gravity acceleration.

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- 1- Using the dimensional equations, find the dimension of  $\sigma$ .
- 2- Determine relative uncertainty on  $\sigma$  based on absolute uncertainties  $\Delta R$ ,  $\Delta g$ ,  $\Delta M$  and  $\Delta \alpha$ .

### Exercise 14

The resonance frequency  $f$  of an electric circuit is given by the formula:

$$f = \frac{1}{2\pi\sqrt{L \cdot C}}$$

$L$  and  $C$  are known with absolute uncertainties  $\Delta L$  and  $\Delta C$ .

Determine as a function of  $L$ ,  $C$ ,  $\Delta L$  and  $\Delta C$  absolute and relative uncertainties on  $f$  with the two differential methods.

### Exercise 15

A) The sound emitted by a guitar wire is characterized by its frequency  $f$ . This frequency is a function of the force  $F$  of the wire tension, the length  $L$  and the density  $\rho$  of the wire.

Or :  $f = K F^a L^b \rho^c$  such  $K$  is constant dimensionless.

Determine the relationship of  $f$ .

B) The focal length  $f$  of a lens is determined from the formula:

$$f = \frac{D^2 - a^2}{4D}$$

Calculate the absolute uncertainty  $\Delta f$  as a function of  $\Delta D$  and  $\Delta a$ .

## Correction of exercises about chapter I

### Exercise 1

- Surface :

We have:  $[l]=L$ ,  $[t]=T$  and  $[m]=M$ .

[Physical quantities] =  $M^x L^y T^z$

$$S = l \times l \Rightarrow [S] = L.L = L^2 \Rightarrow [S] = L^2 \quad (\text{m}^2)$$

- Volume :

$$V = l \times l \times l \Rightarrow [V] = L.L.L = L^3 \Rightarrow [V] = L^3 \quad (\text{m}^3)$$

- Density:

$$\rho = \frac{m}{V} \text{ so } [\rho] = \frac{[m]}{[V]} = \frac{M}{L^3} = ML^{-3} \Rightarrow [\rho] = ML^{-3} \quad (\text{kg/m}^3)$$

- Frequency:

$$f = \frac{1}{T} \Rightarrow [f] = \frac{1}{[T]} = \frac{1}{T} = T^{-1} \Rightarrow [f] = T^{-1} \quad (\text{s}^{-1} \text{ or Hertz})$$

(Note: Period  $[T] = T$  ; unit is « s »)

- Linear velocity:

$$v = \frac{dx}{dt} \Rightarrow [v] = \frac{[x]}{[t]} = \frac{L}{T} = LT^{-1} \Rightarrow [v] = LT^{-1} \quad (\text{m./s})$$

- Angulaire velocity :

$$\omega = \theta \cdot = \frac{d\theta}{dt} = \frac{v}{R} \Rightarrow [\omega] = \frac{[\theta]}{[t]} = \frac{1}{T} = T^{-1} \Rightarrow [\omega] = T^{-1} \quad (\text{Rd/s})$$

$[angle] = 1$  i and its unit is the radian (rad).

- Linear acceleration:

$$a = \frac{dv}{dt} \Rightarrow [a] = \frac{[dv]}{[dt]} = \frac{LT^{-1}}{T} = LT^{-2} \Rightarrow [a] = LT^{-2} \quad (\text{m./s}^2)$$

- Angulaire acceleration :

$$\omega \cdot = \theta \cdot \cdot = \frac{d\theta \cdot}{dt} \Rightarrow [\omega \cdot] = \frac{[d\theta \cdot]}{[dt]} = \frac{T^{-1}}{T} = T^{-2} \Rightarrow [\omega \cdot] = T^{-2} \quad (\text{Rd./s}^2)$$



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- Force :

$$F = m \times a \Rightarrow [F] = [m] \times [a] = M.L.T^{-2} \Rightarrow [F] = MLT^{-2} \text{ (kg.m.s}^{-2} \text{ or Newton)}$$

- Work :

$$W = F \times d \times \cos\alpha \Rightarrow [W] = [F] \times [d] \times [\cos\alpha] = MLT^{-2}.L.1 = ML^2T^{-2} \text{ (kg.m}^2\text{.s}^{-2} \text{ or Joule)}$$

- Energy :

$$E_C = (\frac{1}{2}).m.v^2 \Rightarrow [E] = [1/2].[m].[v]^2 = ML^2T^{-2} \text{ (Joule)}$$

$$E_p = m.g.h \Rightarrow [E] = [m].[g].[h] = M.LT^{-2}.L = ML^2T^{-2} \text{ (Joule)}$$

- Power:

$$P = W/t \Rightarrow [P] = [W]/[t] = (ML^2T^{-2})/T = ML^2T^{-3} \text{ (kg.m}^2\text{.s}^{-3} \text{ or Watt)}$$

- Pressure:

$$P = F/S \Rightarrow [P] = [F]/[S] = (MLT^{-2})/L^2 = ML^{-1}T^{-2} \text{ (kg.m}^{-1}\text{.s}^{-2} \text{ or Pascal)}$$

**Summary :**

Physical Quantities	Symbol	Formula used	Dimension	Unit (SI)
Surface	S	l×l	L <sup>2</sup>	m <sup>2</sup>
Volume	V	l×l×l	L <sup>3</sup>	m <sup>3</sup>
Density	ρ	m/V	ML <sup>-3</sup>	Kg./m <sup>3</sup>
Frequency	F	1/T	T <sup>-1</sup>	s <sup>-1</sup> or hertz
Linear velocity	V	dx/dt	LT <sup>-1</sup>	m/s <sup>1</sup>
Angular Velocity	Ω	dθ/dt	T <sup>-1</sup>	Rd./s <sup>1</sup>
Linear Acceleration	γ	dv/dt	LT <sup>-2</sup>	m./s <sup>2</sup>
Angular Acceleration	ω	dθ'/dt	T <sup>-2</sup>	Rd./s <sup>2</sup>
Force	F	m.a	MLT <sup>-2</sup>	Newton
Work	W	F.d	ML <sup>2</sup> T <sup>-2</sup>	Joule
Energy	E	(1/2)mv <sup>2</sup>	ML <sup>2</sup> T <sup>-2</sup>	Joule
Power	P	W/t	ML <sup>2</sup> T <sup>-3</sup>	Watt
Pressure	℘	F/S	ML <sup>-1</sup> T <sup>-2</sup>	Pascal

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### Exercise 2

We have  $\left(P + \frac{a}{v^2}\right) \times (V - b) = C$

$G=A+B$  or  $G=A-B$  then  $[G]=[A]=[B]$

$[b] = [V] = L^3$

$\left[\frac{a}{v^2}\right] = [P] = \frac{[a]}{[V]^2} \Rightarrow [a] = [P] \times [V]^2 = M.L^{-1}T^{-2}.L^6 = M.L^5.T^{-2}$

$$[C] = \left[P + \frac{a}{v^2}\right] \times [V - b]$$

On the other hand:  $\left[P + \frac{a}{v^2}\right] = [p] = \left[\frac{a}{v^2}\right]$  et  $[V - b] = [V] = [b]$

Et  $[C]=[P] \times [V] = M.L^{-1}T^{-2} . L^3 = M.L^2T^{-2}$

### Exercise 3

Check the homogeneity of this formula:  $p = \rho g h_1 + h_2 F$

Such as:  $p$  a pressure,  $g$  an acceleration of gravity,  $h_1$  and  $h_2$  are heights and  $F$  a force.

$$\text{We have: } \begin{cases} [P] = ML^{-1}T^{-2} \\ [g] = LT^{-2} \\ [h_1] = [h_2] = L \\ [F] = MLT^{-2} \\ [\rho] = ML^{-3} \end{cases}$$

This expression is homogeneous if:  $[p] = [\rho g h_1] = [h_2 F]$

$$[\rho g h_1] = ML^{-3}.L.LT^{-2} = ML^{-1}T^{-2} = [P]$$

$$\text{and } [h_2 F] = ML^2T^{-2} \neq ML^{-1}T^{-2}$$

So the equation is heterogeneous (not homogeneous).

### Exercise 4

This expression is homogeneous if:  $[y] = \left[\frac{g}{2v_0^2} x^2\right] = [h]$

We have :

$$[y] = [h] = L, \left[\frac{g}{2v_0^2} x^2\right] = \begin{bmatrix} 1 \\ 2 \end{bmatrix} \begin{bmatrix} g \\ v_0^2 \end{bmatrix} [x]^2 = 1 \frac{LT^{-2}}{(LT^{-1})^2} L^2 = L,$$

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So :  $[y] = \left[ \frac{g}{2v_0^2} x^2 \right] = [h]$  is checked.

Hence the equation  $y = \frac{g}{2v_0^2} x^2 + h$  is homogeneous.

### Exercise 5

1.  $F = \frac{Gm}{r} = ?$  such that: F is a force, G a constant expressed in  $\left( \frac{m^3}{kg s^2} \right)$ , m a mass and r a length.

$$F = ma \Rightarrow [F] = [m][a] = MLT^{-2}.$$

The unit of g is  $\left( \frac{m^3}{kg s^2} \right)$  so  $[G] = \frac{L^3}{MT^2} = M^{-1}L^3T^{-2}$

$$\text{Then } \frac{G.m}{r} = \frac{[G][m]}{[r]} = \frac{M^{-1}L^3T^{-2} M}{L} = L^2T^{-2}$$

In conclusion  $[F] \neq \left[ \frac{Gm}{r} \right]$  therefore the relationship  $F = \frac{Gm}{r}$  is not valid

2.  $p = \rho g h_1 + h_2 F$  such as: P: a pressure, g: the acceleration of gravity,  $h_1$  and  $h_2$ : heights and F: a force.

To demonstrate that the relation  $p = \rho g h_1 + h_2 F$  is valid it is necessary that:

$$\begin{aligned} [p] &= [\rho g h_1] = [h_2 F] \\ p = \frac{F}{S} &\Rightarrow [p] = \frac{[F]}{[S]} = \frac{MLT^{-2}}{L^2} = ML^{-1}T^{-2} \end{aligned}$$

$$[\rho g h_1] = [\rho][g][h_1]$$

$$\text{And } [\rho] = \frac{[m]}{[V]} = ML^{-3}, [g] = [a] = LT^{-2} \text{ et } [h_1] = L$$

$$\text{So } [\rho g h_1] = [\rho][g][h_1] = ML^{-3} LT^{-2} L = ML^{-1}T^{-2}$$

$$[h_2 F] = [h_2][F] = MLT^{-2} L = ML^2T^{-2}$$

In conclusion  $[p] = [\rho g h_1] \neq [h_2 F]$  therefore the relationship  $p = \rho g h_1 + h_2 F$  is not valid.

3.  $\theta = \frac{b \sin(a)}{t \cos(c)}$ , such as: b and t have a dimension of one length.

To demonstrate that  $[\theta] = \frac{[b \sin(a)]}{[t \cos(c)]} = \frac{[b \sin(a)]}{[t \cos(c)]}$  is valid, it is necessary to demonstrate that:

$$[\theta] = \frac{[b \sin(a)]}{[t \cos(c)]} = \frac{[b \sin(a)]}{[t \cos(c)]} ?$$

We have  $[\theta] = 1$  and  $[\sin(a)] = [\cos(c)] = 1$

$$[b \sin(a)] = [b][\sin(a)] = [b] = L$$

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$$[t \cos(c)] = [t][\cos(c)] = [t] = L$$

$$\text{So } \frac{[b \sin(a)]}{[t \cos(c)]} = \frac{L}{L} = 1$$

In conclusion  $[\theta] = \frac{[b \sin(a)]}{[t \cos(c)]} = \frac{[b \sin(a)]}{[t \cos(c)]}$  is verified therefore the relation  $\theta = \frac{b \sin(a)}{t \cos(c)}$  is valid.

### Exercise 6

We have:  $F = -6\pi\eta rv$

1-  $[\eta] = ?$

$$F = -6\pi\eta rv \Rightarrow \eta = -\frac{F}{6\pi rv}$$

$$[\eta] = \frac{[F]}{[r][v]} \text{ with } \begin{cases} [r] = L \\ [F] = MLT^{-2} \\ [v] = LT^{-1} \\ [-6\pi] = 1 \end{cases}$$

Where

$$[\eta] = \frac{MLT^{-2}}{L \cdot LT^{-1}} = ML^{-1}T^{-1}$$

2- We have:  $v = a \left(1 - \exp\left(-\frac{t}{b}\right)\right)$

we're looking for the dimension of  $[a]$  and  $[b]$ :

The argument of the exponential is therefore dimensionless:

$$\text{so } [v] = LT^{-1} = [a] \Rightarrow [a] = LT^{-1}$$

$$\left[\exp\left(-\frac{t}{b}\right)\right] = 1 \Rightarrow \left[-\frac{t}{b}\right] = \left[-1 \cdot \frac{t}{b}\right] = [-1] \left[\frac{t}{b}\right] = \left[\frac{t}{b}\right] = 1$$

$$\Rightarrow \left[\frac{t}{b}\right] = \frac{[t]}{[b]} = 1$$

$$[b] = [t] = T$$

### Exercise 7 :

We have :  $v = k\rho^x \chi^y$  so  $[v] = [k][\rho]^x [\chi]^y$ .

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$$\text{And } \begin{cases} \text{the speed } v : [v] = LT^{-1} \\ \text{the constant } k : [k] = 1 \\ \text{the density } \rho : [\rho] = ML^{-3} \\ \chi \text{ is homogeneous in contrast to a pressure: } [\chi] = \frac{1}{[\rho]} = M^{-1}LT^{+2} \end{cases}$$

$$\Rightarrow [v] = LT^{-1} = (ML^{-3})^x (M^{-1}LT^2)^y$$

$$\Rightarrow M^0 L T^{-1} = M^x L^{-3x} M^{-1y} L^y T^{2y}$$

$$\Rightarrow M^0 L^1 T^{-1} = M^{x-y} L^{-3x+y} T^{2y}$$

**By identification:**

$$\begin{cases} x - y = 0 \\ -3x + y = 1 \\ 2y = -1 \end{cases} \Rightarrow \begin{cases} x = y = -\frac{1}{2} \\ y = -1/2 \end{cases} \Rightarrow v = k \rho^{-1/2} \chi^{-1/2} = \frac{k}{\sqrt{\rho \chi}}$$

Then:

$$v = \frac{k}{\sqrt{\rho \chi}}$$

**Exercise 8:**

$f = K F^a L^b \rho^c$ ; This function is therefore homogeneous  $[f] = [k][F]^a [L]^b [\rho]^c$

$$\text{with: } \begin{cases} [F] = [m \cdot a] = [m][a] = M \cdot LT^{-2} \\ [L] = L \text{ et } [k] = 1 \\ [\rho] = \left[\frac{m}{V}\right] = ML^{-3} \\ [f] = T^{-1} \end{cases}$$

so:  $[f] = (MLT^{-2})^a (L)^b (ML^{-3})^c = T^{-1}$

$$\Rightarrow M^0 L^0 T^{-1} = M^{a+c} L^{a+b-3c} T^{-2a}$$

$$\text{By identification: } \begin{cases} a + c = 0 \\ a + b - 3c = 0 \\ -2a = -1 \end{cases}$$

$$\Rightarrow \begin{cases} a = 1/2 \\ b = -a + 3c = -\frac{4}{2} = -2 \\ c = -a = -1/2 \end{cases}$$

$$F = K F^{1/2} L^{-2} \rho^{-1/2} = K \sqrt{F} \frac{1}{L^2 \sqrt{\rho}} \quad \text{So} \quad f = k \frac{\sqrt{F}}{L^2 \sqrt{\rho}}$$

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### Exercise 9

1- The period of a pendulum is written :

$$T = K\eta^x R^y \rho^z \text{ such } [\eta] = ML^{-1}T^{-1}$$

Suppose the relationship is homogeneous so  $[T] = [k][\eta]^x[R]^y[\rho]^z$

$$\text{With } \begin{cases} [\eta] = ML^{-1}T^{-1} \\ [R] = L \text{ Let } [k] = 1 \\ [\rho] = \left[ \frac{m}{V} \right] = \frac{M}{L^3} = ML^{-3} \\ [T] = T \end{cases}$$

$$\text{So } [T] = (ML^{-1}T^{-1})^x L^y (ML^{-3})^z = T$$

$$(A^X \cdot A^Y = A^{X+Y})$$

$$\Rightarrow T = M^x L^{-x} T^{-x} L^y M^z L^{-3z}$$

$$\Rightarrow M^0 L^0 T^1 = M^{x+z} L^{-x+y-3z} T^{-x}$$

$$\text{by identification: } \begin{cases} x + z = 0 \\ -x + y - 3z = 0 \\ -x = 1 \end{cases}$$

$$\Rightarrow \begin{cases} x = -1 \\ y = x + 3z = 2 \\ z = -x = 1 \end{cases}$$

$$\Rightarrow T = K\eta^{-1} R^2 \rho^1$$

$$\text{So } T = k \frac{\rho R^2}{\eta}$$

2- The relative uncertainty on  $T = f(\Delta\eta, \Delta R, \Delta m)$  ?

$$T = \frac{K\rho R^2}{\mu} \text{ With } \rho = \frac{m}{V} = \frac{m}{\frac{4}{3}\pi R^3} = \frac{3m}{4\pi R^3} \text{ so } T = \frac{3Km}{4\pi R\mu}$$

$$\Rightarrow \log T = \log\left(\frac{3mK}{4\pi R\mu}\right) = \log 3K + \log(m) - \log(4\pi) - \log(R) - \log(\mu)$$

$$\Rightarrow \frac{dT}{T} = \frac{dm}{m} - \frac{dR}{R} - \frac{d\mu}{\mu}$$

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$$\Rightarrow \frac{\Delta T}{T} = \left| \frac{\Delta m}{m} \right| + \left| -\frac{\Delta R}{R} \right| + \left| -\frac{\Delta \mu}{\mu} \right|$$

m, R, and  $\mu$  are positive quantities, hence:

$$\Rightarrow \frac{\Delta T}{T} = \frac{\Delta m}{m} + \frac{\Delta R}{R} + \frac{\Delta \mu}{\mu}$$

### Exercise 10

1- We have  $G = \frac{T^2 g a}{4\pi^2} - a^2$ , with: 
$$\begin{cases} [T] = T \\ [a] = L \\ [g] = LT^{-2} \\ [4\pi] = 1 \end{cases}$$

The dimension of G :

$$[G] = \left[ \frac{T^2 g a}{4\pi^2} \right] = [a^2] = L^2$$

Or  $[G] = \left[ \frac{T^2 g a}{4\pi^2} \right] = \frac{[T]^2 [g] [a]}{[4\pi^2]} = \frac{T^2 L T^{-2} L}{1} = L^2$  So  $[G]=L^2$

where G has a unit area ( $m^2$ ).

2- Calculation of the absolute uncertainty on G as a function of  $\Delta T$  and  $\Delta a$ :

$$\begin{aligned} G &= \frac{T^2 g a}{4\pi^2} - a^2 \\ \Rightarrow dG &= \left( \frac{\partial G}{\partial T} \right) dT + \left( \frac{\partial G}{\partial a} \right) da \\ \Rightarrow dG &= \left( \frac{2g a T}{4\pi^2} \right) dT + \left( \frac{g T^2}{4\pi^2} - 2a \right) da \end{aligned}$$

So the absolute uncertainty on G is:

$$\Delta G = \left| \frac{g a T}{2\pi^2} \right| \Delta T + \left| \frac{g T^2}{4\pi^2} - 2a \right| \Delta a$$

### Exercise 11

$$n = \sqrt{N^2 - \sin^2 \alpha} :$$

1. Calculation of the absolute uncertainty on n.

The total differential of n is written:

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$$dn = \frac{\partial n}{\partial N} dN + \frac{\partial n}{\partial \alpha} d\alpha$$

The partial differentials of the function “n” with respect to the two variables N and  $\alpha$  are:

$$\frac{\partial n}{\partial N} = \frac{2N}{2\sqrt{N^2 - \sin^2\alpha}}$$

And

$$\frac{\partial n}{\partial \alpha} = \frac{-2\sin\alpha \cdot \cos\alpha}{2\sqrt{N^2 - \sin^2\alpha}}$$

So

$$dn = \frac{N}{\sqrt{N^2 - \sin^2\alpha}} dN + \frac{-\sin\alpha \cdot \cos\alpha}{\sqrt{N^2 - \sin^2\alpha}} d\alpha$$

Then the absolute uncertainty on n is:

$$\Delta n = \frac{N}{\sqrt{N^2 - \sin^2\alpha}} \Delta N + \frac{|-\sin\alpha \cdot \cos\alpha|}{\sqrt{N^2 - \sin^2\alpha}} \Delta \alpha$$

Let for  $\alpha < \frac{\pi}{2}$  the relative uncertainty on n is:

$$\frac{\Delta n}{n} = \frac{N\Delta N + \sin\alpha \cdot \cos\alpha \cdot \Delta \alpha}{(|N^2 - \sin^2\alpha|)}$$

### Exercise 12

1- The dimension of k :

$$\text{We have } \begin{cases} [p] = M \cdot L \cdot T^{-2} \\ [S] = L^2 \\ [k] = 1 \\ [v] = L \cdot T^{-1} \end{cases} \text{ and } k = \frac{p}{v^2 \cdot s} \Rightarrow [k] = \frac{[p]}{[v]^2 \cdot [s]}$$

$$\Rightarrow [k] = [p] \cdot [v]^{-2} \cdot [s]^{-1} \Rightarrow [k] = M \cdot L^{-3}$$

2- N.A :  $v = \sqrt{\frac{P}{K \cdot S}} = 3.097 \text{ m/s}$

3-  $\frac{\Delta P}{P} = 2\% = 0.02$  and  $\frac{\Delta S}{S} = 3\% = 0.03$

The logarithmic method is used to calculate the relative uncertainty on  $v$  :

$$v = \sqrt{\frac{P}{K \cdot S}} \Rightarrow \log v = \log \sqrt{\frac{P}{K \cdot S}} = \frac{1}{2} \log P - \frac{1}{2} \log k - \frac{1}{2} \log S$$



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$$\Rightarrow d \log v = \frac{1}{2} d \log P - \frac{1}{2} d \log k - \frac{1}{2} d \log S$$

$$\frac{dv}{v} = \frac{1}{2} \frac{dp}{p} - \frac{1}{2} \frac{dS}{S} \Rightarrow \frac{\Delta v}{v} = \frac{1}{2} \left| \frac{\Delta p}{p} \right| + \frac{1}{2} \left| -\frac{\Delta S}{S} \right| \Rightarrow \frac{\Delta v}{v} = \frac{1}{2} \frac{\Delta p}{p} + \frac{1}{2} \frac{\Delta S}{S}$$

A.N:  $\frac{\Delta v}{v} = 0.025$

Absolute uncertainty on v is given by:

$$\Delta v = v \cdot \frac{\Delta v}{v} = v * \left( \frac{1}{2} \frac{\Delta p}{p} + \frac{1}{2} \frac{\Delta S}{S} \right) = 0.077 \text{ m/s}$$

Hence the condensed writing of v is given by :  $v = (3.097 \pm 0.077) \text{ m/s}$

### Exercise 13

1-  $H = \frac{(2 \cdot \sigma \cdot \cos \alpha)}{(R \cdot g \cdot \rho)} \Rightarrow \sigma = \frac{HR\rho g}{2 \cos \alpha}$

$$\text{Hence } [\sigma] = \frac{[H][R][\rho][g]}{[2][\cos]}$$

$$\text{On a } \begin{cases} [H] = L \\ [R] = L \\ [\rho] = ML^{-3} \\ [g] = LT^{-2} \\ [2] = [\cos \alpha] = 1 \end{cases}$$

$$\Rightarrow [\sigma] = L \cdot L \cdot M \cdot L^{-3} \cdot LT^{-2} = MT^{-2}$$

So  $[\sigma] = MT^{-2}$

2- Let's calculate relative uncertainty on  $\sigma$ ,  $\Delta\sigma/\sigma = f(\Delta m, \Delta R, \Delta g, \Delta \alpha)$  :

We have,  $\sigma = \frac{HR\rho g}{2 \cos \alpha}$  or  $\rho = \text{the density} = \frac{M}{V} = \frac{M}{H\pi R^2}$

So:  $\sigma = \frac{HR \frac{M}{H\pi R^2} g}{2 \cos \alpha} = \frac{Mg}{2\pi R \cos \alpha}$

a- The Logarithmic method :

$$\log \sigma = \log \left( \frac{Mg}{2\pi R \cos \alpha} \right) = \log Mg - \log(2\pi R \cos \alpha)$$

$$\Rightarrow \log \sigma = \log M + \log g - \log 2\pi - \log R - \log \cos \alpha$$

$$\Rightarrow d \log \sigma = d \log M + d \log g - d \log 2\pi - d \log R - d \log \cos \alpha$$

## Chapter I: Dimensional Analysis and Uncertainty Calculation

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$$\frac{d\sigma}{\sigma} = \frac{dM}{M} + \frac{dg}{g} - \frac{dR}{R} - \frac{d\cos\alpha}{\cos\alpha}$$

$$= \frac{d\sigma}{\sigma} = \frac{dM}{M} + \frac{dg}{g} - \frac{dR}{R} - \frac{\sin\alpha}{\cos\alpha} d\alpha$$

$$\Rightarrow \frac{\Delta\sigma}{\sigma} = \frac{\Delta M}{M} + \frac{\Delta g}{g} + \frac{\Delta R}{R} + |\tan\alpha|\Delta\alpha$$

b- The total differential method:

$$d\sigma = \frac{\partial\sigma}{\partial M} dM + \frac{\partial\sigma}{\partial R} dR + \frac{\partial\sigma}{\partial g} dg + \frac{\partial\sigma}{\partial\alpha} d\alpha$$

$$\text{With } \begin{cases} \frac{\partial\sigma}{\partial g} = \frac{M}{2\pi R \cos\alpha} \\ \frac{\partial\sigma}{\partial M} = \frac{g}{2\pi R \cos\alpha} \\ \frac{\partial\sigma}{\partial R} = \frac{-Mg}{2\pi R^2 \cos\alpha} \\ \frac{\partial\sigma}{\partial\alpha} = \frac{Mg}{2\pi R} \cdot \frac{\sin\alpha}{\cos^2\alpha} \end{cases}$$

Hence

$$d\sigma = \frac{g}{2\pi R \cos\alpha} dM + \frac{-Mg}{2\pi R^2 \cos\alpha} dR + \frac{M}{2\pi R \cos\alpha} dg + \frac{Mg}{2\pi R} \cdot \frac{\sin\alpha}{\cos^2\alpha} d\alpha$$

$$\Rightarrow d\sigma = \left( \frac{Mg}{2\pi R \cos\alpha} \right) \left[ \left( \frac{1}{M} \right) dM + \left( \frac{-1}{R} \right) dR + \left( \frac{1}{g} \right) dg + \left( \frac{\sin\alpha}{\cos\alpha} \right) d\alpha \right]$$

$$\underbrace{\hspace{10em}}_{= \sigma}$$

$$\Rightarrow \frac{d\sigma}{\sigma} = \frac{dM}{M} + \frac{-dR}{R} + \frac{dg}{g} + \frac{\sin\alpha}{\cos\alpha} d\alpha$$

$$\text{So } \frac{\Delta\sigma}{\sigma} = \frac{\Delta M}{M} + \frac{\Delta R}{R} + \frac{\Delta g}{g} + |\tan\alpha|\Delta\alpha$$

### Exercise 14

$$\text{We have } f = \frac{1}{2\pi\sqrt{L\cdot C}}$$

We will calculate the absolute uncertainty on f.

#### 1<sup>st</sup> Method: The total differential

$$df = \left( \frac{\partial f}{\partial L} \right) dL + \left( \frac{\partial f}{\partial C} \right) dC$$

$$\text{With } \begin{cases} \left(\frac{\partial f}{\partial L}\right) = \frac{1}{2\pi} \cdot \frac{-C}{2LC\sqrt{LC}} = \frac{1}{2\pi} \cdot \left(\frac{-1}{2L\sqrt{LC}}\right) \\ \left(\frac{\partial f}{\partial C}\right) = \frac{1}{2\pi} \cdot \frac{-1}{2C\sqrt{LC}} \end{cases}$$

$$\text{So: } df = \frac{1}{2\pi} \cdot \left(\frac{-1}{2L\sqrt{LC}}\right) dL + \frac{1}{2\pi} \cdot \frac{-1}{2C\sqrt{LC}} dC$$

$$\Rightarrow df = \left(\frac{1}{2\pi\sqrt{LC}}\right) \cdot \left[\left(\frac{-1}{2L}\right) dL + \frac{-1}{2C} dC\right]$$

$$\Rightarrow \frac{df}{f} = \left[\left(\frac{-1}{2L}\right) dL + \frac{-1}{2C} dC\right]$$

$$\Rightarrow \frac{\Delta f}{f} = \left|\frac{-1}{2L}\right| \Delta L + \left|\frac{-1}{2C}\right| \Delta C$$

$$\Rightarrow \frac{\Delta f}{f} = \frac{1}{2L} \Delta L + \frac{1}{2C} \Delta C \text{ is the relative uncertainty on } f.$$

$$\text{And } \Delta f = \left(\frac{1}{2\pi\sqrt{LC}}\right) \cdot \left(\left|\frac{-1}{2L}\right| \Delta L + \left|\frac{-1}{2C}\right| \Delta C\right)$$

$$\Delta f = \left(\frac{1}{2\pi\sqrt{LC}}\right) \cdot \left(\frac{1}{2L} \Delta L + \frac{1}{2C} \Delta C\right) \text{ is the absolute uncertainty on } f.$$

### 2<sup>nd</sup> Method: The logarithmic differential

$$\text{We have } f = \frac{1}{2\pi\sqrt{LC}} \Rightarrow \log f = \log \frac{1}{2\pi} - \frac{1}{2} \log L - \frac{1}{2} \log C$$

$$\Rightarrow d \log f = d \log \frac{1}{2\pi} - \frac{1}{2} d \log L - \frac{1}{2} d \log C$$

$$\Rightarrow \frac{df}{f} = -\frac{1}{2} \frac{dL}{L} - \frac{1}{2} \frac{dC}{C} \Rightarrow \frac{\Delta f}{f} = \left|-\frac{1}{2}\right| \frac{\Delta L}{L} + \left|-\frac{1}{2}\right| \frac{\Delta C}{C}$$

Hence the relative uncertainty on f is:

$$\frac{\Delta f}{f} = \frac{1}{2} \frac{\Delta L}{L} + \frac{1}{2} \frac{\Delta C}{C}$$

And absolute uncertainty about f:

$$\Delta f = f \left(\frac{1}{2} \frac{\Delta L}{L} + \frac{1}{2} \frac{\Delta C}{C}\right)$$

## Chapter I: Dimensional Analysis and Uncertainty Calculation

### Exercise 15

A) The expression of the frequency  $f$ :

$$\text{We have: } \begin{cases} [f] = T^{-1} \\ [F] = [m][a] = MLT^{-2} \\ [L] = L \\ [\rho] = \frac{[m]}{[V]} = ML^{-3} \\ [K] = 1 \end{cases}$$

We consider that the formula of "f" is homogeneous.

$$\begin{aligned} [f] = [K][F]^a[L]^b[\rho]^c &\Rightarrow T^{-1} = (MLT^{-2})^a(L)^b(ML^{-3})^c \\ &\Rightarrow T^{-1} = M^{a+c} L^{a+b-3c} T^{-2a} \end{aligned}$$

By identification

$$\begin{cases} a + c = 0 \\ a + b - 3c = 0 \\ -1 = -2a \end{cases} \Rightarrow \begin{cases} a = \frac{1}{2} \\ b = -2 \\ c = -\frac{1}{2} \end{cases} \Rightarrow f = KF^{\frac{1}{2}}L^{-2}\rho^{-\frac{1}{2}}$$

or

$$f = \frac{K}{L^2} \sqrt{\frac{F}{\rho}}$$

B) Let us calculate the absolute uncertainty  $\Delta f$  as a function of  $\Delta D$  and  $\Delta a$ .

$$\begin{aligned} f &= \frac{D^2 - a^2}{4D} \\ &\Rightarrow df = \left(\frac{\partial f}{\partial D}\right) dD + \left(\frac{\partial f}{\partial a}\right) da \\ &\Rightarrow df = \left(\frac{2D(4D) - 4(D^2 - a^2)}{16D^2}\right) dT + \left(\frac{-2a}{4D}\right) da \\ &\Rightarrow df = \left(\frac{D^2 - a^2}{4D^2}\right) dT + \left(\frac{-2a}{4D}\right) da \end{aligned}$$

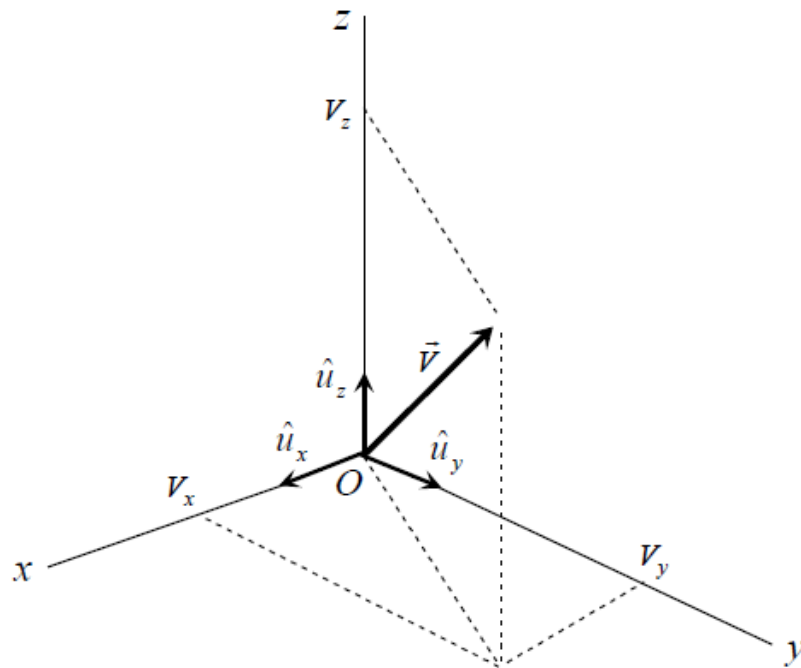
So the absolute uncertainty on  $f$  is:

$$\begin{aligned} \Delta f &= \left|\frac{D^2 - a^2}{4D^2}\right| \Delta T + \left|\frac{-2a}{4D}\right| \Delta a \\ &\Rightarrow \Delta f = \left|\frac{D^2 - a^2}{4D^2}\right| \Delta T + \frac{a}{2D} \Delta a \end{aligned}$$

# COURSE OF MECHANICS OF THE MATERIAL POINT

## *Chapter II: Vector Analysis*

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**Glossary**

<b>In English</b>	<b>In French</b>	<b>In Arabic</b>
Vector	Vecteur	شعاع
Vector quantity	Grandeurs vectorielles	مقدار شعاعي
Displacement vectors	Vecteur de déplacement	شعاع التنقل الموضعي
Position vector	Vecteur position	شعاع موضعي
Velocity vector	Le vecteur vitesse	شعاع السرعة
Acceleration vector	Le vecteur accélération	شعاع التسارع
Unit vector	Vecteur unité	شعاع الوحدة
Magnitude	Module, l'intensité ou la norme du vecteur	طويلة شعاع
Direction	La direction, le sens	اتجاه شعاع
Addition	Addition	جمع
Subtraction	Soustraction	الطرح
Scalar multiplication	Multiplication scalaire	جداء سلمي
Coordinate systems	Systèmes de coordonnées	نظام احداثيات
Force vectors	Le vecteur force	شعاع القوة
Opposite direction	La direction opposée	اتجاه معاكس
Zero vector	Vecteur nul	شعاع منعدم
Equal vectors	Vecteurs égaux	اشعة متساوية
Vector parallel	Vecteur parallèle	شعاع موازي
Vector perpendicular	Vecteur perpendiculaire	شعاع معامد
Free vector	Vecteurs libres	اشعة حرة
Sliding vector	Vecteurs glissants	اشعة منزلقة
Linked vectors	Vecteurs liés	اشعة متصلة
Opposite vector	Vecteurs opposés	شعاعان متعاكسان
Algebraic value	Valeur algébrique	قيمة جبرية
Equal magnitudes	Même amplitude (module)	نفس الطويلة
Graphic representations	Représentation graphique	تمثيل في الفضاء
The sum of two vectors	la somme de deux vecteurs	مجموع شعاعين
The axis	Un axe	محور

## Chapter II: Vector Analysis

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The xy-plane	Un plan (Oxy)	معلم مستوي
The scalar product	Produit scalaire	جاء سلمي
The vector product	Produit vectoriel	جاء شعاعي
The mixed product	Produit mixte	جاء مختلط
The parallelepiped	Un parallélépipède	متوازي المستطيلات
The parallelogramme	Un parallélogramme	متوازي الاضلاع
The norme	La norme	طول الشعاع

### 1. Introduction

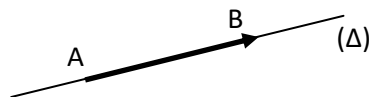
Vectors are fundamental mathematical entities used to represent quantities that have both magnitude and direction. Unlike scalars, which only have magnitude (e.g., distance, time, temperature), vectors provide a more comprehensive description of physical quantities by including information about their orientation or direction.

In other words, in physics, two types of quantities are used: scalar quantities and vector quantities:

- Scalar quantity **المقدار السلمي** : defined by a number (a scalar) and an appropriate unit such as: volume, mass, temperature, time ...
- Vector quantity **المقدار الشعاعي**: this is a quantity defined by a scalar, a unit and a direction such as : Displacement vector, velocity  $\vec{v}$ , weight  $\vec{p}$ , electric field  $\vec{E}$  ...

### 2. Definition

Vectors are physical or mathematical quantities carrying two properties: magnitude and direction. It is an oriented segment. Symbolically, a vector is usually represented by an arrow.



- Origin (**المبدأ**): presents the point of application "A".
- Support (**الحامل**): the straight line that carries the vector ( $\Delta$ ).
- Direction (**الاتجاه**): Vectors have a specific direction or orientation in space, often indicated by angles or coordinate systems (from A to B).
- Modulus (**الطويلة**): The size or length of a vector represents its magnitude. This is typically represented by a positive numerical value gives the algebraic value of the vector  $\overrightarrow{AB}$  noted.

### 3. Vector types

- **Free vector**: the origin is not fixed.
- **Sliding vector**: the support is fixed, but the origin is not.
- **Linked vectors**: the origin is fixed.



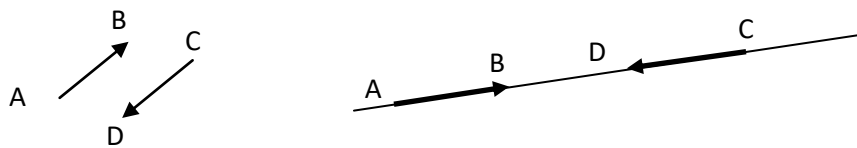
## Chapter II: Vector Analysis

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- **Equal vectors:** if they have the same direction, the same support or parallel supports and the same modulus.



- **Opposite vector:** if they have the same support or parallel supports, the same modulus but the direction is opposite.

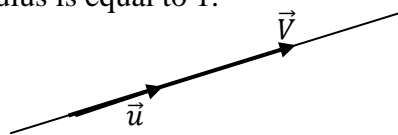


### 4. Unit Vector شعاع الوحدة

A vector is said to be unitary if its modulus is equal to 1.

We write:  $|\vec{u}|=1$

and  $\vec{V} = |\vec{V}|\vec{u}=V \cdot \vec{u}$



### 5. Algebraic Measurement

Consider an axis ( $\Delta$ ) bearing points O and A. O is the origin, and the abscissa of point A is the algebraic measure of the vector  $\vec{OA}$ .

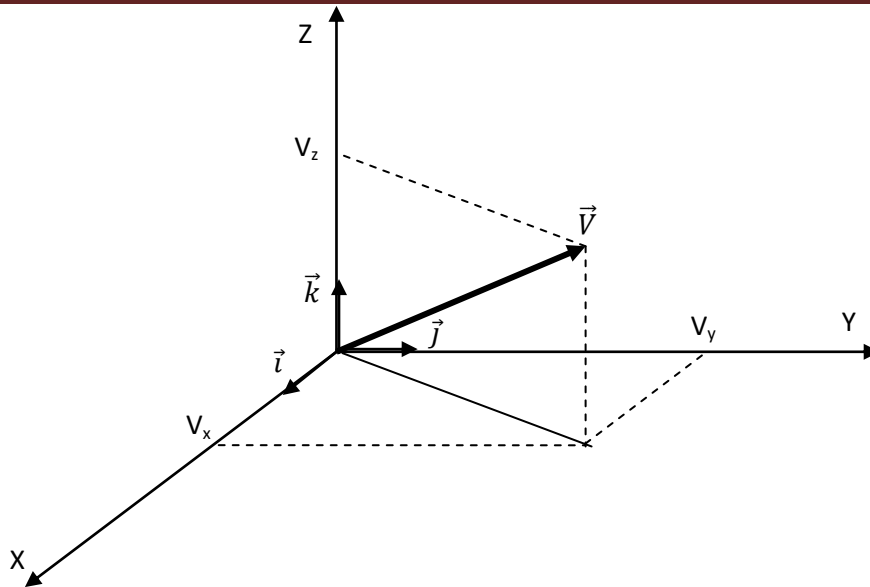


### 6. Components of a Vector مركبات شعاع

The coordinates of a vector in space, represented in an orthonormal base frame  $R(O, \vec{i}, \vec{j}, \vec{k})$  are :  $V_x, V_y$  et  $V_z$  such that:

$$\vec{V} = V_x \vec{i} + V_y \vec{j} + V_z \vec{k}$$

Where a **position vector**  $\vec{V} = \vec{OM}$  is a vector used to determine the position of a point M in space, relative to a fixed reference point O which, typically, is chosen to be the origin of our coordinate system.



The modulus of the vector  $\vec{V}$  is :  $V = \sqrt{V_x^2 + V_y^2 + V_z^2}$

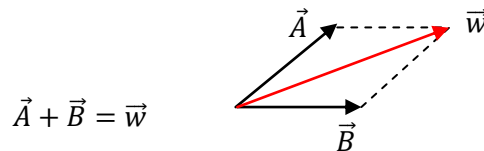
In cartesian coordinates, a vector is written as:

$$\vec{V} = x\vec{i} + y\vec{j} + z\vec{k} \Rightarrow V = \|\vec{V}\| = \sqrt{x^2 + y^2 + z^2}$$

### 7. Elementary Operations on Vectors

#### 7.1. Vector addition

The sum of two vectors  $\vec{A}$  and  $\vec{B}$  is  $\vec{w}$ , obtained using the parallelogram:

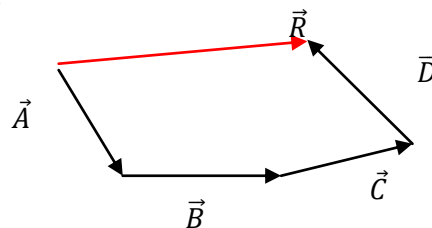


Let two vectors  $\vec{A}$  and  $\vec{B}$ :  $\vec{A} = x\vec{i} + y\vec{j} + z\vec{k}$  and  $\vec{B} = x'\vec{i} + y'\vec{j} + z'\vec{k}$

$$\vec{A} \begin{pmatrix} x \\ y \\ z \end{pmatrix} \text{ and } \vec{B} \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} \text{ so } \vec{A} + \vec{B} = \vec{w} = (x + x')\vec{i} + (y + y')\vec{j} + (z + z')\vec{k}$$

#### Note :

1. For several vectors:  $\vec{A} + \vec{B} + \vec{C} + \vec{D} = \vec{R}$



## Chapter II: Vector Analysis

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2. Properties :

$$\vec{A} + \vec{B} = \vec{B} + \vec{A}, \quad (\vec{A} + \vec{B}) + \vec{C} = \vec{A} + (\vec{B} + \vec{C}), \quad \vec{A} - \vec{B} = \vec{A} + (-\vec{B})$$

3. Charles relationship:

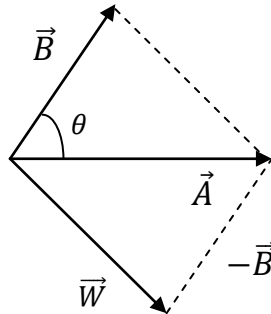
Or the three points: A, B and C, we have:  $\overline{AB} + \overline{BC} = \overline{AC}$

### 7.2. Subtracting two vectors

This is an anticommutative operation such that:  $\vec{W} = \vec{A} - \vec{B} = \vec{A} + (-\vec{B})$

Let two vectors:  $\vec{A}$  and  $\vec{B}$ ,  $\vec{A} = x\vec{i} + y\vec{j} + z\vec{k}$  and  $\vec{B} = x'\vec{i} + y'\vec{j} + z'\vec{k}$

$\vec{A} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$  and  $\vec{B} \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix}$  so  $\vec{A} - \vec{B} = \vec{w} = (x - x')\vec{i} + (y - y')\vec{j} + (z - z')\vec{k}$



### 7.3. Product of a vector and a scalar

The product of a vector  $\vec{v}$  by a scalar  $\alpha$  is the vector  $\alpha\vec{v}$ , this vector has the same support as  $\vec{v}$ .

The two vectors ( $\vec{v}$  and  $\alpha\vec{v}$ ) have the same direction if  $\alpha > 0$  and they are opposite supports if  $\alpha < 0$ .

$$\alpha\vec{v} = \alpha \begin{pmatrix} x \\ y \\ z \end{pmatrix} = \alpha x\vec{i} + \alpha y\vec{j} + \alpha z\vec{k}$$

**Notes:**  $|\alpha\vec{v}| = |\alpha||\vec{v}|$ ,  $\alpha(\vec{u} + \vec{v}) = \alpha\vec{u} + \alpha\vec{v}$  and  $(\alpha + \beta)\vec{u} = \alpha\vec{u} + \beta\vec{u}$

## 8. Products

### 8.1. Scalar product الجداء السلمي

Given two vectors  $\vec{A}$  and  $\vec{B}$  making an angle  $\theta$  between them, the scalar product  $\vec{A} \cdot \vec{B} = m$  with  $m$  is a scalar such that:

$$\vec{A} \cdot \vec{B} = m = |\vec{A}| \cdot |\vec{B}| \cos(\vec{A}, \vec{B})$$

With :  $(\vec{A}, \vec{B}) = \theta$

**Note :** The properties of the scalar product are:

- The scalar product is commutative  $\vec{A} \cdot \vec{B} = \vec{B} \cdot \vec{A}$
- The scalar product isn't associative  $\vec{V}_1 \cdot (\vec{V}_2 \cdot \vec{V}_3)$ , doesn't exist, because the result would be a vector.
- $\vec{A} \cdot \vec{B} = 0$  when both vectors are perpendicular ( $\vec{A} \perp \vec{B}$ ).
- If  $\vec{A} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$  and  $\vec{B} \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix}$  so  $\vec{A} \cdot \vec{B} = x \cdot x' + y \cdot y' + z \cdot z'$

### 8.2. Vector product الجداء الشعاعي

The vector product of two vectors  $\vec{A}$  and  $\vec{B}$  is a vector  $\vec{C}$  and is written as:

$$\vec{C} = \vec{A} \wedge \vec{B}$$

To calculate the vector product of two vectors  $\vec{A} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$  and  $\vec{B} \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix}$  we have :

$$\vec{A} \wedge \vec{B} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ x & y & z \\ x' & y' & z' \end{vmatrix} = \vec{i} \begin{vmatrix} y & z \\ y' & z' \end{vmatrix} - \vec{j} \begin{vmatrix} x & z \\ x' & z' \end{vmatrix} + \vec{k} \begin{vmatrix} x & y \\ x' & y' \end{vmatrix} = \vec{C}$$

$$\vec{A} \wedge \vec{B} = \vec{i}(yz' - zy') - \vec{j}(xz' - zx') + \vec{k}(xy' - yx') = \vec{C}$$

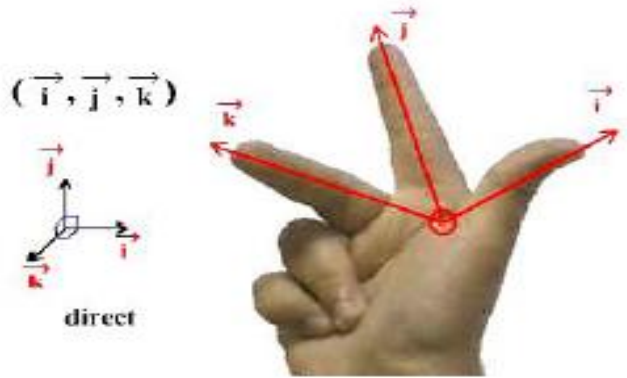
So the modulus of the vector product can be given by another method such as:

$$W = \sqrt{(yz' - zy')^2 + (xz' - zx')^2 + (xy' - yx')^2}$$

**Characteristics of vector  $\vec{C}$  :**

**The support:**  $\vec{C}$  is perpendicular to the plane formed by the two vectors  $\vec{A}$  and  $\vec{B}$ .

**The direction:** The three vectors  $\vec{A}$ ,  $\vec{B}$  and  $\vec{C}$  form a direct trihedron. The direction is given by the rule of the three fingers of the right hand.



**The modulus :**

$$|\vec{C}| = |\vec{A}| \cdot |\vec{B}| \sin(\angle \vec{A}, \vec{B})$$

The modulus of the vector product corresponds to the area (the surface *مساحة*) of the parallelogram (*متوازي الاضلاع*) formed by the two vectors  $\vec{A}$  and  $\vec{B}$ .

**Example:**

In an orthonormal Cartesian coordinate base  $(\vec{i}, \vec{j}, \vec{k})$  :

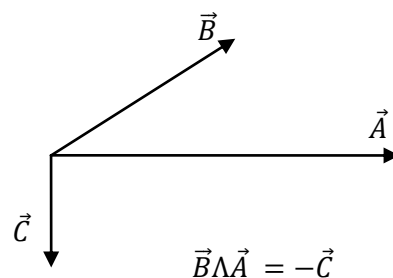
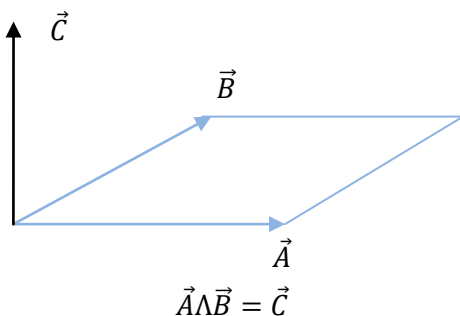
$$\vec{i} \wedge \vec{j} = \vec{k}, \vec{j} \wedge \vec{k} = \vec{i}, \vec{k} \wedge \vec{i} = \vec{j}. \text{ On the other hand } \vec{i} \wedge \vec{k} = -\vec{j}$$

**Notes :** The properties of the vector product are:

- The vector product is not commutative (Anticommutative).
- Not associative :  $\vec{V}_1 \wedge (\vec{V}_2 \wedge \vec{V}_3) \neq (\vec{V}_1 \wedge \vec{V}_2) \wedge \vec{V}_3$ .
- Distributive with respect to vector sum:  $\vec{A} \wedge (\vec{B}_1 + \vec{B}_2) = \vec{A} \wedge \vec{B}_1 + \vec{A} \wedge \vec{B}_2$

But:  $\vec{V}_1 \wedge (\vec{V}_2 + \vec{V}_3) \neq (\vec{V}_1 \wedge \vec{V}_2) + (\vec{V}_1 \wedge \vec{V}_3)$

- $\vec{A} \wedge \vec{B} = -\vec{B} \wedge \vec{A}$  because  $\sin(\angle \vec{A}, \vec{B}) = -\sin(\angle \vec{B}, \vec{A})$



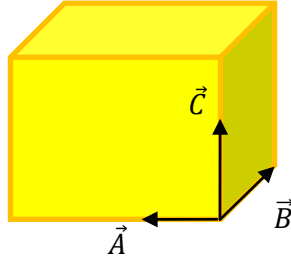
- $\vec{A} \wedge \vec{B} = \vec{0}$  when the two vectors are parallel ( $\vec{A} \parallel \vec{B}$ ).

### 8.3. Mixed product

The mixed product of three vectors  $\vec{A}, \vec{B}$  and  $\vec{C}$  is a scalar quantity  $\mathbf{m}$  such that:

$$m = (\vec{A} \wedge \vec{B}) \cdot \vec{C}$$

Where  $\mathbf{m}$  represents the volume of the parallelepiped (حجم متوازي المستطيلات) constructed by the three vectors :



**Note:** The mixed product is commutative,  $(\vec{A} \wedge \vec{B}) \cdot \vec{C} = \vec{A} \cdot (\vec{B} \wedge \vec{C}) = (\vec{C} \wedge \vec{A}) \cdot \vec{B}$

### 9. Derivative of a vector

Let the vector  $\vec{A} = x\vec{i} + y\vec{j} + z\vec{k}$  which varies with time:

Its first derivative in relation to time is:

$$\vec{A}' = \frac{d\vec{A}}{dt} = \frac{dx}{dt}\vec{i} + \frac{dy}{dt}\vec{j} + \frac{dz}{dt}\vec{k}$$

The second derivative is:

$$\vec{A}'' = \frac{d^2\vec{A}}{dt^2} = \frac{d^2x}{dt^2}\vec{i} + \frac{d^2y}{dt^2}\vec{j} + \frac{d^2z}{dt^2}\vec{k}$$

**Note :**

- Derivative of a scalar product  $(\vec{A} \cdot \vec{B})' = \vec{A}' \cdot \vec{B} + \vec{A} \cdot \vec{B}'$
- If  $\vec{B}$  is constant  $(\vec{A} \cdot \vec{B})' = \vec{A}' \cdot \vec{B}$
- $(\vec{A}^2)' = 0$  because  $(\vec{A}^2)' = 2\vec{A}' \cdot \vec{A} = 0$
- The derivative vector is perpendicular to the vector.
- A vector is written as  $\vec{A} = |\vec{A}|\vec{u} = A\vec{u}$ , if  $\vec{u}$  is a variable vector, then  $\vec{A}' = A'\vec{u} + A\vec{u}'$ .

**Example:** The position vector on Cartesian Coordinates is written as:

$$\vec{A} = x\vec{i} + y\vec{j} + z\vec{k}$$

The velocity vector in Cartesian Coordinates is written as:

## Chapter II: Vector Analysis

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$$\vec{v} = \frac{d\overrightarrow{OM'}}{dt} = \frac{dx}{dt}\vec{i} + \frac{dy}{dt}\vec{j} + \frac{dz}{dt}\vec{k}$$

The acceleration vector in Cartesian Coordinates is written as:

$$\vec{a} = \frac{d^2\overrightarrow{OM'}}{dt^2} = \frac{d^2x}{dt^2}\vec{i} + \frac{d^2y}{dt^2}\vec{j} + \frac{d^2z}{dt^2}\vec{k}$$

### Proposed exercises about chapter II

#### Exercise 1

$\vec{i}$ ,  $\vec{j}$  and  $\vec{k}$  being the unit vectors of the rectangular axes Oxyz, we consider the vectors:

$$\vec{r}_1 = \vec{i} + 3\vec{j} - 2\vec{k}, \quad \vec{r}_2 = 4\vec{i} - 2\vec{j} + 2\vec{k} \quad \text{and} \quad \vec{r}_3 = 3\vec{i} - \vec{j} + 2\vec{k}$$

1. Show these 3 vectors graphically.
2. Calculate their modulus
3. Calculate products  $\vec{r}_1 \cdot \vec{r}_2$  and  $\vec{r}_1 \wedge \vec{r}_2$ .

#### Exercise 2

1. Let the points be  $M_1 (+1,+1,+1)$ ,  $M_2 (+2,+2,+1)$  and  $M_3 (+2,+1,0)$ ; calculate the angle  $\widehat{M_1 M_2 M_3}$
2. Determine the equation of the plane (p) passing through the point M2 and perpendicular to the vector  $\vec{A} = 3\vec{i} - 2\vec{j} + \vec{k}$

#### Exercise 3

$\vec{i}$ ,  $\vec{j}$  et  $\vec{k}$  being the unit vectors in the orthonormal frame (Oxyz). Let two vectors  $\vec{A}$  and  $\vec{B}$  be defined by:

$$\vec{A} = -\vec{i} + \vec{j} - 2\vec{k} \quad \text{and} \quad \vec{B} = 2\vec{i} + 4\vec{j} - 5\vec{k}$$

- 1- Calculate  $(\vec{A} \cdot \vec{B})$  and deduce the angle  $\theta = (\vec{A}, \vec{B})$
- 2- give  $(\vec{A} \wedge \vec{B})$ , deduce the area of the parallelogram formed by the two vectors

#### Exercise 4

We give the three vectors  $\vec{V}_1(1, 1, 0)$ ,  $\vec{V}_2(0, 1, 0)$  and  $\vec{V}_3(0, 0, 2)$ .

1. Calculate norms  $\|\vec{V}_1\|$ ,  $\|\vec{V}_2\|$  and  $\|\vec{V}_3\|$ , deduce the unit vectors  $\vec{v}_1$ ,  $\vec{v}_2$  and  $\vec{v}_3$  respectively from  $\vec{V}_1$ ,  $\vec{V}_2$  and  $\vec{V}_3$ .
2. Calculate  $\cos(\widehat{\vec{v}_1, \vec{v}_2})$ , knowing that the corresponding angle is between 0 and  $\pi$ .
3. Calculate the mixed product  $\vec{v}_1 \cdot (\vec{v}_2 \wedge \vec{v}_3)$ . What does this product represent?



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### Exercise 5

Consider in space, referred to the direct orthonormal reference frame  $(O, \vec{i}, \vec{j}, \vec{k})$  the points  $A(2, 0, 0)$ ,  $B(2, -2, 0)$  and  $C(2, 3, -1)$ .

1. Calculate the vector product  $\overrightarrow{OA} \wedge \overrightarrow{OB}$
2. Calculate the area of triangle OAB.
3. Calculate the mixed product  $(\overrightarrow{OA}, \overrightarrow{OB}, \overrightarrow{OC})$ , Deduce the volume of the parallelepiped built on the vectors.
4. Between these products, what are the mixed products that can be calculated  $(\overrightarrow{OA} \wedge \overrightarrow{OB}) \cdot \overrightarrow{OC}$ ;  $\overrightarrow{OA} \cdot (\overrightarrow{OB} \wedge \overrightarrow{OC})$ ;  $(\overrightarrow{OA} \cdot \overrightarrow{OB}) \wedge \overrightarrow{OC}$ ;  $\overrightarrow{OA} \wedge (\overrightarrow{OB} \cdot \overrightarrow{OC})$ ;

### Exercise 6

A) Let  $\vec{A}(1, 2, 1)$ ,  $\vec{B}(1, 0, c)$  be two vectors where  $c \in \mathbb{R}$

1. Calculate the scalar product  $\vec{A} \cdot \vec{B}$  and the modulus of the two vectors as a function of  $c$ .
2. Determine the values of  $c$  for which the angle  $(\vec{A}, \vec{B})$  is equal to  $\pi/3$ .

B) Consider the points  $A(3, 5, 4)$ ,  $B(3, 1, 3)$ ,  $C(8, 5, 5)$  and  $D(1, 2, 3)$  in space.

Calculate the mixed product  $(\overrightarrow{DA}, \overrightarrow{DB}, \overrightarrow{DC})$ , deduce the volume of the parallelepiped formed by the three vectors.

### Exercise 7

Let be three vectors  $\vec{A}$ ,  $\vec{B}$  and  $\vec{C}$ , such as;  $\vec{A} = -2\vec{i} + \vec{j} + 3\vec{k}$ ,  $\vec{B} = 2\vec{i} - \vec{j} + \vec{k}$ ,  $\vec{C} = x\vec{i} + 1\vec{j} + z\vec{k}$

1- Calculate  $x$  and  $z$  so that the vector  $\vec{C}$  or :

- a- Parallel to  $\vec{A}$       b- Parallel to  $\vec{B}$

2- If,  $\vec{C} = x\vec{i} + y\vec{j} + z\vec{k}$

Calculate  $x$ ,  $y$  and  $z$  so that the vector  $\vec{C}$  or : Perpendicular to  $\vec{A}$  and  $\vec{B}$  at the same time.

### Exercise 8

Let be a vector  $\vec{U} = (t\vec{i} + 3\vec{j}) / (\sqrt{t^2 + 9})$

1. Show that  $\vec{U}$  is a unit vector?

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2. Calculate its derivative with respect to time?

### Exercise 9

Let be the points A (+1,+1,+1), B (+2,+2,+1) and C (+2,+1,0)

1. Calculate the scalar product  $\overrightarrow{AB} \cdot \overrightarrow{AC}$  and the vector product  $\overrightarrow{AB} \wedge \overrightarrow{AC}$ .
2. What do these two products represent? Deduce the angle between the vectors  $\overrightarrow{AB}$  et  $\overrightarrow{AC}$ .

### Exercise 10

$\vec{i}$ ,  $\vec{j}$  and  $\vec{k}$  being the unit vectors of an orthonormal reference frame (Oxyz), consider the vectors.  $\vec{r}_1 = 2\vec{i} - 2\vec{j} + 3\vec{k}$  ,  $\vec{r}_2 = \vec{i} + \vec{j} + \vec{k}$

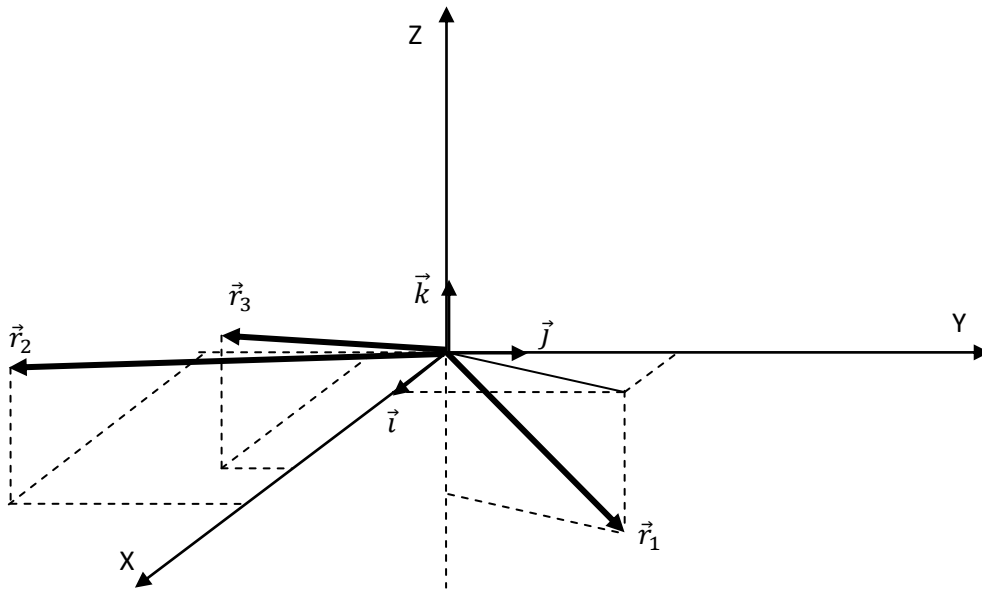
- 1- Calculate the vector product  $\vec{r}_1 \wedge \vec{r}_2$ .
- 2- Deduce the angle  $\theta$  formed by the two vectors  $\vec{r}_1$  and  $\vec{r}_2$ .

**Correction of exercises about chapter II**

**Exercise 1**

We are  $\vec{r}_1 = \vec{i} + 3\vec{j} - 2\vec{k}$   $\vec{r}_2 = 4\vec{i} - 2\vec{j} + 2\vec{k}$  and  $\vec{r}_3 = 3\vec{i} - \vec{j} + 2\vec{k}$

1- Vector representation  $\vec{r}_1, \vec{r}_2$  and  $\vec{r}_3$  :



2- The magnitudes of :

$$\vec{A} \begin{pmatrix} x \\ y \\ z \end{pmatrix} \Rightarrow |\vec{A}| = \|\vec{A}\| = \sqrt{x^2 + y^2 + z^2}$$

$$|\vec{r}_1| = \sqrt{x_1^2 + y_1^2 + z_1^2} = \sqrt{1 + 9 + 4} = \sqrt{14}$$

$$|\vec{r}_2| = \sqrt{x_2^2 + y_2^2 + z_2^2} = \sqrt{16 + 4 + 4} = \sqrt{24}$$

$$|\vec{r}_3| = \sqrt{x_3^2 + y_3^2 + z_3^2} = \sqrt{9 + 1 + 4} = \sqrt{14}$$

3-  $\vec{r}_1 \cdot \vec{r}_2 = x_1x_2 + y_1y_2 + z_1z_2 = 4 - 6 - 4 = -6$

$$\text{and } \vec{r}_1 \wedge \vec{r}_2 = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ 1 & 3 & -2 \\ 4 & -2 & 2 \end{vmatrix} = \begin{vmatrix} 3 & -2 \\ -2 & 2 \end{vmatrix} \vec{i} - \begin{vmatrix} 1 & -2 \\ 4 & 2 \end{vmatrix} \vec{j} + \begin{vmatrix} 1 & 3 \\ 4 & -2 \end{vmatrix} \vec{k}$$

$$\Rightarrow \vec{r}_1 \wedge \vec{r}_2 = (3 \cdot 2 - (-2) \cdot (-2))\vec{i} - (1 \cdot 2 - ((-2) \cdot 4))\vec{j} + (1 \cdot (-2) - (3 \cdot 4))\vec{k}$$

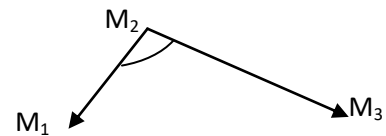
$$\Rightarrow \vec{r}_1 \wedge \vec{r}_2 = 2\vec{i} - 10\vec{j} - 14\vec{k}$$

**Exercise 2**

1- Let  $M_1(1,1,1)$ ,  $M_2(2,2,1)$  et  $M_3(2,1,0)$ .

Calculate the angle  $\widehat{M_1M_2M_3}$  which is the angle between the two vectors  $\overrightarrow{M_2M_1}$ ;  $\overrightarrow{M_2M_3}$ :

$$\text{We have } \overrightarrow{M_2M_1} = \begin{pmatrix} 1-2 \\ 1-2 \\ 1-1 \end{pmatrix} = \begin{pmatrix} -1 \\ -1 \\ 0 \end{pmatrix} = -\vec{i} - \vec{j}$$



$$\text{and } \overrightarrow{M_2M_3} = \begin{pmatrix} 2-2 \\ 1-2 \\ 0-1 \end{pmatrix} = \begin{pmatrix} 0 \\ -1 \\ -1 \end{pmatrix} = -\vec{j} - \vec{k}$$

$$\Rightarrow \begin{cases} \overrightarrow{M_2M_1} \cdot \overrightarrow{M_2M_3} = ((-1) \cdot 0) + (-1 \cdot (-1)) + (0 \cdot (-1)) = 1 \\ \overrightarrow{M_2M_1} \cdot \overrightarrow{M_2M_3} = |\overrightarrow{M_2M_1}| \cdot |\overrightarrow{M_2M_3}| \cdot \cos\theta = \sqrt{2} \cdot \sqrt{2} \cdot \cos\theta = 2\cos\theta \end{cases}$$

$$\Rightarrow 2\cos\theta = 1 \Rightarrow \cos\theta = \frac{1}{2} \Rightarrow \theta = \pm \frac{\pi}{3} + 2k\pi$$

$$\text{So } \widehat{M_1M_2M_3} = \pm \frac{\pi}{3} + 2k\pi$$

2- The equation of the plane passing through  $M_2(2,2,1)$  and perpendicular to the vector

$$\vec{A} = 3\vec{i} - 2\vec{j} + \vec{k}.$$

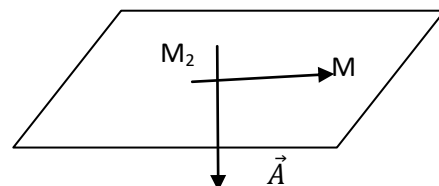
Let X be a point with coordinates (x,y,z) which belongs to the plane (p).

$$\overrightarrow{M_2X} = \begin{pmatrix} x-2 \\ y-2 \\ z-1 \end{pmatrix}$$

We know that  $\vec{A}$  perpendicular to this plane therefore:

$$\vec{A} \cdot \overrightarrow{M_2X} = (x-2)3 - 2(y-2) + (z-1)1 = 0$$

$$\Rightarrow 3x - 2y + z - 3 = 0 (*)$$



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(\*)s the equation of the plane passing through  $M_2(2,2,1)$  and perpendicular to the vector

$$\vec{A} = 3\vec{i} - 2\vec{j} + \vec{k}$$

### Exercise 3

A- Let the two vectors  $\vec{A} \begin{pmatrix} -1 \\ 1 \\ -2 \end{pmatrix}$  and  $\vec{B} \begin{pmatrix} 2 \\ 4 \\ -5 \end{pmatrix}$

1- The scalar product

$$\vec{A} \cdot \vec{B} = -2 + 4 + 10 = 12$$

According to the second writing of the scalar product  $\vec{A} \cdot \vec{B} = \|\vec{A}\| \|\vec{B}\| \cos(\vec{A}, \vec{B})$

$$\|\vec{A}\| = \sqrt{6}, \|\vec{B}\| = \sqrt{45} \text{ so } \cos(\vec{A}, \vec{B}) = \frac{\vec{A} \cdot \vec{B}}{\|\vec{A}\| \|\vec{B}\|} = \frac{12}{\sqrt{6} \cdot \sqrt{45}} = 0.730$$

So the angle  $\theta(\vec{A}, \vec{B}) = 43.08^\circ$

2- Vector product

$$\begin{aligned} \vec{A} \wedge \vec{B} &= \begin{vmatrix} \vec{i} & -\vec{j} & \vec{k} \\ -1 & 1 & -2 \\ 2 & 4 & -5 \end{vmatrix} = (-5 - (-8))\vec{i} - (5 - (-4))\vec{j} + (-4 - 2)\vec{k} \\ &= 3\vec{i} - 9\vec{j} - 6\vec{k} \end{aligned}$$

The vector product  $\vec{A} \wedge \vec{B}$  gives a vector perpendicular to the plane formed by the two vectors  $\vec{A}$  and  $\vec{B}$  and the module of this product ( $\|\vec{A} \wedge \vec{B}\|$ ) presents the surface of the parallelogram formed by the two vectors  $\vec{A}$  and  $\vec{B}$ .

$$\|\vec{A} \wedge \vec{B}\| = \sqrt{9 + 81 + 36} = \sqrt{126}$$

The area of the parallelogram is  $\sqrt{126}$ .

### Exercise 4

We give the three vectors  $\vec{V}_1(1, 1, 0)$ ,  $\vec{V}_2(0, 1, 0)$  and  $\vec{V}_3(0, 0, 2)$ .

1. Calculates the norms  $\|\vec{V}_1\|$ ,  $\|\vec{V}_2\|$  and  $\|\vec{V}_3\|$  :

Let's calculate the norms of the various vectors and the unit vectors of their respective directions.

$$\|\vec{V}_1\| = \sqrt{1^2 + 1^2 + 0^2} = \sqrt{2} \Rightarrow \vec{v}_1 = \frac{\vec{V}_1}{\|\vec{V}_1\|}; \vec{v}_1\left(\frac{\sqrt{2}}{2}, \frac{\sqrt{2}}{2}, 0\right)$$

$$\|\vec{V}_2\| = \sqrt{0^2 + 1^2 + 0^2} = 1 \Rightarrow \vec{v}_2 = \frac{\vec{V}_2}{\|\vec{V}_2\|} = \vec{j}; \vec{v}_2(0, 1, 0)$$

$$\|\vec{V}_3\| = \sqrt{0^2 + 0^2 + 2^2} = \sqrt{4} = 2 \Rightarrow \vec{v}_3 = \frac{\vec{V}_3}{\|\vec{V}_3\|}; \vec{v}_3(0, 0, 1)$$

2. Let's calculate  $\cos(\widehat{(\vec{v}_1, \vec{v}_2)})$  as follows :

$$\begin{aligned} \vec{v}_1 \cdot \vec{v}_2 &= \frac{\sqrt{2}}{2} \quad \text{and} \quad \vec{v}_1 \cdot \vec{v}_2 = \|\vec{v}_1\| \cdot \|\vec{v}_2\| \cdot \cos(\widehat{(\vec{v}_1, \vec{v}_2)}) \\ &\Rightarrow \cos(\widehat{(\vec{v}_1, \vec{v}_2)}) = \frac{\sqrt{2}}{2} \end{aligned}$$

3. We have a :

$$\vec{v}_1 \cdot \vec{v}_2 = \frac{\sqrt{2}}{2}$$

$$\vec{v}_2 \wedge \vec{v}_3 = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{vmatrix} = \begin{vmatrix} 1 & 0 \\ 0 & 1 \end{vmatrix} \vec{i} + 0\vec{j} + 0\vec{k} = \vec{i}$$

$$\Rightarrow \vec{v}_2 \wedge \vec{v}_3 = \vec{i}(1, 0, 0)$$

$$\vec{v}_1 \cdot (\vec{v}_2 \wedge \vec{v}_3) = 1 \times 1 + 1 \times 0 + 0 \times 0 = 1$$

- The first term represents the scalar product between the vectors  $\vec{v}_1$  and  $\vec{v}_2$  is equal to the product of the projection modulus of  $\vec{v}_1$  on  $\vec{v}_2$  multiplied by the magnitude of  $\vec{v}_2$ .
- The second term is the vector product between  $\vec{v}_2$  and  $\vec{v}_3$ .
- The last term is the mixed product between  $(\vec{v}_1, \vec{v}_2, \vec{v}_3)$  and is none other than the volume of the parallelepiped built on the basis of the three vectors.

### Exercise 5

A(2, 0, 0), B(2, -2, 0) and C(2, 3, -1).

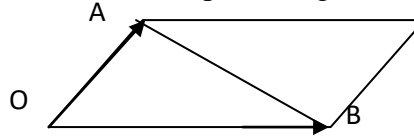
1. The vector product  $\overrightarrow{OA} \wedge \overrightarrow{OB}$  :

$$\overrightarrow{OA} \begin{pmatrix} 2 \\ 0 \\ 0 \end{pmatrix}; \overrightarrow{OB} \begin{pmatrix} 2 \\ -2 \\ 0 \end{pmatrix} \text{ so } \overrightarrow{OA} \wedge \overrightarrow{OB} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ 2 & 0 & 0 \\ 2 & -2 & 0 \end{vmatrix} = -4\vec{k}$$

## Chapter II: Vector Analysis

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The area of the triangle (OAB) is half the area of the parallelogram formed by the two vectors  $\vec{OA}$  and  $\vec{OB}$ .



$$S(OAB) = \frac{|\vec{OA} \wedge \vec{OB}|}{2} = \frac{4}{2} = 2$$

2. The mixed product  $(\vec{OA}, \vec{OB}, \vec{OC})$ , and the volume of the parallelepiped built on the vectors.

$$(\vec{OA} \wedge \vec{OB}) \cdot \vec{OC} = \begin{pmatrix} 0 \\ 0 \\ -4 \end{pmatrix} \cdot \begin{pmatrix} 2 \\ 3 \\ -1 \end{pmatrix} = 4$$

So the volume of the parallelepiped built on the vectors equal 4.

The products that can be calculated are

$$(\vec{OA} \wedge \vec{OB}) \cdot \vec{OC} = \vec{OA} \cdot (\vec{OB} \wedge \vec{OC});$$

On the other hand, these two products are false because the vector product can only be between two vectors

$$(\vec{OA} \cdot \vec{OB}) \wedge \vec{OC}; \quad \vec{OA} \wedge (\vec{OB} \cdot \vec{OC});$$

### Exercise 6

A) Let  $\vec{A}(1, 2, 1)$ ,  $\vec{B}(1, 0, c)$  be two vectors where  $c \in \mathbb{R}$

1. The scalar product  $\vec{A} \cdot \vec{B}$  and the modulus of the two vectors as a function of  $c$ .

$$\vec{A} \cdot \vec{B} = 1 + 0 + c = c + 1$$

$$\|\vec{A}\| = \sqrt{1^2 + 2^2 + 1^2} = \sqrt{6}$$

$$\|\vec{B}\| = \sqrt{1^2 + 0^2 + c^2} = \sqrt{1 + c^2}$$

2. The values of  $c$  for which the angle  $(\vec{A}, \vec{B})$  is equal to  $\pi/3$ .

According to the second writing of the scalar product  $\vec{A} \cdot \vec{B} = \|\vec{A}\| \|\vec{B}\| \cos(\vec{A}, \vec{B})$

$$(\vec{A}, \vec{B}) = \pi/3 \Rightarrow \cos(\vec{A}, \vec{B}) = 1/2$$

$$\cos(\vec{A}, \vec{B}) = \frac{\vec{A} \cdot \vec{B}}{\|\vec{A}\| \|\vec{B}\|} = \frac{c+1}{\sqrt{6} \cdot \sqrt{1+c^2}} = \frac{1}{2}$$

## Chapter II: Vector Analysis

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Then

$$c = 2 \pm \sqrt{3}$$

B) Consider the points A (3, 5, 4), B (3, 1, 3), C (8, 5, 5) and D (1, 2, 3) in space.

Calculate the mixed product  $(\overrightarrow{DA}, \overrightarrow{DB}, \overrightarrow{DC})$ ,

$$\overrightarrow{DA} = \begin{pmatrix} 3-1 \\ 5-2 \\ 4-3 \end{pmatrix} = \begin{pmatrix} 2 \\ 3 \\ 1 \end{pmatrix}, \overrightarrow{DB} = \begin{pmatrix} 3-1 \\ 1-2 \\ 3-3 \end{pmatrix} = \begin{pmatrix} 2 \\ -1 \\ 0 \end{pmatrix}, \overrightarrow{DC} = \begin{pmatrix} 8-1 \\ 5-2 \\ 5-3 \end{pmatrix} = \begin{pmatrix} 7 \\ 3 \\ 2 \end{pmatrix}$$

$$\overrightarrow{DA} \wedge \overrightarrow{DB} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ 2 & 3 & 1 \\ 2 & -1 & 0 \end{vmatrix} = \vec{i} + 2\vec{j} - 8\vec{k}$$

$$(\overrightarrow{DA} \wedge \overrightarrow{DB}) \cdot \overrightarrow{DC} = \begin{pmatrix} 1 \\ 2 \\ -8 \end{pmatrix} \cdot \begin{pmatrix} 7 \\ 3 \\ 2 \end{pmatrix} = -3$$

The volume of the parallelepiped formed by the three vectors is  $3 \text{ m}^3$  (We take the absolute value).

### Exercise 7

Let there be three vectors  $\vec{A}$ ,  $\vec{B}$  and  $\vec{C}$  such that

$$\vec{A} = -2\vec{i} + \vec{j} + 3\vec{k}, \quad \vec{B} = 2\vec{i} - \vec{j} + \vec{k}, \quad \vec{C} = x\vec{i} + 1\vec{j} + z\vec{k}$$

1- Calculate x,y and z so that the vector  $\vec{C}$  is:

a-  $\vec{C}$  Parallel to  $\vec{A}$  if  $\vec{A} \wedge \vec{C} = \vec{0}$

$$\vec{A} \wedge \vec{C} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ -2 & 1 & 3 \\ x & 1 & z \end{vmatrix} = \vec{0}$$

$$\Rightarrow (z - 3)\vec{i} - (-2z - 3x)\vec{j} + (-2 - x)\vec{k} = 0\vec{i} + 0\vec{j} + 0\vec{k}$$

$$\Rightarrow \begin{cases} z - 3 = 0 \\ 2z + 3x = 0 \\ -2 - x = 0 \end{cases}$$

$$\Rightarrow \begin{cases} z = 3 \\ 2z + 3x = 0 \\ x = -2 \end{cases}$$

$$\Rightarrow \vec{C} = -2\vec{i} + \vec{j} + 3\vec{k}$$



b-  $\vec{C}$  Parallel to  $\vec{B}$  if  $\vec{B} \wedge \vec{C} = \vec{0}$

$$\vec{B} \wedge \vec{C} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ 2 & -1 & 1 \\ x & 1 & z \end{vmatrix} = \vec{0}$$

$$\Rightarrow (-z - 1)\vec{i} - (2z - x)\vec{j} + (2 + x)\vec{k} = 0\vec{i} + 0\vec{j} + 0\vec{k}$$

$$\Rightarrow \begin{cases} -z - 1 = 0 \\ 2z - x = 0 \\ 2 + x = 0 \end{cases}$$

$$\Rightarrow \vec{C} = -2\vec{i} + \vec{j} - 1\vec{k}$$

2- If,  $\vec{C} = x\vec{i} + y\vec{j} + z\vec{k}$

$\vec{C}$  Perpendicular to  $\vec{A}$  and  $\vec{B}$  at the same time if  $\vec{C}' = \vec{A} \wedge \vec{B}$

$$\vec{A} \wedge \vec{B} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ -2 & 1 & 3 \\ 2 & -1 & 1 \end{vmatrix}$$

$$\Rightarrow (1 + 3)\vec{i} - (-2 - 6)\vec{j} + (2 - 2)\vec{k}$$

$$\Rightarrow \vec{C}' = 4\vec{i} + 8\vec{j}$$

### Exercise 8

Let be a vector  $\vec{U} = (t\vec{i} + 3\vec{j})/(\sqrt{t^2 + 9})$

1-  $\vec{U}$  is a unit vector ?

Check that  $|\vec{U}| = 1$  or  $|\vec{U}| = \sqrt{\frac{1}{(t^2+9)}(t^2 + 9)} = 1$

So  $\vec{U}$  is an unit vector.

2- The derivative of  $\vec{U}$  :

$$\frac{d\vec{u}}{dt} = \frac{d}{dt} \left( \frac{t}{(\sqrt{t^2 + 9})} \right) \vec{i} + \frac{d}{dt} \left( \frac{3}{(\sqrt{t^2 + 9})} \right) \vec{j}$$

$$\Rightarrow \frac{d\vec{u}}{dt} = \left( \frac{t^2 - t^2 + 9}{(t^2 + 9)^{3/2}} \right) \vec{i} + \left( \frac{-3t}{(t^2 + 9)^{3/2}} \right) \vec{j}$$

$$\Rightarrow \frac{d\vec{u}}{dt} = \left( \frac{9}{(t^2 + 9)^{3/2}} \right) \vec{i} + \left( \frac{-3t}{(t^2 + 9)^{3/2}} \right) \vec{j}$$

### Exercise 9

A. let be the points A (+1,+1,+1), B (+2,+2,+1) and C (+2,+1,0)

The Scalar product  $\overrightarrow{AB} \cdot \overrightarrow{BC}$

$$\overrightarrow{AB} \begin{pmatrix} 2-1 \\ 2-1 \\ 1-1 \end{pmatrix} = \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \text{ and } \overrightarrow{BC} \begin{pmatrix} 2-2 \\ 1-2 \\ 0-1 \end{pmatrix} = \begin{pmatrix} 0 \\ -1 \\ -1 \end{pmatrix}$$

$$\overrightarrow{AB} \cdot \overrightarrow{BC} = 1 \times 0 + 1 \times (-1) + 0 \times (-1) = -1 = |\overrightarrow{AB}| \cdot |\overrightarrow{BC}| \cos \theta$$

The vector product  $\overrightarrow{AB} \wedge \overrightarrow{AC}$  with  $\overrightarrow{AC} \begin{pmatrix} 2-1 \\ 1-1 \\ 0-1 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix}$

$$\overrightarrow{AB} \wedge \overrightarrow{AC} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ 1 & 1 & 0 \\ 1 & 0 & -1 \end{vmatrix} = -\vec{i} - (-1)\vec{j} - \vec{k} = -\vec{i} + \vec{j} - \vec{k}$$

The angle between  $\overrightarrow{AB} \cdot \overrightarrow{BC}$  will;  $\cos \theta = \frac{\overrightarrow{AB} \cdot \overrightarrow{BC}}{|\overrightarrow{AB}| \cdot |\overrightarrow{BC}|} = \frac{-1}{2}$  so  $\theta = -\pi/6$

### Exercise 10

We have:  $\vec{r}_1 = 2\vec{i} - 2\vec{j} + 3\vec{k}$  and  $\vec{r}_2 = \vec{i} + \vec{j} + \vec{k}$

1- Calculation of vector product  $\vec{r}_1 \wedge \vec{r}_2$ .

$$\vec{r}_1 \wedge \vec{r}_2 = -5\vec{i} + \vec{j} + 4\vec{k}$$

2- The modulus of vector product  $\vec{r}_1 \wedge \vec{r}_2$  is:

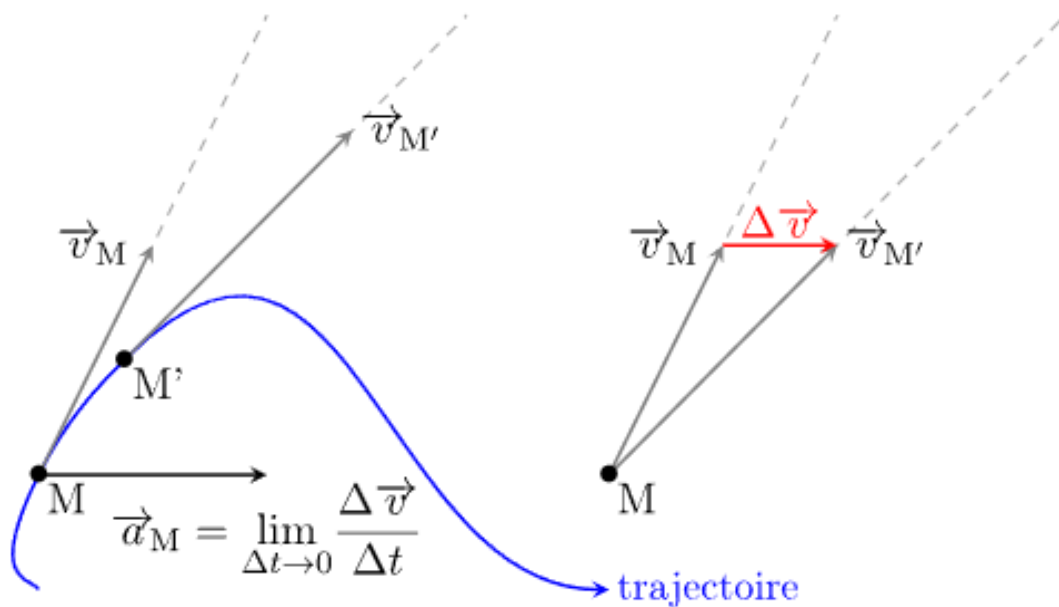
$$|\vec{r}_1 \wedge \vec{r}_2| = |\vec{r}_1| \cdot |\vec{r}_2| \cdot \sin \theta$$

$$\Rightarrow \sin \theta = \frac{|\vec{r}_1 \wedge \vec{r}_2|}{|\vec{r}_1| \cdot |\vec{r}_2|} = 0.9$$

$$\Rightarrow \theta \sim 64^\circ$$

## COURSE OF MECHANICS OF THE MATERIAL POINT

### *Chapter III: Kinematics of material point*



**Glossary**

In English	In French	In Arabic
Kinematics	La cinématique	الحركيات
Material point	Un point matériel	نقطة مادية
Reference system	Un système référentiel	نظام مرجعي
Velocity (speed)	La vitesse	السرعة
Acceleration	L'accélération	التسارع
Motion characteristics	Caractéristiques d'un mouvement	خصائص الحركة
Position vector	Vecteur position	شعاع الموضعي
Time equation/ Hourly equation	Equation horaire	المعادلة الزمنية
Trajectory	Trajectoire	المسار
Trajectory equation	Equation de la trajectoire	معادلة المسار
Velocity vector	Vecteur vitesse	شعاع السرعة
Acceleration vector	Vecteur accélération	شعاع التسارع
Coordinate systems	Système de coordonnées	نظام الإحداثيات
Cartesian coordinates	Coordonnées cartésiennes	الإحداثيات الكارتيزية
Polar coordinates	Coordonnées polaire	الإحداثيات القطبية
Cylindrical coordinates	Coordonnées cylindriques	الإحداثيات الأسطوانية
Spherical coordinates	Coordonnées sphériques	الإحداثيات الكروية
Intrinsic coordinates	Coordonnées intrinsèques	الإحداثيات الجوهرية
Rectilinear movement	Movement rétiligne	حركة مستقيمة
Uniform rectilinear movement	Movement rétiligne uniforme MRU	حركة مستقيمة منتظمة
Uniformly varied rectilinear movement	Movement rétiligne uniformément varié MRUV	حركة مستقيمة متغيرة بانتظام
Circular movement	Movement circulaire MC	حركة دائرية
Uniform circular movement	Movement circulaire uniforme MCU	حركة دائرية منتظمة
Uniformly varied circular movement	Movement circulaire uniformément varié MCUV	حركة دائرية متغيرة بانتظام

## Chapter III: Kinematics of material point

Sinusoidal or harmonic movement	Mouvement sinusoïdal ou harmonique	حركة جيبية أو توافقية
A frame	Un referential	معلم او مرجع
The equation of motion	Equation de mouvement	معادلة الحركة
A mobile	Un mobile	متحرك
Average velocity	La vitesse moyenne	السرعة المتوسطة
Instantaneous velocity	La vitesse instantanée	السرعة اللحظية
Average acceleration	L'accélération moyenne	التسارع المتوسط
Instantaneous acceleration	L'accélération instantanée	التسارع اللحظي
The orthonormal coordinate system	Un système de coordonnées orthogonal	نظام الإحداثيات المتعامد
The Frenet frame	Le repère de Frenet / trièdre de Frenet.	معلم فرينال
The moving point	Un point en mouvement	نقطة مادية في حالة حركة
The normal acceleration	L'accélération normal	التسارع الناطمي
tangential acceleration	L'accélération tangentielle	التسارع المماسي
Motion	Le Mouvement	الحركة
Weight	Le poids	الوزن
Linear velocity	La vitesse linéaire	السرعة الخطية
Angular velocity	La vitesse angulaire	السرعة الزاوية
Linear Acceleration	L'accélération linéaire	التسارع الخطي
Angular Acceleration	L'accélération angulaire	التسارع الزاوي
Acceleration of gravity	Accélération de pesanteur	تسارع الجاذبية
Height	La hauteur	الارتفاع
The period of a pendulum	La période d'une pendule simple	دور نواس بسيط
The sound	Le son	الصوت
Radius	Le rayon	نصف قطر
The abscissa	L'abscisse	الفاصلة
Radius of curvature	Le rayon de courbure	نصف قطر المسار المنحني
The right triangle	Un triangle droit	مثلث قائم
Amplitude	Amplitude	السعة
Frequency	Fréquence	التواتر
Average speed	La vitesse moyenne	السرعة المتوسطة
Instantaneous speed	La vitesse instantanée	السرعة اللحظية

### 1. Introduction

The theory of General Relativity invented by A. Einstein in 1915 is a relativistic theory of gravitation. This theory challenges the idea of an inert Euclidean space, independent of its material content. Kinematics studies the movement of a material point independently of the causes that give rise to it. It is based on a Euclidean description of space and absolute time. The material point is any material body whose dimensions are theoretically zero and practically negligible in relation to the distance it travels. The state of movement or rest of a body is two essentially relative notions: for example, a mountain is at rest in relation to the earth, but in movement in relation to an observer looking at the earth from afar, for whom the globe (with all that it contains) is in perpetual movement. In this course, we illustrate the notions of velocity and acceleration by restricting ourselves to movements in the plane.

### 2. Reference System مرجع

The concept of motion is relative. A body can be in motion with respect to one object and at rest with respect to another (relative motion), hence the necessity of choosing a reference frame. A reference frame is a system of coordinate axes linked to an observer.

This study of motion is carried out in two forms:

- **Vectorial:** using vectors: position  $\overrightarrow{OM}$ , velocity  $\vec{v}$ , and acceleration  $\vec{a}$ .
- **Algebraic:** by defining the equation of motion along a given trajectory.

### 3. Characteristics of a movement

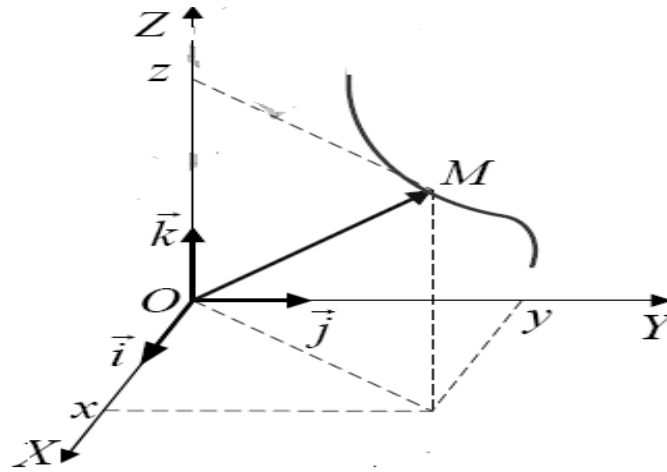
#### 3.1. Vector position and time equation شجاع الموضعي و المعادلة الزمنية للحركة

We define the position of a material point M in a reference frame by the position vector  $\overrightarrow{OM}$ , where O is a fixed point and serves as the origin of the reference frame. The components of point M or the vector  $\overrightarrow{OM}$  are given in the chosen coordinate system's basis (Cartesian coordinates, polar coordinates, etc.).

The point M moves through time, and this movement is described by an equation known as the "time equation" (معادلة زمنية), translated as the "time equation."

#### 3.2. Trajectory المسار

The trajectory is the geometric path of successive positions occupied by the material point over time with respect to the considered reference system.



**Example:**

The position of a material point M identified by its coordinates (x, y, z) at time t in a coordinate system  $R(O, \vec{i}, \vec{j}, \vec{k})$  with a position vector:

$$\vec{OM} = (t - 1)\vec{i} + \frac{t^2}{2}\vec{j}$$

$$\vec{OM} = (t - 1)\vec{i} + \frac{t^2}{2}\vec{j} \Rightarrow \begin{cases} x = t - 1 \\ y = \frac{t^2}{2} \end{cases}$$

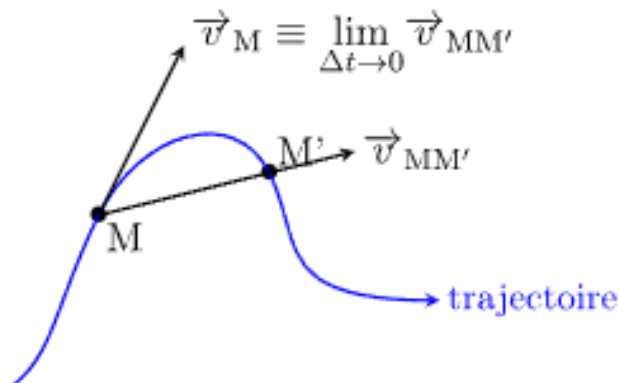
So  $t = x + 1$

The trajectory equation of the material point is

$$y = \frac{(x + 1)^2}{2}$$

**3.3. Velocity vector** شعاع السرعة

Consider a mobile that is located at position M(t) at time t, and it evolves at the point M'(t+Δt) at instant(t+Δt).



## Chapter III: Kinematics of material point

- The **average velocity** السرعة المتوسطة between the two instants  $t$  and  $t+\Delta t$  is called:

$$\overrightarrow{v_{moy}} = \frac{\overrightarrow{MM'}}{(t + \Delta t) - t} = \frac{\overrightarrow{MM'}}{\Delta t}$$

- If the time interval  $\Delta t$  is very small ( $\Delta t \rightarrow 0$ ), we then refer to it as **instantaneous velocity** السرعة اللحظية:

$$\vec{v} = \lim_{\Delta t \rightarrow 0} \overrightarrow{v_{moy}} = \lim_{\Delta t \rightarrow 0} \frac{\overrightarrow{MM'}}{\Delta t}$$

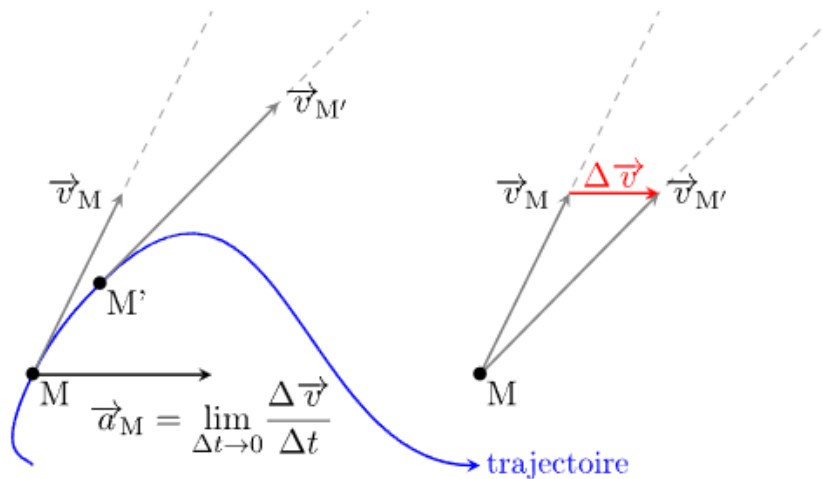
$$\overrightarrow{MM'} = \overrightarrow{MO} + \overrightarrow{OM'} = \overrightarrow{OM'} - \overrightarrow{OM} = \Delta \overrightarrow{OM}$$

So:

$$\vec{v} = \lim_{\Delta t \rightarrow 0} \frac{\Delta \overrightarrow{OM}}{\Delta t} \Rightarrow \vec{v} = \frac{d\overrightarrow{OM}}{dt}$$

### 3.4. Acceleration vector شعاع التسارع

When velocity varies over time  $v=f(t)$ , point  $M$  is subjected to an acceleration.



- The average acceleration التسارع المتوسط is written:

$$\overrightarrow{a_{moy}} = \frac{\vec{v}(t + \Delta t) - \vec{v}(t)}{(t + \Delta t) - t} = \frac{\Delta \vec{v}(t)}{\Delta t}$$

- When the time is very small  $\Delta t \rightarrow 0$  instantaneous acceleration التسارع اللحظي is written by :

$$\vec{a} = \lim_{\Delta t \rightarrow 0} \frac{\Delta \overrightarrow{OM}}{\Delta t}$$

$$\Rightarrow \vec{a} = \frac{d\vec{v}(t)}{dt} = \frac{d^2 \overrightarrow{OM}}{dt^2}$$



## 4. Expression of velocity and acceleration in different coordinate systems

To solve a problem in physics, we must locate the position of the moving point M in space  $\overrightarrow{OM}(t)$ .

The position must be located from a frame of reference (reference), we are required to choose the appropriate reference to use it according to the problem we want to solve

Generally, we use Cartesian, Polar, Cylindrical or Spherical coordinates

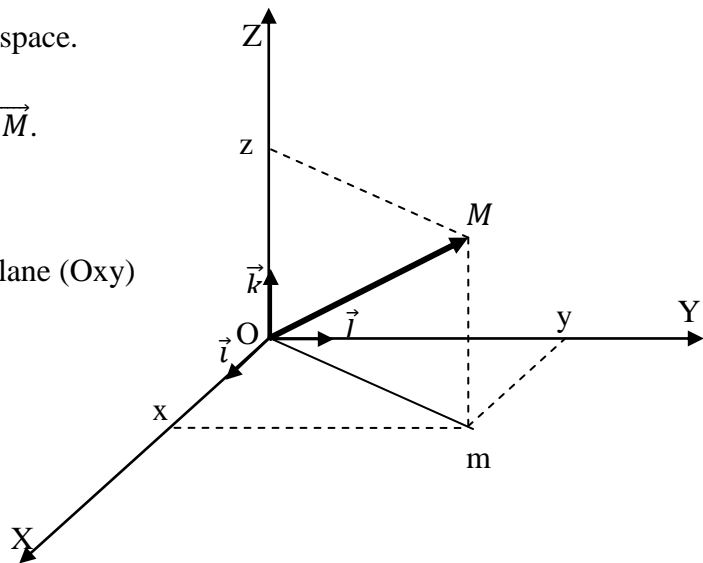
### 4.1 Cartesian coordinates

Let the frame be  $R(O,x,y,z)$  with the unit vectors  $\vec{i}, \vec{j}$  and  $\vec{k}$ . With  $x, y$  and  $z$  are the coordinates of point M which gives its position in space.

They are also the vector components  $\overrightarrow{OM}$ .

$x$ : abscissa;  $y$ : ordinate and  $z$ : height

$m$  is the projection of point M in the plane (Oxy)



- **Vecteur position**

$$\overrightarrow{OM} = x\vec{i} + y\vec{j} + z\vec{k}$$

The unit vectors  $\vec{i}, \vec{j}$  et  $\vec{k}$  constitute a basis linked to the axes (Ox), (Oy) and (Oz)

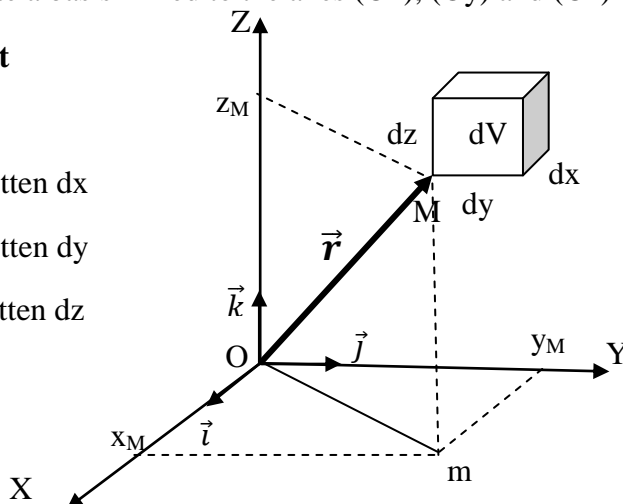
- **Elementary displacement**

The elementary displacement  $d\vec{l}$ :

Next (Ox) the displacement is written  $dx$

Next (Oy) the displacement is written  $dy$

Next (Oz) the displacement is written  $dz$



By fixing  $y$  and  $z$ , M moves along  $\vec{i}$ , the elementary displacement is then written  $d\vec{l}_1 = dx\vec{i}$ .

By fixing  $x$  and  $z$ , M moves along  $\vec{j}$ , the elementary movement is then written  $d\vec{l}_2 = dy\vec{j}$ .

By fixing  $x$  and  $y$ , M moves along  $\vec{k}$ , the elementary displacement is then written  $d\vec{l}_3 = dz\vec{k}$ .

## Chapter III: Kinematics of material point

The total displacement of point M is:

$$\vec{dl} = \vec{dl}_1 + \vec{dl}_2 + \vec{dl}_3 = dx\vec{i} + dy\vec{j} + dz\vec{k}$$

Or mathematically :

$$\vec{OM} = x\vec{i} + y\vec{j} + z\vec{k} \Rightarrow d\vec{OM} = dx\vec{i} + dy\vec{j} + dz\vec{k}$$

The elementary volume  $dV=dl_1.dl_2.dl_3=dx.dy.dz$

- **Velocity vector**

$$\vec{v} = \frac{d\vec{OM}}{dt} = \frac{dx}{dt}\vec{i} + \frac{dy}{dt}\vec{j} + \frac{dz}{dt}\vec{k} \Rightarrow \begin{cases} v_x = \frac{dx}{dt} \\ v_y = \frac{dy}{dt} \\ v_z = \frac{dz}{dt} \end{cases}$$

The velocity module is written:  $|\vec{v}| = \sqrt{v_x^2 + v_y^2 + v_z^2}$

Note: The magnitude of the velocity, equal to  $|\vec{v}|$ , is called the speed. In S.I. units, v is expressed in (m/s) or ( $m.s^{-1}$ ).

- **Acceleration vector:**

$$\vec{a} = \frac{d\vec{v}}{dt} = \frac{d^2\vec{OM}}{dt^2} \Rightarrow \begin{cases} a_x = \frac{dv_x}{dt} = \frac{d^2x}{dt^2} \\ a_y = \frac{dv_y}{dt} = \frac{d^2y}{dt^2} \\ a_z = \frac{dv_z}{dt} = \frac{d^2z}{dt^2} \end{cases}$$

The acceleration module is written:

$$|\vec{a}| = \sqrt{a_x^2 + a_y^2 + a_z^2}$$

The unit of acceleration in S.I units is ( $m/s^2$ ) or ( $m.s^{-2}$ ).

### 4.2. Polar coordinates الاحداثيات القطبية

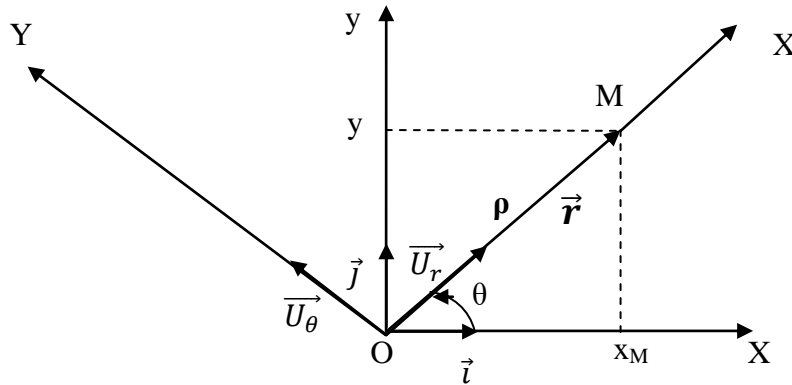
When the motion is in a plane, it's also possible to locate the position of point M using its polar coordinates ( $\rho, \theta$ ).

$\rho$ : polar radius  $\rho = |\vec{OM}|$  ( $0 \leq \rho \leq R$ )

$\theta$ : polar angle  $\theta = (\text{ox}, \vec{OM})$  ( $0 \leq \theta \leq 2\pi$ )

## Chapter III: Kinematics of material point

Let's consider point M moving in space, identified by its polar coordinates  $(\rho, \theta)$  in the orthonormal coordinate system (OXY) with unit vectors  $\vec{u}_r, \vec{u}_\theta$ .



- **Position Vector**

The position vector of a material point M in polar coordinates is written:  $\vec{r} = \overrightarrow{OM} = \rho \vec{U}_r$

The unit vectors  $\vec{U}_r$  is following  $\overrightarrow{OM}$  and  $\vec{U}_\theta$  is perpendicular to  $\vec{U}_r$  ( $\vec{U}_r \perp \vec{U}_\theta$ ).

### Transit relations between cartesian coordinates and polar coordinates

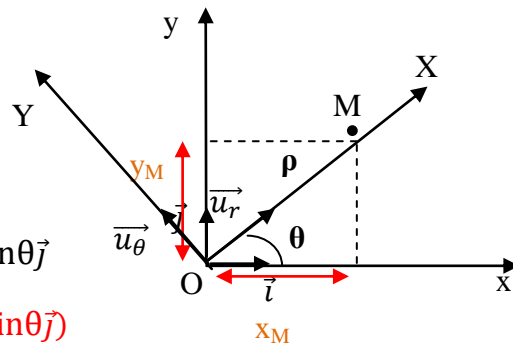
We project the point M into the plane (Oxy)

$$\begin{cases} x_M = |\overrightarrow{OM}| \cos\theta = \rho \cos\theta \\ y_M = |\overrightarrow{OM}| \sin\theta = \rho \sin\theta \end{cases}$$

$$\overrightarrow{OM} = x_M \vec{i} + y_M \vec{j} \Rightarrow \overrightarrow{OM}/cart = \rho \cos\theta \vec{i} + \rho \sin\theta \vec{j}$$

$$\overrightarrow{OM}/pol = \rho \vec{U}_r \text{ and } \overrightarrow{OM}/cart = \rho (\cos\theta \vec{i} + \sin\theta \vec{j})$$

By identification:  $\vec{u}_r = \cos\theta \vec{i} + \sin\theta \vec{j}$



**Rule: Note: The derivative of a unit vector with respect to an angle is a unit vector perpendicular to the angle in the positive direction.**

*مشتقة شعاع وحدة بالنسبة إلى الزاوية هي شعاع وحدة عمودي على هذا الأخير في الاتجاه الموجب*

The vector  $\vec{u}_\theta \perp \vec{u}_r$  in the direction of  $\theta$  which corresponds to the direct direction therefore

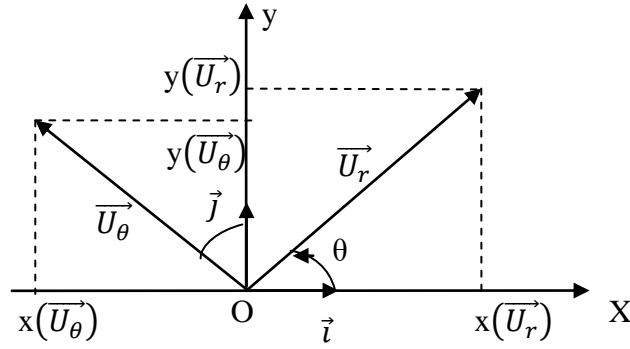
$$\vec{u}_\theta = \frac{d\vec{u}_r}{d\theta}$$

So  $\vec{u}_\theta = -\sin\theta \vec{i} + \cos\theta \vec{j}$

## Chapter III: Kinematics of material point

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By projecting the unit vectors we will have the same results



$$\vec{u}_r = x(\vec{u}_r)\vec{i} + y(\vec{u}_r)\vec{j} \Rightarrow \vec{u}_r = |\vec{u}_r|\cos\theta\vec{i} + |\vec{u}_r|\sin\theta\vec{j}$$

with  $|\vec{u}_r| = 1$  since it is a unit vector therefore  $\vec{u}_r = \cos\theta\vec{i} + \sin\theta\vec{j}$

$$\vec{u}_\theta = x(\vec{u}_\theta)\vec{i} + y(\vec{u}_\theta)\vec{j} \Rightarrow \vec{u}_\theta = -|\vec{u}_\theta|\sin\theta\vec{i} + |\vec{u}_\theta|\cos\theta\vec{j}$$

with  $|\vec{u}_\theta| = 1$  since it is a unit vector therefore  $\vec{u}_\theta = -\sin\theta\vec{i} + \cos\theta\vec{j}$

To write the unit vectors  $\vec{u}_r$  and  $\vec{u}_\theta$  as a function of  $\vec{i}$  and  $\vec{j}$  we use the passage table

	$\vec{i}$	$\vec{j}$
$\vec{u}_r$	$\cos\theta$	$\sin\theta$
$\vec{u}_\theta$	$-\sin\theta$	$\cos\theta$

$$\vec{i} = \cos\theta\vec{u}_r - \sin\theta\vec{u}_\theta \quad \text{and} \quad \vec{j} = \sin\theta\vec{u}_r + \cos\theta\vec{u}_\theta$$

### **Example :**

Write the vector  $\vec{A} = 2x\vec{i} + y\vec{j}$

$$\begin{cases} x = \rho \cos\theta \\ y = \rho \sin\theta \end{cases} \quad \text{and} \quad \begin{cases} \vec{i} = \cos\theta\vec{u}_r - \sin\theta\vec{u}_\theta \\ \vec{j} = \sin\theta\vec{u}_r + \cos\theta\vec{u}_\theta \end{cases}$$

$$\text{So } \vec{A} = 2\rho \cos\theta (\cos\theta\vec{u}_r - \sin\theta\vec{u}_\theta) + \rho \sin\theta (\sin\theta\vec{u}_r + \cos\theta\vec{u}_\theta)$$

$$\Rightarrow \vec{A} = 2\rho \cos^2\theta\vec{u}_r - 2\rho(\cos\theta\sin\theta)\vec{u}_\theta + \rho \sin^2\theta \vec{u}_r + \rho(\sin\theta \cos\theta)\vec{u}_\theta$$

$$\Rightarrow \vec{A} = \rho (2\cos^2\theta + \sin^2\theta)\vec{u}_r - \rho(\cos\theta\sin\theta)\vec{u}_\theta$$

$$\Rightarrow \vec{A} = (\rho \cos^2\theta + 1)\vec{u}_r - \rho(\cos\theta\sin\theta)\vec{u}_\theta$$

## Chapter III: Kinematics of material point

- **Elementary displacement**

The variables  $\rho$  and  $\theta$  are independent: we fix one and change the other

- We fix  $\theta$ , and we change  $\rho$ , then the moving point moves from the point  $M(\rho, \theta)$  to the point  $M'(\rho+d\rho, \theta)$

$$\overrightarrow{dl_1} = \overrightarrow{MM'} = d\rho \overrightarrow{u_r}$$

- We fix  $\rho$ , and we change  $\theta$ , then the moving point moves from the point  $M(\rho, \theta)$  to the point  $M'(\rho, \theta+d\theta)$

The angle  $\theta$  varies by  $d\theta$ , causing a linear displacement of point  $M$  towards point  $M''$   
 $\overrightarrow{U_\theta}, (\overrightarrow{MM''} \perp \overrightarrow{u_\theta})$

In the right triangle  $OMM''$ ,  $MM'' = \rho \sin\theta d\theta$ .

Since  $d\theta$  is very small, we can approximate

$\sin(d\theta)$  as  $d\theta$ .

Therefore,  $MM'' = \rho d\theta$ , so

$$\overrightarrow{dl_2} = \overrightarrow{MM''} = \rho d\theta \overrightarrow{u_\theta}$$

so

$$\overrightarrow{dl} = \overrightarrow{dl_1} + \overrightarrow{dl_2} = d\rho \overrightarrow{u_r} + \rho d\theta \overrightarrow{u_\theta}$$

We can obtain the same result mathematically:

$$\overrightarrow{OM} = \rho \overrightarrow{U_r} \Rightarrow d\overrightarrow{OM} = d\rho \overrightarrow{U_r} + \rho d\overrightarrow{U_r}$$

To make the derivative of a unit vector  $d\overrightarrow{U_r}$ , we must bring out the derivative with respect to an angle  $\frac{d\overrightarrow{U_r}}{d\theta}$  for this we multiply and divide by  $d\theta$

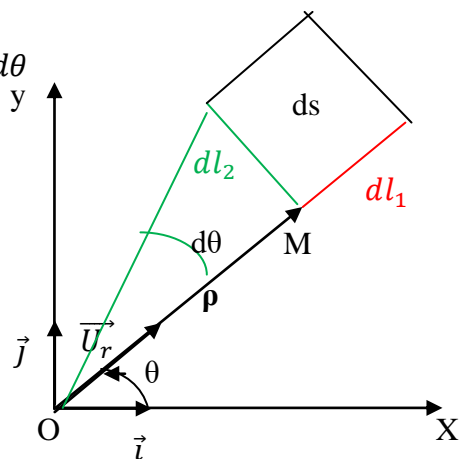
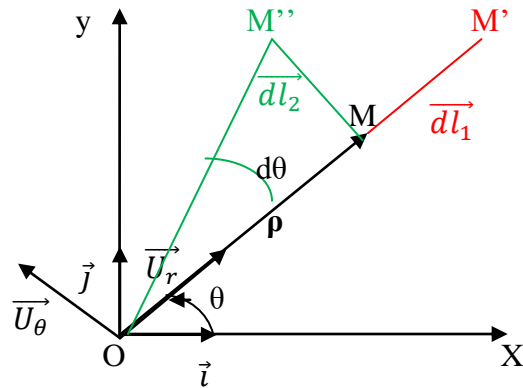
$$d\overrightarrow{U_r} = \frac{d\overrightarrow{U_r}}{d\theta} d\theta = \overrightarrow{U_\theta} d\theta$$

With  $\frac{d\overrightarrow{U_r}}{d\theta} = \overrightarrow{U_\theta}$  so  $d\overrightarrow{OM} = d\rho \overrightarrow{U_r} + \rho d\theta \overrightarrow{U_\theta}$

**Calculation of the surface:**

$$ds = |\overrightarrow{dl_1}| \cdot |\overrightarrow{dl_2}| = d\rho \cdot \rho d\theta \Rightarrow s = \iint d\rho \cdot \rho d\theta$$

We can separate the variables since they are independent



$$s = \int_0^R \rho d\rho \cdot \int_0^{2\pi} d\theta = \frac{R^2}{2} 2\pi$$

with  $\rho$  varies from 0 to R and  $\theta$  varies from 0 to  $2\pi \Rightarrow s = \pi R^2$

- **Velocity vector**

$$\vec{v} = \frac{d\overline{OM}}{dt} = \frac{d\rho}{dt} \vec{U}_r + \rho \frac{d\vec{U}_r}{dt}$$

We have:  $\frac{d\vec{U}_r}{dt} = \frac{d\vec{U}_r}{dt} \frac{d\theta}{d\theta} = \frac{d\vec{U}_r}{d\theta} \frac{d\theta}{dt}$

With :  $\frac{d\vec{U}_r}{d\theta} = \vec{U}_\theta$  donc  $\frac{d\vec{U}_r}{dt} = \frac{d\theta}{dt} \vec{U}_\theta$  so  $\vec{v} = \frac{d\overline{OM}}{dt} = \frac{d\rho}{dt} \vec{U}_r + \rho \frac{d\theta}{dt} \vec{U}_\theta$

$$\Rightarrow \vec{v} = \rho \cdot \dot{\vec{U}}_r + \rho \theta \cdot \vec{U}_\theta \text{ with } \rho \cdot = \frac{d\rho}{dt} \text{ and } \theta \cdot = \frac{d\theta}{dt}$$

- **Acceleration vector**

$$\vec{a} = \frac{d\vec{v}}{dt} = \frac{d^2\overline{OM}}{dt^2} = \frac{d^2\rho}{dt^2} \vec{U}_r + \frac{d\rho}{dt} \frac{d\vec{U}_r}{dt} + \frac{d\rho}{dt} \frac{d\theta}{dt} \vec{U}_\theta + \rho \frac{d^2\theta}{dt^2} \vec{U}_\theta + \rho \frac{d\theta}{dt} \frac{d\vec{U}_\theta}{dt}$$

$$\Rightarrow \vec{a} = \frac{d^2\rho}{dt^2} \vec{U}_r + \frac{d\rho}{dt} \frac{d\theta}{dt} \vec{U}_\theta + \frac{d\rho}{dt} \frac{d\theta}{dt} \vec{U}_\theta + \rho \frac{d^2\theta}{dt^2} \vec{U}_\theta - \rho \left( \frac{d\theta}{dt} \right)^2 \vec{U}_r$$

With :  $\frac{d\vec{U}_r}{d\theta} = \vec{U}_\theta$  et  $\frac{d\vec{U}_\theta}{d\theta} = -\vec{U}_r$

$$\text{So : } \vec{a} = \rho \cdot \ddot{\vec{U}}_r + 2\rho \cdot \dot{\theta} \cdot \ddot{\vec{U}}_\theta + \rho \theta \cdot \ddot{\vec{U}}_\theta - \rho (\dot{\theta})^2 \vec{U}_r$$

### 4.3. Cylindrical Coordinates الاحداثيات الاسطوانية

If the spatial trajectory involves  $\rho$  and  $z$  playing a specific role in determining the position vector ( $\overline{OM}$ ); for example, the movement of air molecules in a whirlwind; it is preferable to use cylindrical coordinates ( $\rho, \theta, z$ ). With:

$\rho$ : polar radius

$\theta$ : polar angle

$z$ : altitude or height

$$\text{and } \begin{cases} \rho = |\overline{Om}|, & 0 < \rho < R \\ \theta = ((ox), \overline{Om}), & 0 < \theta < 2\pi \\ z = z_M, & 0 < z < H \end{cases}$$

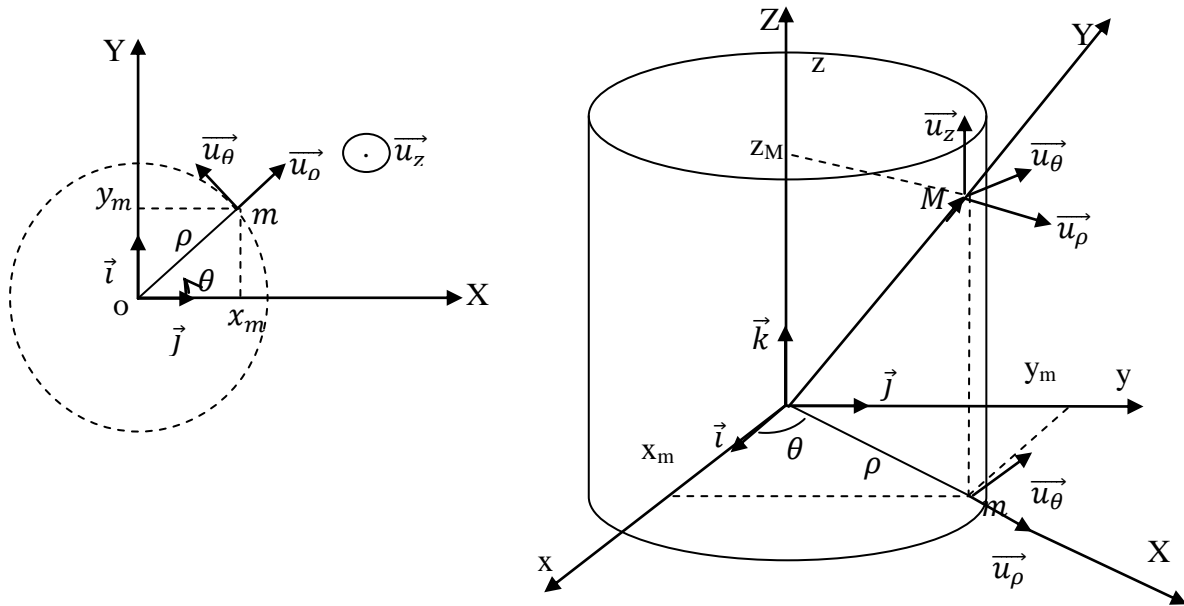
Where m is the projection of point M onto the plane (Oxy), and R is the radius of the cylinder, and H is the height of the cylinder.

## Chapter III: Kinematics of material point

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If we add the 'z' component to polar coordinates in space, we obtain what is known as cylindrical coordinates. Consider  $R(Oxyz)$  and a point  $M$  belonging to a cylinder.

Point  $M$  is identified by three coordinates  $\rho$ ,  $\theta$  (polar coordinates), and  $z$ .



- **Position vector**

The position vector in cylindrical coordinates  $(\rho, \theta, z)$  in the orthonormal frame  $R(O, \vec{u}_\rho, \vec{u}_\theta, \vec{u}_z)$  is written:

$$\begin{cases} \vec{r} = \overline{OM} = \overline{Om} + \overline{mM} \text{ (Relation de Charles)} \\ \overline{Om} = \rho \vec{u}_\rho \text{ (Coordonnées polaires)} \\ \overline{mM} = z \vec{u}_z \text{ (hauteur du cylindre)} \end{cases} \Rightarrow \vec{r} = \overline{OM} = \rho \vec{u}_\rho + z \vec{u}_z$$

- **Unit vectors**

The unit vectors  $\vec{U}_\rho$  is following  $\overline{Om}$  ( $m$  is the projection of the point  $M$  on the plane  $(Oxy)$ ) and  $\vec{U}_\theta$  is perpendicular to  $\vec{U}_r$  and  $\overline{Om}$  in the direction of  $\theta$  ( $\vec{U}_\rho \perp \vec{U}_\theta$ ) and  $\vec{U}_z$  is following  $(Oz)$ , ( $\vec{U}_z \parallel \vec{k}$ ) and it is perpendicular to the plane formed by the two other unit vectors ( $\vec{U}_\rho$  and  $\vec{U}_\theta$ ).

### Transit relations between cylindrical coordinates and Cartesian coordinates:

By projecting the point  $m$  onto the axes  $(Ox)$  and  $(Oy)$  (like polar coordinates)  $z$  is the height

$$\begin{cases} x = \rho \cos \theta \\ y = \rho \sin \theta \\ z = z \end{cases}$$

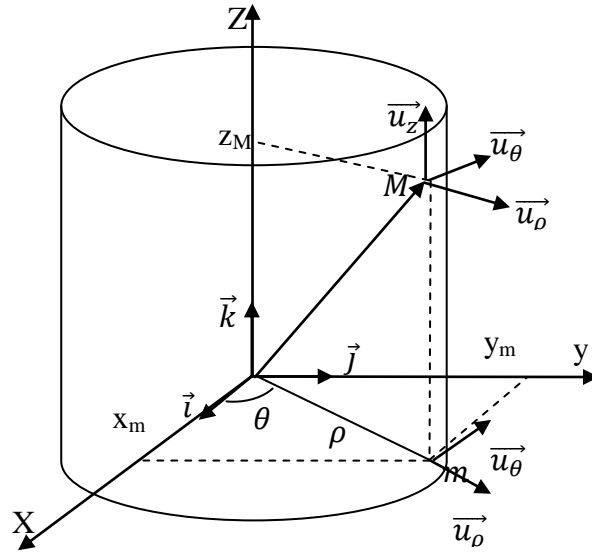
$$\overrightarrow{OM}/_{cylin} = \overrightarrow{Om} + \overrightarrow{mM} = \rho \overrightarrow{u}_\rho + z \overrightarrow{u}_z$$

$$\overrightarrow{OM}/_{cart} = x \vec{i} + y \vec{j} + z \vec{k}$$

$$\overrightarrow{OM}/_{cart} = \rho (\cos \theta \vec{i} + \sin \theta \vec{j}) + z \vec{k}$$

**By identification**

$$\begin{cases} \overrightarrow{u}_\rho = \cos \theta \vec{i} + \sin \theta \vec{j} \\ \overrightarrow{u}_\theta = \frac{d\overrightarrow{u}_\rho}{d\theta} = -\sin \theta \vec{i} + \cos \theta \vec{j} \\ \overrightarrow{u}_z = \vec{k} \end{cases}$$



Using the passage table :

	$\vec{i}$	$\vec{j}$	$\vec{k}$
$\overrightarrow{u}_\rho$	$\cos \theta$	$\sin \theta$	0
$\overrightarrow{u}_\theta$	$-\sin \theta$	$\cos \theta$	0
$\overrightarrow{u}_z$	0	0	1

$$\vec{i} = \cos \theta \overrightarrow{u}_\rho - \sin \theta \overrightarrow{u}_\theta$$

$$\vec{j} = \sin \theta \overrightarrow{u}_\rho + \cos \theta \overrightarrow{u}_\theta \quad \text{and} \quad \vec{k} = \overrightarrow{u}_z$$

- Elementary displacement**

The variables  $\rho$ ,  $\theta$  and  $z$  are independent: we fix one and change the other

- We fix  $\theta$ ,  $z$  and we change  $\rho$ , then the moving point moves from the point  $M(\rho, \theta, z)$  to the point  $M'(\rho+d\rho, \theta, z)$

$$d\vec{l}_1 = \overrightarrow{MM'} = d\rho \overrightarrow{u}_\rho$$

- We fix  $\rho$ ,  $z$  and we change  $\theta$ , then the moving point moves from the point  $M(\rho, \theta, z)$  to the point  $M'(\rho, \theta+d\theta, z)$

The angle  $\theta$  varies by  $d\theta$ , this leads to a linear movement from point  $M$  towards point  $M''$

following  $\overrightarrow{U}_\theta$ , ( $\overrightarrow{MM''} \parallel \overrightarrow{u}_\theta$ )

In the right triangle  $OMM''$ ,  $MM'' = \rho \sin d\theta$



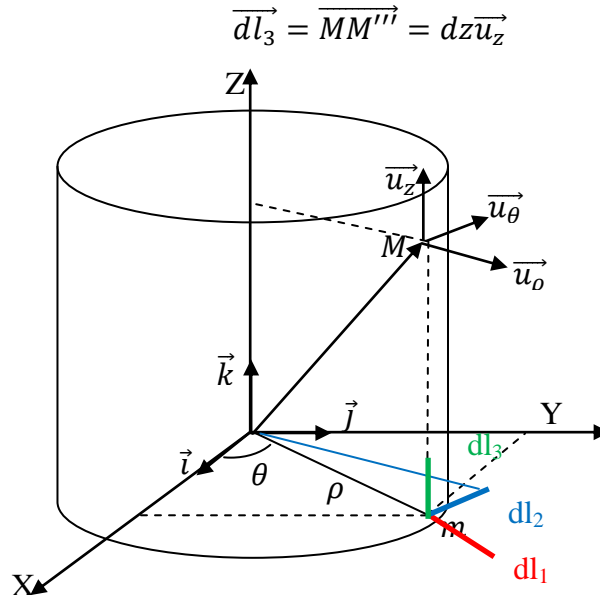
## Chapter III: Kinematics of material point

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$d\theta$  is so small then  $\sin d\theta \approx d\theta$ .

Then  $MM'' = \rho d\theta$  therefore  $\overrightarrow{dl_2} = \overrightarrow{MM''} = \rho d\theta \overrightarrow{u_\theta}$

- We fix  $\rho, \theta$  and we change  $z$ , then the moving point moves from the point  $M(\rho, \theta, z)$  to the point  $M'''(\rho, \theta, z+dz)$



So

$$\overrightarrow{dl} = \overrightarrow{dl_1} + \overrightarrow{dl_2} + \overrightarrow{dl_3} = d\rho \overrightarrow{u_\rho} + \rho d\theta \overrightarrow{u_\theta} + dz \overrightarrow{u_z}$$

We can obtain the same result mathematically

$$\overrightarrow{OM} = \rho \overrightarrow{u_\rho} + z \overrightarrow{u_z} \Rightarrow d\overrightarrow{OM} = d\rho \overrightarrow{U_\rho} + \rho d\overrightarrow{U_\rho} + dz \overrightarrow{u_z} + z d\overrightarrow{u_z}$$

$d\overrightarrow{u_z} = 0$  car  $\overrightarrow{u_z} = \vec{k}$  it's a vector fix.

$$d\overrightarrow{U_\rho} = \frac{d\overrightarrow{U_\rho}}{d\theta} d\theta = \overrightarrow{U_\theta} d\theta$$

With  $\frac{d\overrightarrow{U_\rho}}{d\theta} = \overrightarrow{U_\theta}$  so  $d\overrightarrow{OM} = d\rho \overrightarrow{U_\rho} + \rho d\theta \overrightarrow{U_\theta} + dz \overrightarrow{u_z}$

- **The cylinder volume**

$$dV = |d\overrightarrow{l_1}| \cdot |d\overrightarrow{l_2}| \cdot |d\overrightarrow{l_3}| = d\rho \cdot \rho d\theta \cdot dz \Rightarrow V = \iiint d\rho \cdot \rho d\theta \cdot dz$$

We can separate the variables since they are independent

$$V = \int_0^R \rho d\rho \cdot \int_0^{2\pi} d\theta \int_0^H dz = \frac{R^2}{2} 2\pi H \Rightarrow V = \pi R^2 H$$

(with  $\rho$  varies from 0 to R and  $\theta$  varies from 0 to  $2\pi$  and  $z$  varies from 0 to H)

- **The surface of the base of cylinder**

$$ds_{base} = |\overrightarrow{dl_1}| \cdot |\overrightarrow{dl_2}| = d\rho \cdot \rho d\theta \Rightarrow s_{base} = \iint d\rho \cdot \rho d\theta$$

We can separate the variables since they are independent

$$s = \int_0^R \rho d\rho \cdot \int_0^{2\pi} d\theta = \frac{R^2}{2} 2\pi \Rightarrow s_{base} = \pi R^2$$

- **The lateral surface of cylinder**

$$ds_{lat} = |\overrightarrow{dl_2}| \cdot |\overrightarrow{dl_3}| = d\rho \cdot \rho d\theta \Rightarrow s_{base} = \iint \rho d\theta \cdot dz$$

The radius is constant  $\rho=R$ , the variables are independent so we can separate them

$$s = R \int_0^{2\pi} d\theta \cdot \int_0^H dz = R 2\pi H \Rightarrow s_{base} = 2\pi R H$$

- **Velocity vector**

The velocity in this case is written by:  $\vec{v} = \frac{d\overrightarrow{OM}}{dt} = \frac{d\rho}{dt} \vec{U}_r + \rho \frac{d\vec{U}_r}{dt} + \frac{dz}{dt} \vec{U}_z + z \frac{d\vec{U}_z}{dt}$

$$\frac{d\vec{U}_r}{dt} = \frac{d\vec{U}_r}{dt} \frac{d\theta}{d\theta} = \frac{d\vec{U}_r}{d\theta} \frac{d\theta}{dt}$$

With:  $\frac{d\vec{U}_r}{d\theta} = \vec{U}_\theta$  so  $\frac{d\vec{U}_r}{dt} = \frac{d\theta}{dt} \vec{U}_\theta$  and  $\frac{d\vec{U}_z}{dt} = \vec{0}$

$$\vec{v} = \frac{d\overrightarrow{OM}}{dt} = \frac{d\rho}{dt} \vec{U}_r + \rho \frac{d\theta}{dt} \vec{U}_\theta + \frac{dz}{dt} \vec{U}_z$$

$$\Rightarrow \vec{v} = \rho \cdot \vec{U}_r + \rho \theta \cdot \vec{U}_\theta + z \cdot \vec{U}_z$$

With:  $\rho \cdot = \frac{d\rho}{dt}$ ,  $\theta \cdot = \frac{d\theta}{dt}$  and  $z \cdot = \frac{dz}{dt}$

- **Acceleration vector**

$$\vec{a} = \frac{d\vec{v}}{dt} = \frac{d^2\overrightarrow{OM}}{dt^2} = \frac{d^2\rho}{dt^2} \vec{U}_r + \frac{d\rho}{dt} \frac{d\vec{U}_r}{dt} + \frac{d\rho}{dt} \frac{d\theta}{dt} \vec{U}_\theta + \rho \frac{d^2\theta}{dt^2} \vec{U}_\theta + \rho \frac{d\theta}{dt} \frac{d\vec{U}_\theta}{dt} + \frac{d^2z}{dt^2} \vec{U}_z + \frac{dz}{dt} \frac{d\vec{U}_z}{dt}$$

$$\Rightarrow \vec{a} = \frac{d^2\rho}{dt^2} \vec{U}_r + \frac{d\rho}{dt} \frac{d\theta}{dt} \vec{U}_\theta + \frac{d\rho}{dt} \frac{d\theta}{dt} \vec{U}_\theta + \rho \frac{d^2\theta}{dt^2} \vec{U}_\theta - \rho \left( \frac{d\theta}{dt} \right)^2 \vec{U}_r + \frac{d^2z}{dt^2} \vec{U}_z$$

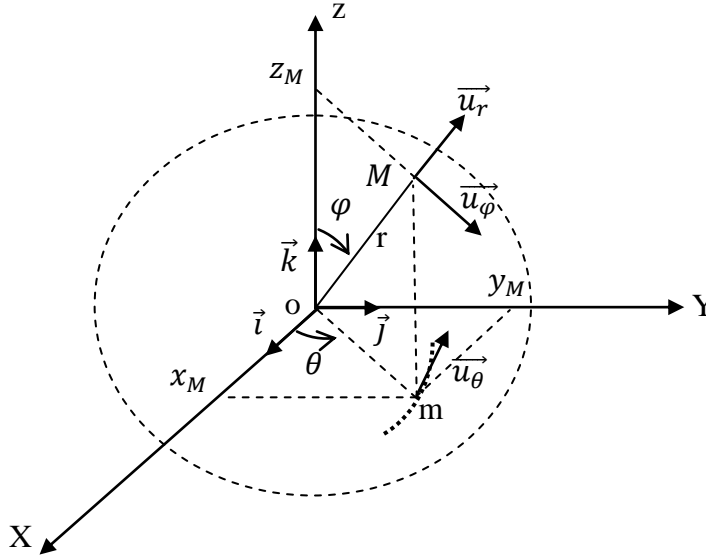
With :  $\frac{d\vec{U}_r}{d\theta} = \vec{U}_\theta$ ,  $\frac{d\vec{U}_\theta}{d\theta} = -\vec{U}_r$  and  $\frac{d\vec{U}_z}{dt} = \vec{0}$

$$\Rightarrow \vec{a} = \rho \ddot{U}_r + 2\dot{\rho}\dot{\theta}\vec{U}_\theta + \rho\ddot{\theta}\vec{U}_\theta - \rho(\dot{\theta})^2\vec{U}_r + z\ddot{U}_z$$

#### 4.4. Spherical coordinates *الإحداثيات الكروية*

When the point O and the distance r separating M and O play a characteristic role, the use of spherical coordinates (r,θ,φ) are the best suited in the orthonormed base ( $\vec{u}_r, \vec{u}_\theta, \vec{u}_\varphi$ ) with:

$$\begin{cases} r = |\overline{OM}|, & 0 < r < R \\ \theta = ((ox), \overline{Om}) & 0 < \theta < 2\pi \\ \varphi = ((oz), \overline{OM}) & 0 < \varphi < \pi \end{cases}$$



With m is the projection of M in the plane (Oxy).

- **Position Vector :**

The position vector in spherical coordinates (r,θ,φ) is written as:  $\vec{r} = \overline{OM} = r\vec{U}_r$

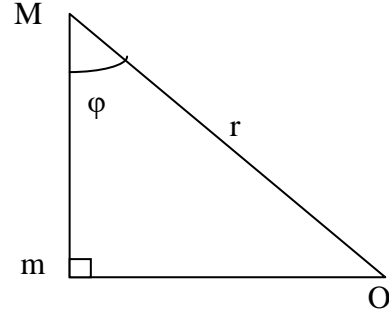
- **The unit vectors**

The unit vectors  $\vec{U}_r$  is following  $\overline{OM}$  and  $\vec{U}_\varphi$  is perpendicular to  $\vec{U}_r$  and  $\overline{OM}$  in the direction of φ ( $\vec{U}_\varphi \perp \vec{U}_r$ ) and  $\vec{U}_\theta$  is perpendicular to  $\overline{Om}$  ( $\vec{U}_\theta \perp \overline{Om}$ ).

#### Transit relations between spherical coordinates and Cartesian coordinates

By projecting m onto the axes (Ox) and (Oy)

$$\begin{cases} x = |\overrightarrow{Om}| \cos \theta \\ y = |\overrightarrow{Om}| \sin \theta \\ z = |\overrightarrow{mM}| \end{cases}$$



By taking the right triangle (OmM):

We have  $Om=r \sin \varphi$  and  $mM=r \cos \varphi$ , replacing them in passing relationships we will have:

$$\begin{cases} x = r \sin \varphi \cos \theta \\ y = r \sin \varphi \sin \theta \\ z = r \cos \varphi \end{cases}$$

$$\overrightarrow{OM} /_{sph} = r \vec{u}_r$$

$$\overrightarrow{OM} /_{cart} = r \sin \varphi \cos \theta \vec{i} + r \sin \varphi \sin \theta \vec{j} + r \cos \varphi \vec{k}$$

$$\overrightarrow{OM} /_{cart} = r (\sin \varphi \cos \theta \vec{i} + \sin \varphi \sin \theta \vec{j} + \cos \varphi \vec{k})$$

**By identification**

$$\vec{u}_r = \sin \varphi \cos \theta \vec{i} + \sin \varphi \sin \theta \vec{j} + \cos \varphi \vec{k}$$

$$\vec{u}_\varphi = \frac{-d\vec{u}_r}{d(-\varphi)} = \frac{d\vec{u}_r}{d\varphi} = \cos \varphi \cos \theta \vec{i} + \cos \varphi \sin \theta \vec{j} - \sin \varphi \vec{k}$$

$$\begin{aligned} \vec{U}_\theta &= \vec{U}_r \wedge \vec{U}_\varphi = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ \sin \varphi \cos \theta & \sin \varphi \sin \theta & \cos \varphi \\ \cos \varphi \cos \theta & \cos \varphi \sin \theta & -\sin \varphi \end{vmatrix} \\ &= \vec{i}(-\sin^2 \varphi \sin \theta - \cos^2 \varphi \sin \theta) - \vec{j}(-\sin^2 \varphi \cos \theta - \cos^2 \varphi \cos \theta) \\ &\quad + \vec{k}(\sin \varphi \cos \theta \cos \varphi \sin \theta - \sin \varphi \sin \theta \cos \varphi \cos \theta) \\ \vec{U}_\theta &= -\sin \theta \vec{i} + \cos \theta \vec{j} \end{aligned}$$

By using the passage table :

	$\vec{i}$	$\vec{j}$	$\vec{k}$
$\vec{u}_r$	$\sin \varphi \cos \theta$	$\sin \varphi \sin \theta$	$\cos \varphi$
$\vec{u}_\varphi$	$\cos \varphi \cos \theta$	$\cos \varphi \sin \theta$	$-\sin \varphi$
$\vec{u}_\theta$	$-\sin \theta$	$\cos \theta$	1

$$\begin{aligned} \vec{i} &= \sin \varphi \cos \theta \vec{u}_r + \cos \varphi \cos \theta \vec{u}_\varphi - \sin \theta \vec{u}_\theta \\ \vec{j} &= \sin \varphi \sin \theta \vec{u}_r + \cos \varphi \sin \theta \vec{u}_\varphi + \cos \theta \vec{u}_\theta \\ \vec{k} &= \cos \varphi \vec{u}_r - \sin \varphi \vec{u}_\varphi \end{aligned}$$

## Chapter III: Kinematics of material point

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- **Elementary displacement**

The variables  $r$ ,  $\varphi$  and  $\theta$  are independent: we fix one and change the other:

- We fix  $\varphi, \theta$  and we change  $r$ , then the moving point moves from the point  $M(r, \varphi, \theta)$  to the point  $M'(r+dr, \varphi, \theta)$  so  $\overrightarrow{dl_1} = \overrightarrow{MM'} = dr \overrightarrow{u_r}$
- We fix  $r, \theta$  and we change  $\varphi$ , then the moving point moves from the point  $M(r, \varphi, \theta)$  to the point  $M''(r, \varphi+d\varphi, \theta)$

The angle  $\varphi$  varies by  $d\varphi$ , this leads to a linear movement from point  $M$  towards point  $M''$  following  $\overrightarrow{U_\varphi}$ ,  $(\overrightarrow{MM''} \perp \overrightarrow{u_\varphi})$ .

In the right triangle  $OMM''$ ,  $MM'' = r \cdot \sin d\varphi$

$d\varphi$  is so small then  $\sin d\varphi \approx d\varphi$

then  $MM'' = r \cdot d\varphi$  therefore  $\overrightarrow{dl_2} = \overrightarrow{MM''} = r d\varphi \overrightarrow{u_\varphi}$

- We fix  $r, \varphi$  and we change  $\theta$ , then the moving point moves from the point  $M(r, \varphi, \theta)$  to the point  $M'''(r, \varphi, \theta+d\theta)$

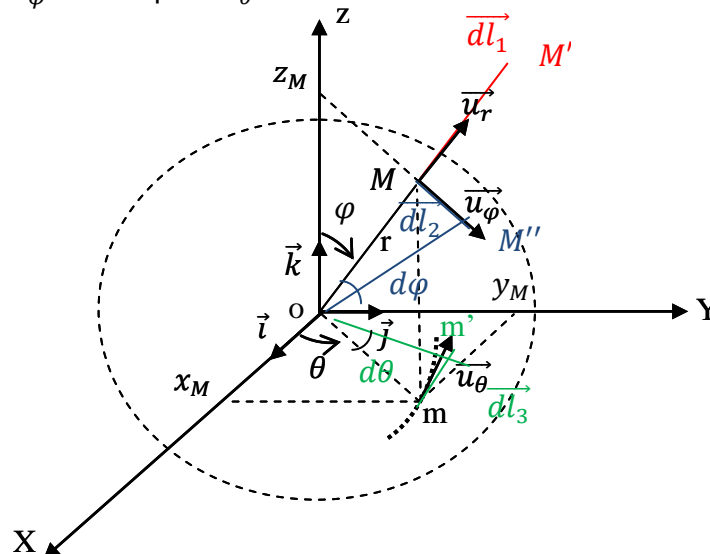
The angle  $\theta$  varies by  $d\theta$ , this leads to a linear displacement of the point  $m$  (projection of the point  $M$  in the plane  $(Oxy)$ ) towards the following point  $m'$   $\overrightarrow{U_\theta}$ ,  $(\overrightarrow{mm'} \perp \overrightarrow{u_\theta})$ .

In the right triangle  $Omm'$ ,  $mm' = Om \cdot \sin d\theta$

$d\theta$  is so small then  $\sin d\theta \approx d\theta$

So  $mm' = Om d\theta$  and  $Om = r \sin \varphi$  therefore  $\overrightarrow{dl_3} = \overrightarrow{mm'} = r \sin \varphi d\theta \overrightarrow{u_\theta}$

$\overrightarrow{dOM} = dr \overrightarrow{u_r} + r d\varphi \overrightarrow{u_\varphi} + r \sin \varphi d\theta \overrightarrow{u_\theta}$



Or mathematically :

$$\begin{aligned}\overline{OM} &= r\overline{U}_r \Rightarrow d\overline{OM} = d(r\overline{U}_r) = dr\overline{U}_r + r d\overline{U}_r \\ d\overline{U}_r &= \frac{\partial \overline{U}_r}{\partial \theta} d\theta + \frac{\partial \overline{U}_r}{\partial \varphi} d\varphi \\ \overline{U}_r &= \sin\varphi \cos\theta \vec{i} + \sin\varphi \sin\theta \vec{j} + \cos\varphi \vec{k} \\ \frac{\partial \overline{U}_r}{\partial \theta} &= -\sin\varphi \sin\theta \vec{i} + \sin\varphi \cos\theta \vec{j} = \sin\varphi (-\sin\theta \vec{i} + \cos\theta \vec{j}) \\ &\Rightarrow \frac{\partial \overline{U}_r}{\partial \theta} = \sin\varphi \overline{U}_\theta \\ \frac{\partial \overline{U}_r}{\partial \varphi} &= \cos\varphi \cos\theta \vec{i} + \cos\varphi \sin\theta \vec{j} - \sin\varphi \vec{k} \\ &\Rightarrow \frac{\partial \overline{U}_r}{\partial \varphi} = \overline{U}_\varphi \\ d\overline{OM} &= dr\overline{U}_r + r d\varphi \overline{U}_\varphi + r \sin\varphi d\theta \overline{U}_\theta\end{aligned}$$

- **Volume of sphere**

$$dV = dl_1 dl_2 dl_3 = dr \ r \sin\varphi \ d\theta \ r \ d\varphi \Rightarrow V = \iiint r^2 dr \sin\varphi \ d\varphi \ d\theta$$

$$\Rightarrow V = \int_0^R r^2 dr \int_0^\pi \sin\varphi \int_0^{2\pi} d\theta = \left. \frac{r^3}{3} \right] (-\cos\varphi) \Big|_0^\pi \theta$$

$$\Rightarrow V = \frac{4}{3} \pi R^3$$

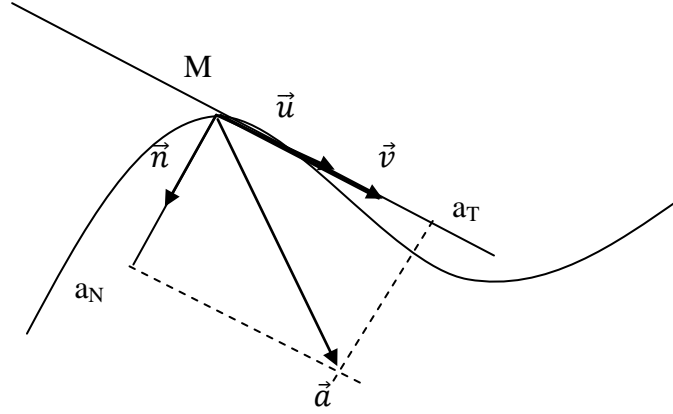
- **Velocity vector**

The velocity vector is written in spherical coordinates (r,θ,φ) by:

$$\vec{v} = \frac{d\overline{OM}}{dt} = \frac{dr}{dt} \overline{U}_r + r \frac{d\varphi}{dt} \overline{U}_\varphi + r \sin\varphi \frac{d\theta}{dt} \overline{U}_\theta$$

### 4.5. Intrinsic coordinates (Frenet frame) احداثيات الحركة المنحنية

We used to work in a fixed frame, but in this case, we study the motion in a moving frame that travels with the moving point "M". This frame is the Frenet frame.



We study the motion in the Frenet frame:

The Frenet frame is a two-dimensional reference frame.

- $\vec{u}$  is the unit vector along the tangent to the trajectory.
- $\vec{n}$  is the unit vector normal to the trajectory and perpendicular to  $\vec{u}$ , directed towards the center of curvature.
- The position remains unchanged (the frame moves with point M).
- The velocity vector is tangent to the trajectory, and it is written as:  $\vec{v} = |\vec{v}|\vec{u}$
- The acceleration vector :

$$\vec{a} = \frac{d\vec{v}}{dt} = \frac{d|\vec{v}|\vec{u}}{dt} = \frac{d|\vec{v}|}{dt}\vec{u} + |\vec{v}|\frac{d\vec{u}}{dt}$$

$$\frac{d\vec{u}}{dt} = \frac{d\vec{u}}{d\theta} \cdot \frac{d\theta}{dt} = \vec{n} \cdot \omega \text{ with } \vec{n} = \frac{d\vec{u}}{d\theta} \text{ and } \omega = \frac{d\theta}{dt}$$

The acceleration vector is written by:  $\vec{a} = a_T\vec{u} + a_N\vec{n}$

$$\text{So : } \vec{a} = \frac{d|\vec{v}|}{dt}\vec{u} + |\vec{v}| \cdot \vec{n} \cdot \omega$$

(the perimeter of a circle (محيط دائرة)  $l = 2\pi R$ , for the length of a segment (طول قوس))

$$x = \theta R \ ; \ \text{from angular velocity to linear velocity by } \frac{dx}{dt} = R \frac{d\theta}{dt} \Rightarrow v = R\omega$$

Hence:

$$\omega = \frac{v}{R} \text{ with } R \text{ is the radius of the curvature of the trajectory.}$$

$$\text{So } \vec{a} = \frac{d|\vec{v}|}{dt}\vec{u} + \frac{v^2}{R}\vec{n}$$

The normal acceleration (التسارع الناطمي) and tangential acceleration (التسارع المماس) are written

$$\text{by : } \begin{cases} a_T = \frac{d|\vec{v}|}{dt} \\ a_N = \frac{v^2}{R} \end{cases}$$

$$|\vec{a}| = \sqrt{a_x^2 + a_y^2} = \sqrt{a_N^2 + a_T^2}$$

### Note :

- $R \rightarrow \infty$  : so the trajectory is a line .
- $R$  is constant: so the trajectory is circular.

## 5. Study of some movements

### 5.1. Rectilinear motion حركة خطية

We have linear motion if the trajectory is a straight line.

We choose a point O as the origin on the trajectory and a unit vector  $\vec{i}$  .

The position of the mobile M, as a function of time, is identified by its abscissa:

$$x(t) = \overline{OM(t)}.$$

The position vector will be:  $\vec{r}(t) = \overline{OM(t)} = x(t)\vec{i}$

#### 5.1.1. Uniform rectilinear motion حركة مستقيمة منتظمة URM

We have uniform rectilinear motion if the trajectory is a straight line and the velocity vector is constant. This is a motion with zero acceleration  $\overline{a(t)} = \vec{0}$ .

The initial conditions to  $t=0$  ;  $x=x_0$ .

- *The velocity*

$$a = \frac{dv}{dt} = 0 \Rightarrow \int_{v_0}^v dv = \int_0^t 0 \cdot dt = cte$$

So  $v=v_0=cte$

- *The position*

$$v = \frac{dx}{dt} = v_0 \Rightarrow \int_{x_0}^x dx = \int_0^t v_0 dt = [v_0 t]_0^t = v_0 t$$

So :  $x=v_0 t+x_0$  *This is the hourly equation of the motion. URM*

#### 5.1.2. Uniformly varied rectilinear motion حركة مستقيمة متغيرة بانتظام UVRM

One has a uniformly varied rectilinear movement if the trajectory is a straight and the acceleration is constant.

The initial conditions to  $t=0$  ;  $v=v_0$  and  $x=x_0$



- **The velocity**

$$a = \frac{dv}{dt} = a_0 \Rightarrow \int_{v_0}^v dv = \int_0^t a_0 dt = [a_0 t]_0^t$$

So  $v = a_0 t + v_0$

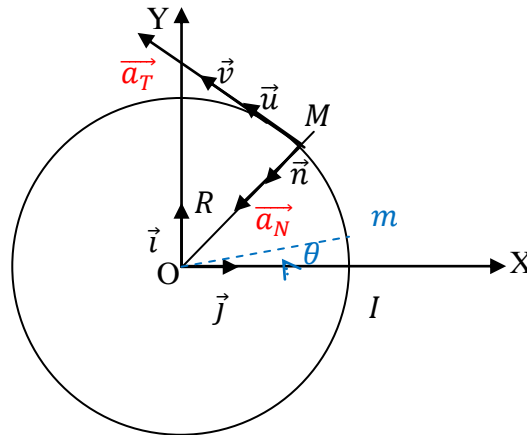
- **The position**

$$v = \frac{dx}{dt} = a_0 t + v_0 \Rightarrow \int_{x_0}^x dx = \int_0^t (a_0 t + v_0) dt = \left[ \frac{1}{2} a_0 t^2 + v_0 t \right]_0^t$$

So  $x = \frac{1}{2} a_0 t^2 + v_0 t + x_0$  this is the hourly equation of the motion UVRM

### 5.2. Circular motion حركة دائرية

Circular motion is plane motion with constant radius of curvature  $\rho = R$ . The trajectory of the moving object is a circle of radius  $R$ .



- **The position**

The moving point travels from point I to point M, thus the trajectory forms an arc  $\widehat{IM}$ .

By considering an elementary displacement of the moving point from point I to point m, we would have a displacement in the form of an elementary arc Im.

In the right triangle OIm,  $\widehat{Im} = R \sin \theta$

In the right triangle. If  $\theta$  is so small then  $\sin \theta \approx \theta$ .

So  $\widehat{Im} = R \cdot \theta$

- **The speed**

$$v = \frac{d\widehat{Im}}{dt} = R \frac{d\theta}{dt}$$

$R$  is constant, the speed is following the trajectory, so it is written  $\vec{v} = v\vec{u}$  so the vector  $\vec{u}$  would be following the tangent.

$\frac{d\theta}{dt} = \dot{\theta} = \omega$  is the angular velocity السرعة الزاوية

$$v = R \frac{d\theta}{dt} = R \cdot \dot{\theta} = R \cdot \omega$$

Note: The relationship between linear velocity and angular velocity is:  $\mathbf{v} = \mathbf{R}\boldsymbol{\omega}$

- **The acceleration**

$$\vec{a} = \frac{d\vec{v}}{dt} = \frac{dv}{dt}\vec{u} + v \frac{d\vec{u}}{dt}$$

$$\frac{d\vec{u}}{dt} = \frac{d\vec{u}}{d\theta} \frac{d\theta}{dt} \text{ with } \frac{d\vec{u}}{d\theta} = \vec{n}$$

(with  $(\vec{u}, \vec{n})$  the unit vectors in the Fresnet frame and  $\frac{d\theta}{dt} = \omega$ )

### 5.2.1. Uniforme circular motion حركة دائرية منتظمة

In this case the angular velocity  $\omega$  is constant and therefore the linear velocity  $v$  is also constant, then  $a_T = 0$ .

The acceleration in this case is:  $\vec{a} = \vec{a}_N = \frac{v^2}{R} \vec{n}$

### 5.2.2. Uniformly variable circular motion حركة دائرية متغيرة بانتظام

In this case the angular velocity  $\omega$  is not constant and therefore the velocity  $v$  is not constant also, then  $\vec{a} = a_T \vec{u} + a_N \vec{n}$ .

The acceleration in this case is:  $\vec{a} = \frac{dv}{dt} \vec{u} + \frac{v^2}{R} \vec{n} = R \frac{d\omega}{dt} \vec{u} + R\omega^2 \vec{n}$

### 5.3. Sinusoidal or harmonic motion حركة جيبيية

The movement is called sinusoidal or harmonic if its evolution over time is written by the equation:

$$x(t) = A \sin(\omega t + \varphi)$$

A: amplitude,  $\omega$ : angular frequency, and  $\varphi$ : phase.

$$\omega = \frac{2\pi}{T} = 2\pi f$$

T: period and f: frequency

- *The speed*

$$v(t) = \frac{dx(t)}{dt} = A\omega \cos(\omega t + \varphi)$$

- *The acceleration*

$$a(t) = \frac{dv(t)}{dt} = \frac{d^2x(t)}{dt^2} = -A\omega^2 \sin(\omega t + \varphi) \Rightarrow a(t) = \frac{d^2x(t)}{dt^2} = -\omega^2 x(t)$$

**Note:**

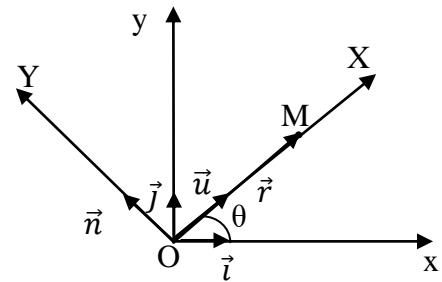
**Another type of movement which is relative movement will be detailed in the next chapter.**

**Proposed exercises about chapter III**

**Exercise 1**

We consider a vector  $\vec{r}$ , of module  $r = OM = a \theta$ , carried by an axis OX making with Ox the variable angle  $\theta$ . We denote by  $\vec{u}$  the unit vector of OX and by  $\vec{n}$  the unit vector directly perpendicular to  $\vec{u}$ .

1. Calculate the express of  $\frac{d\vec{r}}{d\theta}$  in terms of  $a$ ,  $\theta$ ,  $\vec{u}$  and  $\vec{n}$ .
2. Represent this vector  $\frac{d\vec{r}}{d\theta}$ .



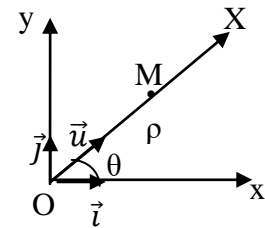
**Exercise 2**

A) A material point M is marked by its cartesian coordinates (x,y).

1. Write x and y in terms of the polar coordinates  $\rho$  and  $\theta$ .
2. Give the expression of the unit vector  $\vec{u}$  as a function of the unit vectors  $\vec{i}$  and  $\vec{j}$ .
3. Calculate  $d\vec{u}/d\theta$ , what does this vector represent?

B) If the position of point M is given by  $\begin{cases} \overline{OM} = t^2 \vec{u} \\ \theta = \omega t \end{cases}$  ( $\omega$  constant)

Find the expression of the velocity vector  $\vec{v}$  in polar coordinates.



**Exercise 3**

Consider a polar coordinate system with the origin O and unit vectors  $\vec{u}_\rho$  and  $\vec{u}_\theta$ .

Let M be a point with coordinates  $(\rho, \theta)$ .

1. Using a detailed diagram, provide the expression for the position vector  $\overline{OM}$  in polar coordinates
2. Give the conversion relationships between polar and Cartesian coordinates.
3. Express the vector  $\vec{A} = 2x\vec{i} + y\vec{j}$  in polar coordinates.
4. Write the elementary displacement vector in polar coordinates.
5. Provide the velocity vector and the acceleration vector in polar coordinates.
6. Find the expression for the elementary area in this coordinate system and deduce the area of a disk with radius R.

## Chapter III: Kinematics of material point

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### Exercise 4

A material point M is identified by its cartesian coordinates (x, y, z).

1. Write down the relationship between cartesian coordinates and cylindrical coordinates (using a diagram).
2. Write the position vector in cylindrical coordinates and deduce the velocity vector in the same coordinate system.

3. If the position of the point is represented in cylindrical coordinates by 
$$\begin{cases} \rho = 4t^2 \\ \theta = \omega t \\ z = \sqrt{t} \end{cases}$$

Find the expression of the velocity vector  $\vec{v}$  in cylindrical coordinates.

### Exercise 5

The differential of the vector  $\vec{r}$ ,  $d\vec{r} = d\vec{l} = dx\vec{i} + dy\vec{j} + dz\vec{k}$  can be expressed in cylindrical coordinates as  $d\vec{r} = \frac{\partial \vec{r}}{\partial \rho} d\rho + \frac{\partial \vec{r}}{\partial \theta} d\theta + \frac{\partial \vec{r}}{\partial z} dz$ .

1. Using the formulas for switching between the two coordinate systems, evaluate the vectors  $\frac{\partial \vec{r}}{\partial \rho}$ ,  $\frac{\partial \vec{r}}{\partial \theta}$  et  $\frac{\partial \vec{r}}{\partial z}$ .
2. Derive the unit vectors  $\vec{U}_\rho, \vec{U}_\theta$  et  $\vec{U}_z$  (cylindrical coordinates) as a function of  $\vec{i}, \vec{j}$  and  $\vec{k}$  (Cartesian coordinates), check that they are orthogonal.
3. Write  $\vec{A} = 2x\vec{i} + y\vec{j} - 2z\vec{k}$  in cylindrical coordinates.

### Exercise 6

Write the vector  $\vec{A} = x \cdot \vec{i} - 2y \cdot \vec{j} + z \cdot \vec{k}$  in cylindrical coordinates i.e. as a function of  $\rho, \theta, z$   $\vec{u}_\rho, \vec{u}_\theta, \vec{u}_z$ . (using passing relations)

If the position of point M is given by 
$$\begin{cases} \vec{OM} = t^3 \vec{u}_\rho + 5t^2 \vec{u}_z \\ \theta = \omega t \end{cases} \quad (\omega \text{ constant})$$

Find the expression of the vectors: speed  $\vec{v}$  and acceleration  $\vec{a}$  in cylindrical coordinates.

### Exercise 7

A body moves along the x axis according to the relation  $x(t) = 2t^3 + 5t^2 + 5$ .

1. Determine the velocity  $v(t)$  and acceleration  $a(t)$  at each instant t.

## Chapter III: Kinematics of material point

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2. Calculate the body's position, velocity and instantaneous acceleration for  $t_1=2s$  and  $t_2=3s$ .
3. Deduce the average velocity and acceleration of the body between  $t_1$  and  $t_2$ .

### Exercise 8

The x and y coordinates of a moving point M in the (oxy) plane vary with time t according to the following relationships:  $x=t+1$  and  $y=(t^2/2)+2$ .

Find :

1. The equation of the trajectory
2. The components of speed and acceleration and their modules.
3. Accelerations: normal  $a_N$  and tangential  $a_T$  and deduce the radius of curvature.
4. The nature of the movement

### Exercise 9

A particle is launched with an initial horizontal speed  $v_0$  according to the time-dependent equations:

$$\begin{cases} x = v_0 t \\ y = \frac{1}{2} g t^2 \end{cases}$$

Determine:

1. The trajectory equation.
2. The components of speed and its module.
3. The components of acceleration and its module.
4. Tangential and normal accelerations.
5. The radius of curvature R of the particle's trajectory.

### Exercise 10

A comet is moving through the solar system. His position is expressed:

$$\overrightarrow{OM} = (t - 1)\vec{i} + \frac{t^2}{2}\vec{j}$$

Where O is the origin of the landmark (the sun) and t represents the time expressed in seconds. We assume that the comet remains in the plane (Oxy)

1. Write the equation of the trajectory
2. Determine the components of the velocity vector  $\vec{v}$  and the acceleration vector  $\vec{a}$  and give the nature of the movement.

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3. Express the expressions for the tangential  $a_T$  and normal  $a_N$  accelerations and deduce the radius of curvature.

### Exercise 11

A particle moves on a trajectory whose trajectory equation is  $y=x^2$  such that at each instant  $v_x=v_0=\text{cst}$ . If  $t=0$ ,  $x_0$ ,  $y_0=0$ .

Determine :

- 1- The  $x(t)$  and  $y(t)$  coordinates of the particle.
- 2- The speed and acceleration of the particle.
- 3- The normal and tangential accelerations as well as the radius of curvature.

### Exercise 12

A body moves on a straight line with an acceleration such that

1. a)  $a=-kv$  ; b)  $a=-kv^2$  where  $k$  is a constant

If at  $t=0$ ;  $v=v_0$  and  $x=x_0$

2. Find for both cases its speed and its displacement in time as well as  $v$  as a function of  $x$

### Exercise 13

A body whose motion is defined by the following velocity components:  $v_x=1$  and  $v_y=2/(t+1)$

Knowing that at  $t=0$   $x=0$  and  $y=2$ .

- 1- What is the equation of the trajectory  $y=f(x)$ .
- 2- Calculate the components of the acceleration.

### Exercise 14

From the ground, a balloon rises with a constant initial speed  $v_0$  (according to  $y$ ). The wind gives the balloon horizontal speed  $V_x=a.y$  ( $a$  constant).

1. Determine the equations of motion  $x(t)$  and  $y(t)$ , Deduce the equation of the trajectory  $y=f(x)$
2. Calculate the accelerations  $a$ ,  $a_N$  and  $a_T$ . Deduce the radius of curvature.

### Exercise 15

A stone is thrown from the top of a 20 m high building, with a horizontal speed of 10m/s.

1. What time does it take for the stone to reach the ground?
  2. At what distance from the building will the stone reach the ground?
- b- With what speed will the stone reach the ground?

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### Exercise 16

The motion of a body is defined by the following velocity components:

$$\begin{cases} v_x = R\omega\cos(\omega t) \\ v_y = R\omega\sin(\omega t) \end{cases}$$

Knowing that  $\omega$  is constant and at  $t=0$ , the moving body is at point M (0, R).

Determine :

1. The components of the acceleration vector and its modulus.
2. The tangential and normal components of acceleration, and deduce the radius of curvature.
3. The components of the position vector and deduce the equation of the trajectory.
4. What is the nature of the motion?

### Exercise 17

A material point M moves along the OX axis with acceleration  $\vec{a} = a\vec{i}$  with  $a > 0$ .

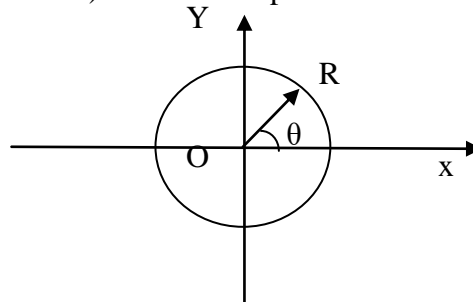
1. Determine the velocity vector knowing that  $v(t=0) = v_0$ .
2. Determine the position vector  $\vec{OM}$  given that  $x(t=0) = x_0$ .
3. Check that:  $v_f^2 - v_i^2 = 2a(x - x_0)$ .
4. What is the condition that  $\vec{a} \cdot \vec{v}$  so that the motion is uniformly accelerated and retarded?

### Exercise 18

Consider a moving point M describing a circle of radius R and center O with an angular speed  $\omega = \frac{d\theta}{dt}$ . At time  $t=0$  point M is at A.

1. Write the coordinates of M as a function of R and  $\theta$ .
2. Calculate the modulus of the speed of point M.
3. Determine the components of the acceleration on the axes Ox and Oy (Cartesian coordinates) on the one hand and on the axes parallel and perpendicular to OM on the other hand (polar coordinates).
4. We assume that  $\alpha = \frac{d\omega}{dt}$  ( $\alpha$  is a non-zero constant). Give the expressions for  $\omega$  and  $\theta$  as a function of time.
5. We recall that at  $t=0$ ,  $\theta_0=0$  et  $\omega=\omega_0$ .

What relationship exists between  $\omega$  and  $\theta$ .





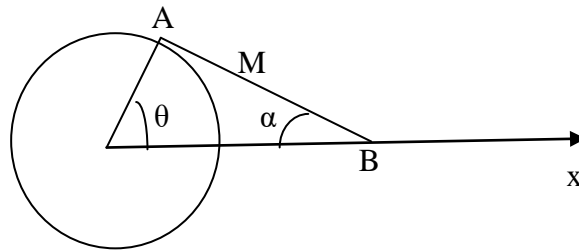
## Chapter III: Kinematics of material point

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### Exercise 19

OA is a rod of length  $L$  animated by a uniform circular movement of angular speed  $\omega$  around the point O. AB is another rod of length  $R$ , articulated at A to OA such that B can move on Ox.

1. Establish the time equations of M (middle of AB)
2. Determine the abscissa of B. Is its movement sinusoidal?
3. Calculate the speed of M
4. Show that if  $r=R$ , the movement of B becomes sinusoidal



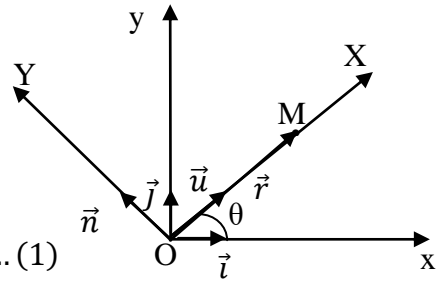
**Correction of exercises about chapter III**

**Exercise 1**

We have  $|\vec{r}| = r = OM = a\theta$

1-  $\frac{d\vec{r}}{d\theta} = f(\theta, a, \vec{u}, \vec{n}) = ?$

According to the diagram  $\begin{cases} \vec{r} = \overline{OM} = a\theta\vec{u} \\ \vec{u} = \cos\theta\vec{i} + \sin\theta\vec{j} \dots \dots (1) \\ \vec{n} = -\sin\theta\vec{i} + \cos\theta\vec{j}. \end{cases}$

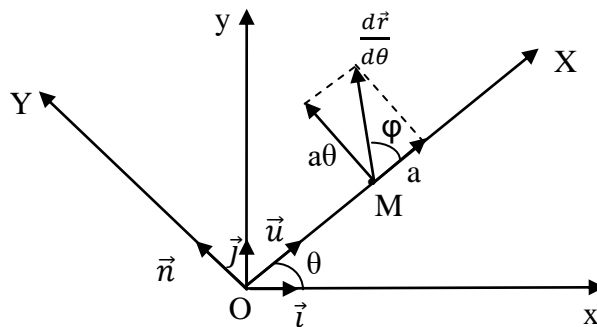


$$\frac{d\vec{r}}{d\theta} = \frac{d}{d\theta}(\theta a \vec{u}) = a\vec{u} + a\theta \frac{d\vec{u}}{d\theta}$$

$$(1) \Rightarrow \frac{d\vec{u}}{d\theta} = -\sin\theta\vec{i} + \cos\theta\vec{j} = \vec{n}$$

$$\Rightarrow \frac{d\vec{r}}{d\theta} = a\vec{u} + a\theta\vec{n}$$

2-  $\frac{d\vec{r}}{d\theta} = a\vec{u} + a\theta\vec{n}$  which originates from point M presents as follows:



3-  $\left| \frac{d\vec{r}}{d\theta} \right| = \sqrt{a^2 + (a\theta)^2} = a\sqrt{1 + \theta^2}$

**Exercise 2**

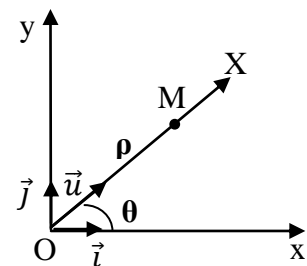
A) A material point M is identified by its Cartesian coordinates (x,y):

Find x and y in terms of polar coordinates rho and theta

$$\overline{OM} = x\vec{i} + y\vec{j} \quad (1)$$

In the other hand  $\overline{OM}$  is written by projection as:

$$\overline{OM} = \rho\cos\theta\vec{i} + \rho\sin\theta\vec{j} \quad (2)$$



$$(1) \text{ and } (2) \Rightarrow \begin{cases} x = \rho \cos \theta \\ y = \rho \sin \theta \end{cases}$$

1- The unit vector  $\vec{u}$  as a function of the unit vectors  $\vec{i}$  and  $\vec{j}$  :

$$\text{we have } \overrightarrow{OM} = |\overrightarrow{OM}| \vec{u} = \rho \vec{u} = \rho \cos \theta \vec{i} + \rho \sin \theta \vec{j}$$

$$\text{so } \vec{u} = \cos \theta \vec{i} + \sin \theta \vec{j} \quad \text{and} \quad \vec{n} = -\sin \theta \vec{i} + \cos \theta \vec{j}$$

$\vec{n}$  and  $\vec{u}$  represent the unit vectors of the polar coordinate basis.

2- Calculate the expression of  $d\vec{u}/d\theta$ , which this vector represents?

$$\frac{d\vec{u}}{d\theta} = \frac{d(\cos \theta \vec{i} + \sin \theta \vec{j})}{d\theta} = -\sin \theta \vec{i} + \cos \theta \vec{j} = \vec{n}$$

$\frac{d\vec{u}}{d\theta}$  represents a unit vector perpendicular to  $\vec{u}$  in the direct direction .

2- The position of point M is given by  $\begin{cases} \overrightarrow{OM} = t^2 \vec{u} \\ \theta = \omega t \end{cases}$  ( $\omega$  constant)

The expression of the velocity vector  $\vec{v}$  in polar coordinates is :

$$\vec{v} = \frac{d\overrightarrow{OM}}{dt} = \frac{d(t^2 \vec{u})}{dt} = 2t \vec{u} + t^2 \frac{d\vec{u}}{dt}$$

$$\frac{d\vec{u}}{dt} = \frac{d\vec{u}}{d\theta} \cdot \frac{d\theta}{dt} = \vec{n} \cdot \omega$$

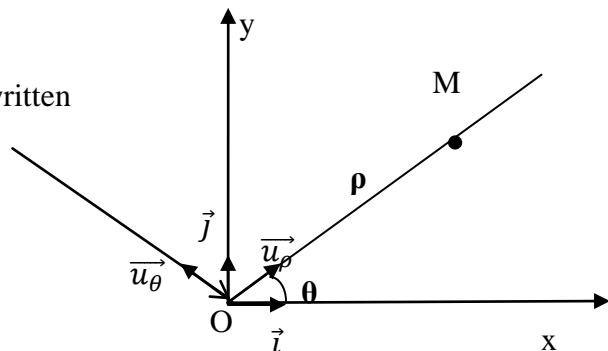
$$\vec{v} = 2t \cdot \vec{u} + t^2 \cdot \omega \cdot \vec{n}$$

### Exercise 3

B- The polar coordinates are  $\rho$  and  $\theta$  ; with  $\rho = \|\overrightarrow{OM}\|$  ;  $0 < \rho < R$  and the  $\theta = (\overrightarrow{Ox}, \overrightarrow{OM})$  with  $0 < \theta < 2\pi$ .

1- The vector  $\overrightarrow{OM}$  in polar coordinates is written

$$\text{as following : } \overrightarrow{OM} = \rho \vec{u}_\rho$$



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2- the transition relationships between polar and Cartesian coordinates.

$$\begin{cases} \cos\theta = \frac{x_M}{\rho} \\ \sin\theta = \frac{y_M}{\rho} \end{cases} \Rightarrow \begin{cases} x_M = \rho \cos\theta \\ y_M = \rho \sin\theta \end{cases}$$

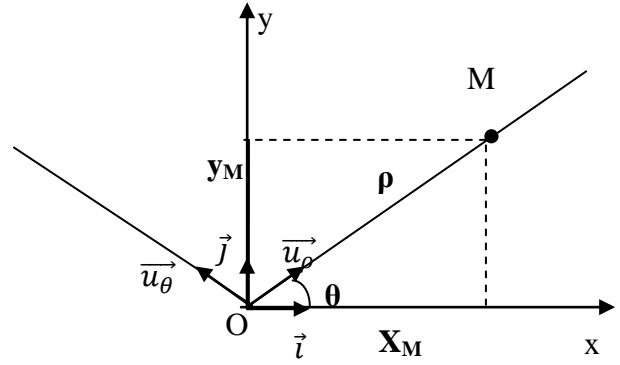
So the vector  $\overrightarrow{OM}$  in coordinates cartesian

$$\text{is written } \overrightarrow{OM} = x_M \vec{i} + y_M \vec{j}$$

We had :  $\overrightarrow{OM} = \rho \vec{u}_\rho$  (in polar coordinates)

$$\text{Then } \overrightarrow{OM} = \rho(\cos\theta \vec{i} + \sin\theta \vec{j})$$

By identification  $\vec{u}_\rho = \cos\theta \vec{i} + \sin\theta \vec{j}$  and  $\vec{u}_\theta = \frac{d\vec{u}_\rho}{d\theta} = -\sin\theta \vec{i} + \cos\theta \vec{j}$



3- The writing of the vector  $\vec{A} = 2x\vec{i} + y\vec{j}$  in polar coordinates

$$\text{we have } \begin{cases} x_M = \rho \cos\theta \\ y_M = \rho \sin\theta \end{cases} \quad \text{and} \quad \begin{cases} \vec{u}_\rho = \cos\theta \vec{i} + \sin\theta \vec{j} \\ \vec{u}_\theta = -\sin\theta \vec{i} + \cos\theta \vec{j} \end{cases}$$

By using the passage table

	$\vec{u}_\rho$	$\vec{u}_\theta$
$\vec{i}$	$\cos\theta$	$-\sin\theta$
$\vec{j}$	$\sin\theta$	$\cos\theta$

$$\text{So } \vec{i} = \cos\theta \vec{u}_\rho - \sin\theta \vec{u}_\theta \quad \text{and} \quad \vec{j} = \sin\theta \vec{u}_\rho + \cos\theta \vec{u}_\theta$$

The vector  $\vec{A}$  is then written as;

$$\vec{A} = 2\rho \cos\theta (\cos\theta \vec{u}_\rho - \sin\theta \vec{u}_\theta) + \rho \sin\theta (\sin\theta \vec{u}_\rho + \cos\theta \vec{u}_\theta)$$

$$\Rightarrow \vec{A} = \rho(1 + \cos^2\theta) \vec{u}_\rho - \rho \sin\theta \cos\theta \vec{u}_\theta$$

4- The vector of elementary displacement in polar coordinates

$$d\overrightarrow{OM} = d(\rho \vec{u}_\rho) = d\rho \vec{u}_\rho + \rho d\vec{u}_\rho \quad \text{with} \quad d\vec{u}_\rho = \frac{d\vec{u}_\rho}{d\theta} d\theta = \vec{u}_\theta d\theta$$

$$\text{So } d\overrightarrow{OM} = d\rho \vec{u}_\rho + \rho d\theta \vec{u}_\theta$$

5- The velocity vector and the acceleration vector in polar coordinates.

$$\text{The velocity vector in polar coordinates: } \vec{v} = \frac{d\overrightarrow{OM}}{dt} = \frac{d\rho}{dt} \vec{u}_\rho + \rho \frac{d\theta}{dt} \vec{u}_\theta$$

The acceleration vector in polar coordinates:

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$$\vec{a} = \frac{d\vec{v}}{dt} = \frac{d^2\rho}{dt^2}\vec{u}_\rho + \frac{d\rho}{dt}\frac{d\vec{u}_\rho}{dt} + \frac{d\rho}{dt}\frac{d\theta}{dt}\vec{u}_\theta + \rho\frac{d^2\theta}{dt^2}\vec{u}_\theta + \rho\frac{d\theta}{dt}\frac{d\vec{u}_\theta}{dt}$$

$$\text{With } \frac{d\vec{u}_\rho}{dt} = \frac{d\vec{u}_\rho}{d\theta} \cdot \frac{d\theta}{dt} = \frac{d\theta}{dt}\vec{u}_\theta \text{ et } \frac{d\vec{u}_\theta}{dt} = \frac{d\vec{u}_\theta}{d\theta} \cdot \frac{d\theta}{dt} = -\frac{d\theta}{dt}\vec{u}_\rho$$

$$\text{so } \vec{a} = \frac{d\vec{v}}{dt} = \frac{d^2\rho}{dt^2}\vec{u}_\rho + \frac{d\rho}{dt}\frac{d\theta}{dt}\vec{u}_\theta + \frac{d\rho}{dt}\frac{d\theta}{dt}\vec{u}_\theta + \rho\frac{d^2\theta}{dt^2}\vec{u}_\theta - \rho\frac{d\theta}{dt}\frac{d\theta}{dt}\vec{u}_\rho$$

$$\Rightarrow \vec{a} = \frac{d^2\rho}{dt^2}\vec{u}_\rho + 2\frac{d\rho}{dt}\frac{d\theta}{dt}\vec{u}_\theta + \rho\frac{d^2\theta}{dt^2}\vec{u}_\theta - \rho\left(\frac{d\theta}{dt}\right)^2\vec{u}_\rho$$

6- The expression of the elementary surface in the polar frame:

$$ds = dl_1 \cdot dl_2 \text{ and } d\vec{OM} = d\rho\vec{u}_\rho + \rho d\theta\vec{u}_\theta = dl_1\vec{u}_\rho + dl_2\vec{u}_\theta$$

with  $dl_1$  is the variation of  $\rho$  along  $\vec{u}_\rho$  which is  $d\rho$  and  $dl_2$  is the variation of  $\theta$  along  $\vec{u}_\theta$

$$ds = d\rho \cdot \rho d\theta$$

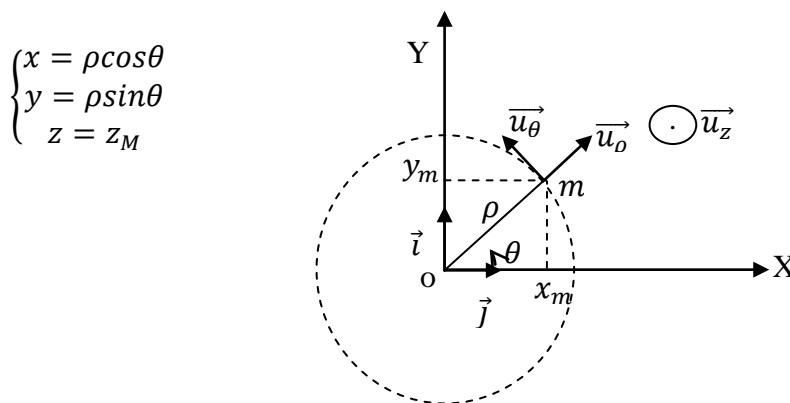
The surface of a disk with radius R.

$$S = \iint d\rho \cdot \rho d\theta = \int_0^R \rho d\rho \int_0^{2\pi} d\theta = \frac{R^2}{2} 2\pi = \pi R^2$$

### Exercise 4

1. A material point M is identified by its Cartesian coordinates (x, y, z).

Write the relationship between Cartesian coordinates and polar coordinates.



2. Find the expression of the position vector and deduce the velocity  $\vec{v}$  of point M in cylindrical coordinates.

$$\vec{OM} = \rho\vec{U}_\rho + z\vec{U}_z$$

$$\Rightarrow \vec{v} = \frac{d\vec{OM}}{dt} = \frac{d\rho}{dt}\vec{U}_\rho + \rho\frac{d\vec{U}_\rho}{dt} + \frac{dz}{dt}\vec{U}_z + z\frac{d\vec{U}_z}{dt}$$

$$\text{We have } \frac{d\vec{U}_z}{dt} = 0 \Rightarrow \vec{v} = \dot{\rho}\vec{U}_\rho + \rho\frac{d\theta}{dt}\frac{d\vec{U}_\rho}{d\theta} + \dot{z}\vec{U}_z$$

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$$\Rightarrow \vec{v} = \dot{\rho}\vec{U}_\rho + \rho\dot{\theta}\vec{U}_\theta + \dot{z}\vec{U}_z$$

3. A velocity vector  $\vec{v}$  of point M in cylindrical coordinates:

$$\text{We have } \begin{cases} \rho = 4t^2 \\ \theta = \omega t \\ z = \sqrt{t} \end{cases} \quad \text{Hence } \begin{cases} \frac{d\rho}{dt} = 8t \\ \frac{d\theta}{dt} = \omega \\ \frac{dz}{dt} = \frac{1}{2\sqrt{t}} \end{cases}$$

$$\Rightarrow \vec{v} = \dot{\rho}\vec{U}_\rho + \rho\dot{\theta}\frac{d\vec{U}_\rho}{d\theta} + \dot{z}\vec{U}_z = 8t\vec{U}_\rho + 4t^2 \cdot \omega \cdot \vec{U}_\theta + \frac{1}{2\sqrt{t}}\vec{U}_z$$

### Exercise 5

The differential of vector  $\vec{r}$ ,  $d\vec{r} = d\vec{l} = dx\vec{i} + dy\vec{j} + dz\vec{k}$  can be expressed in cylindrical coordinates as  $d\vec{r} = \frac{\partial \vec{r}}{\partial \rho} d\rho + \frac{\partial \vec{r}}{\partial \theta} d\theta + \frac{\partial \vec{r}}{\partial z} dz$ .

1. We are looking for the vectors  $\frac{\partial \vec{r}}{\partial \rho}$ ,  $\frac{\partial \vec{r}}{\partial \theta}$  et  $\frac{\partial \vec{r}}{\partial z}$ .

We are  $\vec{r} = x\vec{i} + y\vec{j} + z\vec{k}$

- The displacement vector in cartesian coordinates (x, y, z) :

$$d\vec{r} = d\vec{l} = dx\vec{i} + dy\vec{j} + dz\vec{k}$$

- The displacement vector in cylindrical coordinates ( $\rho, \theta, z$ ) :

$$d\vec{r} = \frac{\partial \vec{r}}{\partial \rho} d\rho + \frac{\partial \vec{r}}{\partial \theta} d\theta + \frac{\partial \vec{r}}{\partial z} dz$$

Relationships between cartesian coordinates (x, y, z) and cylindrical coordinates ( $\rho, \theta, z$ ) are :

$$\begin{cases} x = \rho \cos \theta \\ y = \rho \sin \theta \\ z = z_M \end{cases} \Rightarrow \begin{cases} dx = d\rho \cdot \cos \theta - \rho \cdot \sin \theta \cdot d\theta \\ dy = d\rho \cdot \sin \theta + \rho \cdot \cos \theta \cdot d\theta \\ dz = dz_M \end{cases}$$

$$\Rightarrow d\vec{r} = d\vec{l} = (d\rho \cdot \cos \theta - \rho \cdot \sin \theta \cdot d\theta)\vec{i} + (d\rho \cdot \sin \theta + \rho \cdot \cos \theta \cdot d\theta)\vec{j} + dz\vec{k}$$

$$\Rightarrow d\vec{r} = (\cos \theta \cdot \vec{i} + \sin \theta \cdot \vec{j})d\rho + (-\rho \sin \theta \vec{i} + \rho \cdot \cos \theta \cdot \vec{j})d\theta + dz\vec{k} \dots \dots \dots (1)$$

$$\Rightarrow d\vec{r} = \left(\frac{\partial \vec{r}}{\partial \rho}\right) d\rho + \left(\frac{\partial \vec{r}}{\partial \theta}\right) d\theta + \left(\frac{\partial \vec{r}}{\partial z}\right) dz \dots \dots \dots (2)$$

With identification between (1) and (2) we'll have :

$$\Rightarrow \begin{cases} \frac{\partial \vec{r}}{\partial \rho} = \cos\theta \cdot \vec{i} + \sin\theta \cdot \vec{j} \\ \frac{\partial \vec{r}}{\partial \theta} = -\rho \sin\theta \vec{i} + \rho \cdot \cos\theta \cdot \vec{j} \\ \frac{\partial \vec{r}}{\partial z} = \vec{k} \end{cases}$$

2. Deduce Unit Vectors  $\vec{U}_\rho, \vec{U}_\theta$  and  $\vec{U}_z$  (cylindrical coordinates) as function of  $\vec{i}, \vec{j}$  and  $\vec{k}$  (Cartesian coordinates) :

The displacement vector in cylindrical coordinates is written:

$$d\vec{r} = d\rho \vec{U}_\rho + \rho d\theta \vec{U}_\theta + dz \vec{k} \dots \dots \dots (3)$$

$$(1) \text{ and } (3) \Rightarrow \begin{cases} \vec{U}_\rho = \frac{\partial \vec{r}}{\partial \rho} = \cos\theta \cdot \vec{i} + \sin\theta \cdot \vec{j} \\ \vec{U}_\theta = \frac{1}{\rho} \frac{\partial \vec{r}}{\partial \theta} = -\sin\theta \vec{i} + \cos\theta \cdot \vec{j} \\ \vec{U}_z = \frac{\partial \vec{r}}{\partial z} = \vec{k} \end{cases}$$

**Note :**

The unit vectors of the Cartesian coordinates base can be written as a function of the unit vectors of the cylindrical coordinates base from the table below:

	$\vec{i}$	$\vec{j}$	$\vec{k}$
$\vec{u}_\rho$	Cos $\theta$	Sin $\theta$	0
$\vec{u}_\theta$	-sin $\theta$	Cos $\theta$	0
$\vec{u}_z$	0	0	1

$$\Rightarrow \begin{cases} \vec{i} = \cos\theta \vec{u}_\rho - \sin\theta \vec{u}_\theta \\ \vec{j} = \sin\theta \vec{u}_\rho + \cos\theta \vec{u}_\theta \\ \vec{k} = \vec{u}_z \end{cases}$$

3. Checking that they are orthogonal?

$$\Rightarrow \begin{cases} |\vec{U}_\rho| = \sqrt{\cos^2\theta + \sin^2\theta} = 1 \\ |\vec{U}_\theta| = \sqrt{(-\sin\theta)^2 + \cos^2\theta} = 1 \\ |\vec{U}_z| = |\vec{k}| = 1 \end{cases}$$

Hence  $\vec{U}_\rho, \vec{U}_\theta$  et  $\vec{U}_z$ , are the unit vectos.

We have  $\vec{U}_\rho \cdot \vec{U}_\theta = 0, \vec{U}_\rho \cdot \vec{U}_z = 0$  and  $\vec{U}_z \cdot \vec{U}_\theta = 0$

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So  $\vec{U}_\rho, \vec{U}_\theta,$  and  $\vec{U}_z$  are orthogonal vectors.

Therefore the vectors  $\vec{U}_\rho, \vec{U}_\theta, \vec{U}_z$  form an orthonormal reference frame.

4. Write  $\vec{A} = 2x\vec{i} + y\vec{j} - 2z\vec{k}$  in cylindrical coordinates.

$$\text{We have } \begin{cases} x = \rho \cos \theta \\ y = \rho \sin \theta \\ z = z_M \end{cases} \text{ and } \begin{cases} \vec{i} = \cos \theta \vec{u}_\rho - \sin \theta \vec{u}_\theta \\ \vec{j} = \sin \theta \vec{u}_\rho + \cos \theta \vec{u}_\theta \\ \vec{k} = \vec{u}_z \end{cases}$$

So  $\vec{A} = 2x\vec{i} + y\vec{j} - 2z\vec{k}$  is written by :

$$\begin{aligned} \Rightarrow \vec{A} &= 2\rho \cos \theta (\cos \theta \vec{u}_\rho - \sin \theta \vec{u}_\theta) + \rho \sin \theta (\sin \theta \vec{u}_\rho + \cos \theta \vec{u}_\theta) - 2z\vec{k} \\ \Rightarrow \vec{A} &= (2\rho \cos^2 \theta + \rho \sin^2 \theta) \vec{u}_\rho + (-2\rho \cos \theta \sin \theta + \rho \sin \theta \cos \theta) \vec{u}_\theta - 2z\vec{k} \\ \Rightarrow \vec{A} &= (\cos^2 \theta + 1) \rho \vec{u}_\rho - \rho \cos \theta \sin \theta \vec{u}_\theta - 2z\vec{u}_z \end{aligned}$$

### Exercise 6

Writing the vector  $\vec{A} = x \cdot \vec{i} - 2 \cdot y \cdot \vec{j} + z \cdot \vec{k}$  in cylindrical coordinates:

Transit relations between cylindrical and Cartesian coordinates

$$\begin{cases} x_M = \rho \cos \theta \\ y_M = \rho \sin \theta \\ z_M = z \end{cases}$$

So the vector  $\vec{OM}$  in Cartesian coordinates is written  $\vec{OM} = x_M \vec{i} + y_M \vec{j} + z_M \vec{k}$

$$\vec{OM} = \rho (\cos \theta \vec{i} + \sin \theta \vec{j}) + z \vec{k}$$

We have:  $\vec{OM} = \rho \vec{u}_\rho + z \vec{u}_z$  (in cylindrical coordinates)

By identification  $\vec{u}_\rho = \cos \theta \vec{i} + \sin \theta \vec{j}$ ,  $\vec{u}_z = \vec{k}$

$$\text{and } \vec{u}_\theta = \frac{d\vec{u}_\rho}{d\theta} = -\sin \theta \vec{i} + \cos \theta \vec{j}$$

By using the table passage

$$\vec{i} = \cos \theta \vec{u}_\rho - \sin \theta \vec{u}_\theta$$

$$\vec{j} = \sin \theta \vec{u}_\rho + \cos \theta \vec{u}_\theta$$

$$\vec{k} = \vec{u}_z$$

	$\vec{u}_\rho$	$\vec{u}_\theta$	$\vec{u}_z$
$\vec{i}$	Cos $\theta$	-sin $\theta$	0
$\vec{j}$	Sin $\theta$	cos $\theta$	0
$\vec{k}$	0	0	1



The vector  $\vec{A}$  is written by:

$$\vec{A} = \rho \cos \theta (\cos \theta \vec{u}_\rho - \sin \theta \vec{u}_\theta) - 2 \rho \sin \theta (\sin \theta \vec{u}_\rho + \cos \theta \vec{u}_\theta) + z \vec{u}_z$$

$$\Rightarrow \vec{A} = \rho (\cos^2 \theta - 2 \sin^2 \theta) \vec{u}_\rho - 3 \rho \sin \theta \cos \theta \vec{u}_\theta + z \vec{u}_z$$

B. The position of point M is given by  $\begin{cases} \vec{OM} = t^3 \vec{u}_\rho + 5t^2 \vec{u}_z \\ \theta = \omega t \end{cases}$  ( $\omega$  constant)

1. Speed is written :  $\vec{v} = \frac{d\vec{OM}}{dt} = 3t^2 \vec{u}_\rho + t^3 \frac{d\vec{u}_\rho}{dt} + 10t \vec{u}_z + 10t \frac{d\vec{u}_z}{dt}$

$$\text{With } \frac{d\vec{u}_\rho}{dt} = \frac{d\vec{u}_\rho}{d\theta} \cdot \frac{d\theta}{dt} = \omega \vec{u}_\theta \text{ and } \frac{d\vec{u}_z}{dt} = \vec{0}$$

$$\Rightarrow \vec{v} = 3t^2 \vec{u}_\rho + t^3 \omega \vec{u}_\theta + 10t \vec{u}_z$$

$$\vec{a} = \frac{d\vec{v}}{dt} = \frac{d^2 \vec{OM}}{dt^2} = 6t \vec{u}_\rho + 3t^2 \frac{d\vec{u}_\rho}{dt} + 3t^2 \omega \vec{U}_\theta + t^3 \omega \frac{d\vec{U}_\theta}{dt} + 10 \vec{U}_z + 10t \frac{d\vec{u}_z}{dt}$$

$$\Rightarrow \vec{a} = 6t \vec{u}_\rho + 3t^2 \omega \vec{u}_\theta + 3t^2 \omega \vec{U}_\theta - t^3 \omega \omega \vec{U}_\rho + 10 \vec{U}_z$$

$$\text{With } \frac{d\vec{u}_\rho}{dt} = \frac{d\vec{u}_\rho}{d\theta} \cdot \frac{d\theta}{dt} = \omega \vec{u}_\theta, \quad \frac{d\vec{u}_\theta}{dt} = \frac{d\vec{u}_\theta}{d\theta} \cdot \frac{d\theta}{dt} = -\omega \vec{U}_\rho$$

$$\Rightarrow \vec{a} = 6t \vec{u}_\rho + 6t^2 \omega \vec{u}_\theta - t^3 \omega^2 \vec{U}_\rho + 10 \vec{U}_z$$

$$\text{So } \vec{a} = (6t - t^3 \omega^2) \vec{u}_\rho + 6t^2 \omega \vec{u}_\theta + 10 \vec{U}_z$$

### Exercise 7

a- we have  $x(t) = 2t^3 + 5t^2 + 5$  so :

$$\text{The velocity: } v(t) = \frac{dx}{dt} = 6t^2 + 10t$$

$$\text{The acceleration: } a(t) = \frac{dv(t)}{dt} = 12t + 10$$

b- The body's position at time  $t_1 = 2s$ , as well as its instantaneous velocity and acceleration:

$$\text{The position : } x(2) = 2(2)^3 + 5(2)^2 + 5 = 41m$$

$$\text{Instantaneous speed: } v(2) = 6(2)^2 + 10(2) = 44m/s$$

$$\text{Instantaneous acceleration: } a(2) = 12(2) + 10 = 34m/s^2$$

-The body's position at time  $t_2 = 3s$ , as well as its instantaneous velocity and acceleration:

$$\text{Position : } x(3) = 2(3)^3 + 5(3)^2 + 5 = 104m$$

$$\text{Instantaneous speed: } v(3) = 6(3)^2 + 10(3) = 84m/s$$

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Instantaneous acceleration:  $a(3)=12(3)+10=46\text{m/s}^2$

c- We deduce the speed and average acceleration of the body between  $t_1$  and  $t_2$ :

$$\text{Average speed: } v_{moy} = \frac{\Delta x}{\Delta t} = \frac{x(t_2)-x(t_1)}{t_2-t_1} \Rightarrow v_{moy} = \frac{104-41}{3-2} = 63\text{m/s}$$

Average acceleration :

$$a_{moy} = \frac{\Delta v}{\Delta t} = \frac{v(t_2) - v(t_1)}{t_2 - t_1} \Rightarrow a_{moy} = \frac{84 - 44}{3 - 2} = 40\text{m/s}^2$$

### Exercise 8

The coordinates of a moving point M in the plane (oxy) are written as:

$$x(t)=t+1 \text{ and } y(t)=(t^2/2)+2$$

a- The equation of the trajectory is then written :

To find the equation of the trajectory, simply find the relationship between  $x(t)$  and  $y(t)$ .

To do this, deduce the time from one equation,  $x(t)$  or  $y(t)$ , and replace it in the other equation).

Here, we'll write  $t$  as a function of  $x$  :

$$t=x-1 \text{ so } y = \frac{(x-1)^2}{2} + 2 = \frac{x^2}{2} - x + \frac{5}{2}$$

$$\text{The equation of the trajectory is : } \mathbf{y(x)} = \frac{x^2}{2} - x + \frac{5}{2}$$

b- Components of velocity and acceleration vectors:

- The velocity :  $\overrightarrow{v(t)} = v_x(t)\vec{i} + v_y(t)\vec{j}$

$$\begin{cases} v_x(t) = \frac{dx(t)}{dt} = 1 \\ v_y(t) = \frac{dy(t)}{dt} = t \end{cases}$$

The velocity is written by  $\overrightarrow{v(t)} = \vec{i} + t\vec{j}$

The velocity module:  $|\vec{v}(t)| = \sqrt{1 + t^2}$

- The acceleration:  $\overrightarrow{a(t)} = a_x(t)\vec{i} + a_y(t)\vec{j}$

$$\begin{cases} a_x(t) = \frac{dv_x(t)}{dt} = 0 \\ a_y(t) = \frac{dv_y(t)}{dt} = 1 \end{cases}$$

So  $\overrightarrow{a(t)} = \vec{j}$

The acceleration module is:  $|\vec{a}(t)| = 1$

c- Normal and tangential acceleration:

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- Tangential acceleration

$$a_T = \frac{d|\vec{v}(t)|}{dt} \quad \text{with} \quad |\vec{v}(t)| = \sqrt{v_x^2 + v_y^2} = \sqrt{1 + t^2}$$

$$a_T = \frac{d(\sqrt{1 + t^2})}{dt} = \frac{2t}{2\sqrt{1 + t^2}}$$

$$a_T = \frac{t}{\sqrt{1+t^2}} \quad \text{because} \quad (U^n)' = nU' U^{n-1}$$

- Normal acceleration :

The accelerations  $a_N$  and  $a_T$  are the normal and tangential components of the acceleration.  $\vec{a}$

$$(\vec{a} = a_T \vec{U}_T + a_N \vec{U}_N \Rightarrow |\vec{a}| = \sqrt{a_T^2 + a_N^2})$$

We have the shape of a right triangle, by applying Pitagort's relation.

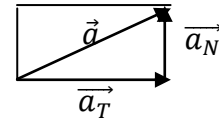
$$a^2 = a_T^2 + a_N^2$$

$$\text{So} \quad a_N^2 = a^2 - a_T^2 \quad \text{or} \quad |\vec{a}| = \sqrt{a_T^2 + a_N^2}$$

$$a_N^2 = 1 - \left(\frac{t}{\sqrt{1+t^2}}\right)^2 = 1 - \frac{t^2}{1+t^2}$$

$$a_N^2 = \frac{1}{1+t^2}$$

$$\text{So} \quad a_N = \frac{1}{\sqrt{1+t^2}} = \frac{1}{v}$$



- The radius of curvature:  $a_N = \frac{v^2}{R} = \frac{1}{v} \Rightarrow R = v^3 = (1 + t^2)^{\frac{3}{2}}$

- c- The nature of movement

$$\vec{a}(t) \cdot \vec{v}(t) = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \cdot \begin{pmatrix} 1 \\ t \end{pmatrix} = 1(0) + t(1) = t > 0$$

The motion is then uniformly accelerated.

### Exercise 9

The x and y coordinates of a mobile point M in the (xy) plane vary with time t according to

$$\text{the following relationships:} \quad \begin{cases} x = v_0 t \\ y = \frac{1}{2} g t^2 \end{cases}$$

- 1- The equation of the trajectory is then written as follows:

$$\text{Here, we will express t as a function of x:} \quad t = \frac{x}{v_0} \quad \text{So} \quad y = \frac{1}{2} g \left(\frac{x}{v_0}\right)^2 = \frac{g}{2v_0^2} x^2$$

$$\text{The equation of the trajectory is:} \quad \mathbf{y(x)} = \frac{g}{2v_0^2} x^2$$

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The components of the velocity

$$\begin{cases} v_x(t) = \frac{dx(t)}{dt} = v_0 \\ v_y(t) = \frac{dy(t)}{dt} = gt \end{cases}$$

The velocity is expressed as:  $\vec{v}(t) = v_0\vec{i} + gt\vec{j}$

The magnitude of the velocity:  $|\vec{v}(t)| = \sqrt{v_0^2 + (gt)^2} = \sqrt{v_0^2 + g^2t^2}$

The components of the acceleration

$$\begin{cases} a_x(t) = \frac{dv_x(t)}{dt} = 0 \\ a_y(t) = \frac{dv_y(t)}{dt} = g \end{cases} \quad \text{The acceleration is expressed as: } \vec{a}(t) = g\vec{j}$$

The magnitude of the acceleration  $|\vec{a}(t)| = g$

2- The nature of the movement

$$\vec{a}(t) \cdot \vec{v}(t) = v_0(0) + gt(g) = g^2t > 0$$

The movement in this case is uniformly accelerated.

3- Normal and tangential accelerations.

- Tangential acceleration:

$$a_T = \frac{d|\vec{v}(t)|}{dt} \quad \text{avec } |\vec{v}(t)| = \sqrt{v_0^2 + g^2t^2} \quad \text{so } a_T = \frac{d(\sqrt{v_0^2 + g^2t^2})}{dt} = \frac{2g^2t}{2\sqrt{v_0^2 + g^2t^2}}$$

$$a_T = \frac{g^2t}{\sqrt{v_0^2 + g^2t^2}} = \frac{g^2t}{v}$$

- Normal acceleration

The accelerations  $a_N$  and  $a_T$  are the normal and tangential components of the acceleration vector  $\vec{a}$ .

$$(\vec{a} = a_T\vec{U}_T + a_N\vec{U}_N) \Rightarrow a^2 = a_T^2 + a_N^2 \quad \text{so } a_N^2 = a^2 - a_T^2$$

$$a_N^2 = g^2 - \left( \frac{g^2t}{\sqrt{v_0^2 + g^2t^2}} \right)^2 = g^2 - \frac{g^4t^2}{v_0^2 + g^2t^2} \Rightarrow a_N^2 = \frac{g^2v_0^2 + g^4t^2 - g^4t^2}{v_0^2 + g^2t^2}$$

$$\text{So } a_N = \sqrt{\frac{g^2v_0^2}{v_0^2 + g^2t^2}} = \frac{gv_0}{v}$$

The radius of curvature

$$a_N = \frac{v^2}{R} = \frac{gv_0}{v} \Rightarrow R = \frac{v^3}{gv_0} = \frac{(v_0^2 + g^2t^2)^{\frac{3}{2}}}{gv_0}$$

### Exercise 10

$$\overrightarrow{OM} = (t - 1)\vec{i} + \frac{t^2}{2}\vec{j}$$

1- The equation of the trajectory

$$\overrightarrow{OM} = (t - 1)\vec{i} + \frac{t^2}{2}\vec{j} \Rightarrow \begin{cases} x = t - 1 \\ y = \frac{t^2}{2} \end{cases}$$

$$t = x + 1 \Rightarrow y = \frac{(x+1)^2}{2}$$

2- The components of velocity and acceleration, and their magnitudes :

- Velocity

$$\begin{cases} v_x = \frac{dx}{dt} \\ v_y = \frac{dy}{dt} \end{cases} \Rightarrow \begin{cases} v_x = 1 \\ v_y = t \end{cases} \quad \vec{v} = v = \vec{i} + t\vec{j} \Rightarrow |\vec{v}| = \sqrt{1 + t^2}$$

- Acceleration

$$\begin{cases} a_x = \frac{dv_x}{dt} \\ a_y = \frac{dv_y}{dt} \end{cases} \Rightarrow \begin{cases} a_x = 0 \\ a_y = 1 \end{cases} \quad \vec{a} = 1\vec{j} \Rightarrow |\vec{a}| = a = 1$$

3- The nature of the movement

$\vec{a} \cdot \vec{v} = t > 0$  so The movement in this case is uniformly accelerated.

Normal and tangential accelerations.

- Tangential acceleration

$$a_T = \frac{d|\vec{v}|}{dt} = \frac{d(\sqrt{1 + t^2})}{dt} = \frac{t}{\sqrt{t^2 + 1}}$$

- Normal acceleration

$$\text{We have } a^2 = a_T^2 + a_N^2 \quad \text{so } a_N^2 = a^2 - a_T^2$$

$$a_N^2 = 1 - \frac{t^2}{t^2 + 1} \Rightarrow a_N^2 = \frac{1}{t^2 + 1} \Rightarrow a_N = \frac{1}{\sqrt{t^2 + 1}}$$

4- The radius of curvature

$$a_N = \frac{v^2}{R} \Rightarrow R = \frac{v^2}{a_N} = \frac{v^3}{1} = v^3$$

### Exercise 11

A particle moves along a trajectory with the equation  $y = x^2$  in such a way that at each moment  $v_x = v_0 = \text{const}$ . If at  $t=0$ ,  $x_0, y_0=0$ .

a- Let's find the coordinates  $x(t)$  and  $y(t)$  of the particle.

We have the following (Ox):

$$v_x = v_0 = \frac{dx}{dt} \Rightarrow \int_0^x dx = \int_0^t v_0 dt$$

$$\Rightarrow \mathbf{x}(t) = v_0 t$$

On the other hand:  $y=x^2 \Rightarrow y(t) = v_0^2 t^2$

$$\text{So } \begin{cases} \mathbf{x}(t) = v_0 t \\ y(t) = v_0^2 t^2 \end{cases}$$

5- The components of velocity and acceleration,

- Velocity

$$\begin{cases} v_x = \frac{dx}{dt} = v_0 \\ v_y = \frac{dy}{dt} = 2v_0^2 t \end{cases} \Rightarrow \overrightarrow{v}(t) = v_0 \vec{i} + 2v_0^2 t \vec{j}$$

The velocity module  $|\overrightarrow{v}(t)| = \sqrt{v_0^2 + 4v_0^4 t^2}$

- The acceleration

$$\begin{cases} a_x = \frac{dv_x}{dt} = 0 \\ a_y = \frac{dv_y}{dt} = 2v_0^2 \end{cases} \Rightarrow \overrightarrow{a}(t) = 2v_0^2 \vec{j}$$

The acceleration module  $|\overrightarrow{a}(t)| = 2v_0^2$

b- Normal and tangential acceleration

- The tangential acceleration

$$a_T = \frac{d|\overrightarrow{v}(t)|}{dt} = \frac{4v_0^4 t}{\sqrt{v_0^2 + 4v_0^4 t^2}}$$

- The normal acceleration

$$a_N^2 = a^2 - a_T^2 \Rightarrow a_N^2 = 4v_0^4 - \frac{16v_0^8 t^2}{v_0^2 + 4v_0^4 t^2}$$

$$\Rightarrow a_N^2 = \frac{4v_0^6}{v_0^2 + 4v_0^4 t^2}$$

$$\text{So } a_N = \frac{2v_0^3}{\sqrt{v_0^2 + 4v_0^4 t^2}} = \frac{2v_0^3}{v}$$

- The radius of curvature

$$a_N = \frac{v^2}{R} = \frac{2v_0^3}{v} \Rightarrow R = \frac{v^3}{2v_0^3}$$

### Exercise 12

- *1<sup>st</sup> case :  $a = -kv$*

$K = \text{constant}$ , at  $t=0$ ,  $v=v_0$  et  $x=x_0$

We have the acceleration, seeking the velocity as a function of time.

$$a = \frac{dv}{dt} = -kv \Rightarrow \frac{dv}{v} = -kdt$$

$$\text{So } \int_{v_0}^v \frac{dv}{v} = \int_0^t -kdt \Rightarrow \ln v - \ln v_0 = -kt$$

$$\Rightarrow \ln \frac{v}{v_0} = -kt$$

$$\text{Then } \frac{v}{v_0} = e^{-kt} \text{ and } v(t) = v_0 e^{-kt}$$

We have found the velocity, now seeking the position as a function of time.

$$v = \frac{dx}{dt} = v_0 e^{-kt} \Rightarrow dx = v_0 e^{-kt} dt$$

$$\text{So } \int_{x_0}^x dx = \int_0^t v_0 e^{-kt} dt \Rightarrow x - x_0 = \frac{v_0}{k} (1 - e^{-kt})$$

$$\text{Then } x(t) = x_0 + \frac{v_0}{k} (1 - e^{-kt})$$

We need to find the relationship between velocity and position.

$$\text{We have } \frac{v}{v_0} = e^{-kt} \text{ and } x = x_0 + \frac{v_0}{k} (1 - e^{-kt})$$

$$\text{So } x = x_0 + \frac{v_0}{k} \left(1 - \frac{v}{v_0}\right) = x_0 + \frac{v_0}{k} \left(\frac{v_0 - v}{v_0}\right)$$

$$x = x_0 + \frac{v_0 - v}{k} \Rightarrow v_0 - v = k(x - x_0)$$

$$\text{Then } v = k(x - x_0) + v_0$$

- *2<sup>nd</sup> case :  $a = -kv^2$*

Like the first case, we first seek the velocity as a function of time.

$$a = \frac{dv}{dt} = -kv^2 \Rightarrow \frac{dv}{v^2} = -kdt$$

$$\text{So } \int_{v_0}^v \frac{dv}{v^2} = \int_0^t -kdt \Rightarrow \frac{-1}{v} + \frac{1}{v_0} = -kt \quad \text{so} \quad \frac{1}{v} = kt + \frac{1}{v_0}$$

$$\text{Then } \frac{1}{v} = \frac{ktv_0 + 1}{v_0} \quad (*)$$

$$\text{The velocity is written as: } v(t) = \frac{v_0}{ktv_0 + 1}$$

We have found the velocity, now seeking the position as a function of time.

$$v = \frac{dx}{dt} = \frac{v_0}{ktv_0 + 1} \quad \text{so} \quad dx = \frac{v_0}{ktv_0 + 1} dt$$

## Chapter III: Kinematics of material point

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$$\int_{x_0}^x dx = \int_0^t \frac{v_0}{ktv_0 + 1} dt \Rightarrow x - x_0 = \frac{1}{k} \ln(ktv_0 + 1)$$

The position as a function of time is given by  $x(t) = x_0 + \frac{1}{k} \ln(ktv_0 + 1)$

The relationship between velocity and position.

$$(*) \Rightarrow ktv_0 + 1 = \frac{v_0}{v} \text{ donc } x = x_0 + \frac{1}{k} \ln \frac{v_0}{v}$$

$$\Rightarrow \ln \frac{v_0}{v} = k(x - x_0) \Rightarrow \frac{v_0}{v} = e^{k(x-x_0)} \text{ so } v = v_0 e^{k(x_0-x)}$$

### Exercise 13

Let's find the equation of the trajectory of a body whose motion is defined by

$$v_x=1 \quad \text{and} \quad v_y=2/(t+1)$$

At (t=0) x=0 and y=2

$$\left\{ \begin{array}{l} v_x = 1 \Rightarrow \frac{dx}{dt} = 1 \\ v_y = \frac{2}{t+1} \Rightarrow \frac{dy}{dt} = \frac{2}{t+1} \end{array} \right. \text{ so } \left\{ \begin{array}{l} dx = dt \Rightarrow x = t \\ dy = \frac{2}{t+1} dt \Rightarrow \int_2^y dy = \int_0^t \frac{2}{t+1} dt \quad (*) \end{array} \right.$$

$$(*) \Rightarrow y - 2 = 2 \ln(t + 1)$$

By replacing t by x we have the equation of the trajectory of the form:

$$Y=2+2.\ln(x+1)$$

a- The components of acceleration

$$\text{b- } \left\{ \begin{array}{l} a_x = \frac{dv_x}{dt} = \frac{d(1)}{dt} = 0 \\ a_y = \frac{dv_y}{dt} = \frac{d(\frac{2}{t+1})}{dt} = \frac{-2}{(t+1)^2} \end{array} \right.$$

$$\text{So } \vec{a}(t) = \frac{-2}{(t+1)^2} \vec{j}$$

$$\text{Its modulus: } |\vec{a}| = \frac{2}{(t+1)^2}$$

### Exercise 14

A balloon rises with a horizontal velocity  $v_0$  ( $v_y=v_0$ ) and the wind gives it a horizontal velocity

$$(v_x = ay)$$

a- Let's determine the equations x(t) and y(t), with  $v_x=ay$  and  $v_y=v_0$

We take at t=0 x=0 and y=2

$$\text{We have } v_y=v_0 \Rightarrow \frac{dy}{dt} = v_0$$



## Chapter III: Kinematics of material point

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$$\Rightarrow dy = v_0 dt \quad \text{donc} \quad \int_0^y dy = \int_0^t v_0 dt$$

$$\Rightarrow y = v_0 t$$

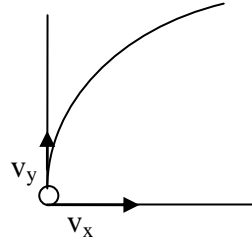
on the other hand

$$v_x = ay = av_0 t \Rightarrow \frac{dx}{dt} = av_0 t$$

$$\Rightarrow dx = av_0 t dt \quad \text{donc} \quad \int_0^x dx = \int_0^t av_0 t dt \Rightarrow x = av_0 \frac{t^2}{2}$$

$$\text{So } y = v_0 t \Rightarrow t = \frac{y}{v_0}$$

$$\Rightarrow x = av_0 \frac{\left(\frac{y}{v_0}\right)^2}{2} = \frac{a}{2v_0} y^2$$



The equation of the trajectory  $y=f(x)$  is of the form:  $y = \sqrt{\frac{2xv_0}{a}}$

b- Normal and tangential accelerations:

$$\vec{v} = v_0 \vec{i} + ay \vec{j} = v_0 \vec{i} + av_0 t \vec{j}$$

The acceleration is of the form:  $\vec{a} = \frac{d\vec{v}}{dt} = av_0 \vec{j}$

$$\text{And } |\vec{v}| = \sqrt{v_0^2 + a^2 v_0^2 t^2}$$

The tangential acceleration

$$a_T = \frac{d|\vec{v}|}{dt} = \frac{d\sqrt{v_0^2 + a^2 v_0^2 t^2}}{dt} \Rightarrow a_T = \frac{2a^2 v_0^2 t}{2\sqrt{v_0^2 + a^2 v_0^2 t^2}}$$

So the tangential acceleration is written  $a_T = \frac{a^2 v_0^2 t}{v}$

The normal acceleration

$$a^2 = a_T^2 + a_N^2 \Rightarrow a_N^2 = a^2 - a_T^2 \Rightarrow a_N^2 = a^2 v_0^2 - \frac{a^4 v_0^4 t^2}{v_0^2 + a^2 v_0^2 t^2} \Rightarrow a_N^2 = \frac{a^2 v_0^4}{v^2}$$

So the normal acceleration is written  $a_N = \frac{a v_0^2}{v}$

Radius of curvature :

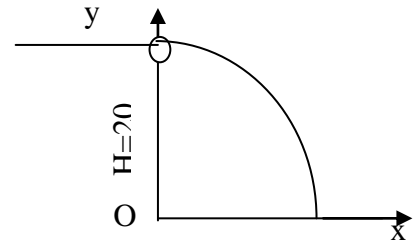
$$a_N = \frac{v^2}{R} = \frac{a v_0^2}{v} \Rightarrow R = \frac{v^3}{a v_0^2}$$

## Chapter III: Kinematics of material point

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### Exercise 15

a-A stone is thrown from the top of a building 20 meters high with a velocity  $v_x=10$  m/s



Calculating the time taken by the stone to reach the ground:

Considering oy direction, the acceleration is gravity  $g$ .

So  $a_y = -g$  (because the motion is in the opposite direction of oy, and  $a_x = 0$  since the velocity  $v_x$  in the direction of Ox is constant).

$$a_y = \frac{dv_y}{dt} = -g \Rightarrow dv_y = -gdt$$

So  $v_y = -gt$

At  $t=0$ , the stone has zero velocity in the y direction.

$$v_y = \frac{dy}{dt} = -gt \Rightarrow dy = -gtdt$$

$$\text{So } y = \frac{-1}{2}gt^2 + y_0$$

At  $t=0$ , the stone is at the top of the building, so  $y_0=H=20$ .

When the stone reaches the ground, its y component will be zero, so  $y=0$ .

Therefore, we are looking for the time at which  $y=0$ .

$$Y=-0,5 g t^2+20 =0 \Rightarrow gt^2=40$$

We take  $g=10$ m/s, so the time taken by the stone to reach the ground is  $t=2$ s.

b- Along (Ox),  $v_x = \frac{dx}{dt} = v_0 \Rightarrow dx = v_0dt$

At  $t=0$ ,  $x=0$ , so  $x=v_0t=10t$

The stone reaches the ground at a distance of 20 m.

c- Along (oy),  $v_y=-g \cdot t=-20$ m/s, and along (ox),  $v_x=10$ m/s,

$$\text{So } \vec{v} = -20\vec{i} + 10\vec{j} \Rightarrow |\vec{v}| = \sqrt{400 + 100} = \sqrt{500} \text{ m/s}$$

### Exercise 16

$$\begin{cases} v_x = R\omega\cos(\omega t) \\ v_y = R\omega\sin(\omega t) \end{cases}$$

Knowing that at  $t=0$ , the moving body is at the origin O (0,0),

## Chapter III: Kinematics of material point

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1. The components of the acceleration vector and its modulus

$$\begin{cases} a_x = \frac{dv_x}{dt} = -R\omega^2 \sin(\omega t) \\ a_y = \frac{dv_y}{dt} = R\omega^2 \cos(\omega t) \end{cases}$$

$$|\vec{a}| = \sqrt{(-R\omega^2 \sin(\omega t))^2 + (R\omega^2 \cos(\omega t))^2} = R\omega^2$$

2. The tangential and normal components of acceleration and deduce the radius of curvature.

- Tangential acceleration:

$$|\vec{v}| = \sqrt{(R\omega \cos(\omega t))^2 + (R\omega \sin(\omega t))^2} = R\omega$$

$$a_T = \frac{dv}{dt} = \frac{dR\omega}{dt} \Rightarrow a_T = 0$$

- Normal acceleration:

$$a_N = \frac{v^2}{R} = a = R\omega^2 \quad \text{and} \quad a_T = 0 \quad \text{so} \quad R = \frac{v^2}{a_N} = \frac{R^2\omega^2}{R\omega^2} = R$$

Radius of curvature is R.

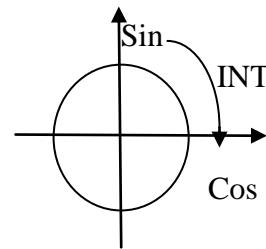
3. The components of the position vector

$$\begin{cases} v_x = R\omega \cos(\omega t) \\ v_y = R\omega \sin(\omega t) \end{cases} \Rightarrow \begin{cases} \frac{dx}{dt} = R\omega \cos(\omega t) \\ \frac{dy}{dt} = R\omega \sin(\omega t) \end{cases}$$

$$\Rightarrow \begin{cases} dx = R\omega \cos(\omega t) dt \\ dy = R\omega \sin(\omega t) dt \end{cases}$$

$$\Rightarrow \begin{cases} \int dx = R \int \omega \cos(\omega t) dt \\ \int dy = R \int \omega \sin(\omega t) dt \end{cases}$$

$$\Rightarrow \begin{cases} x = R \sin(\omega t) \\ y = -R \cos(\omega t) \end{cases}$$



The trajectory equation.

$$x^2 + y^2 = R^2 \sin^2 \omega t + R^2 \cos^2 \omega t \Rightarrow x^2 + y^2 = R^2$$

4. The nature of movement

The acceleration  $a = a_N$  and the equation of the trajectory is  $x^2 + y^2 = R^2$ , so the motion is uniformly circular.

## Chapter III: Kinematics of material point

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### Exercise 17

A material point M moves along the OX axis with acceleration  $\vec{a} = a\vec{i}$  with  $a > 0$ .

1- Determine the velocity vector knowing that  $v(t=0) = v_0$ .

$$a = \frac{dv}{dt} \Rightarrow \int_{v_0}^v dv = a \int_0^t dt$$

$$\Rightarrow v - v_0 = at(1)$$

$$\text{So } \vec{v} = (at + v_0)\vec{i}$$

2- The position vector  $\vec{OM}$  knowing that  $x(t=0) = x_0$ .

$$v = \frac{dx}{dt} = at + v_0 \Rightarrow \int_{x_0}^x dx = \int_0^t (at + v_0) dt = a \int_0^t t dt + v_0 \int_0^t dt$$

$$\Rightarrow x - x_0 = \left[ a \frac{t^2}{2} + v_0 t \right]_0^t$$

$$\int x^n dx = \frac{x^{n+1}}{n+1} \text{ and } \int \frac{dx}{x} = \ln x$$

$$x = \frac{1}{2}at^2 + v_0t + x_0(2)$$

$$\Rightarrow \vec{OM} = \left( \frac{1}{2}at^2 + v_0t + x_0 \right) \vec{i}$$

3. Show that  $v^2 - v_0^2 = 2a(x - x_0)$

$$(1) \Rightarrow t = \frac{v - v_0}{a} \text{ in } (2) \quad x - x_0 = \frac{1}{2}a \left( \frac{v - v_0}{a} \right)^2 + v_0 \left( \frac{v - v_0}{a} \right) = \frac{v^2 + v_0^2 - 2vv_0}{2a} + \frac{vv_0 - v_0^2}{a}$$

$$\Rightarrow x - x_0 = \frac{v^2 + v_0^2 - 2vv_0}{2a} + \frac{2vv_0 - 2v_0^2}{2a}$$

$$\Rightarrow x - x_0 = \frac{v^2 - v_0^2}{2a}$$

$$\text{so } 2a(x - x_0) = v^2 - v_0^2$$

4- For motion to be uniformly accelerated,  $\vec{a} \cdot \vec{v}$  must be positive.

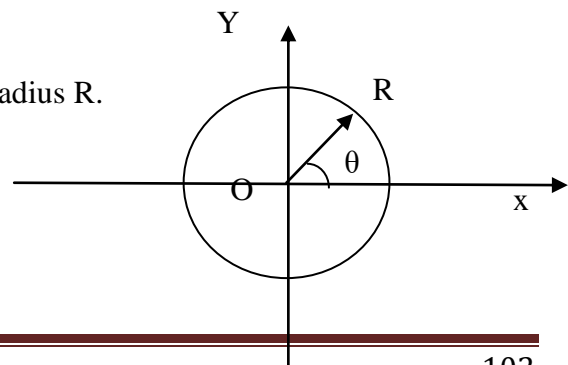
For motion to be uniformly retarded,  $\vec{a} \cdot \vec{v}$  must be negative.

### Exercise 18

a- Point M describes a circle with center O and radius R.

The coordinates of point M are:

$$\begin{cases} x = R \cos \theta \\ y = R \sin \theta \end{cases}$$



b- The velocity at point M

$$\begin{cases} v_x = \frac{dx}{dt} = -R \frac{d\theta}{dt} \sin \theta = -R\dot{\theta} \sin \theta \\ v_y = R \frac{d\theta}{dt} \cos \theta = R\dot{\theta} \cos \theta \end{cases}$$

$$v = \sqrt{v_x^2 + v_y^2} = \sqrt{R^2 \dot{\theta}^2 (-\sin^2 \theta) + R^2 \dot{\theta}^2 (\cos^2 \theta)}$$

$$\Rightarrow v = R\dot{\theta}$$

c- The components of acceleration

$$\begin{cases} a_x = \frac{dv_x}{dt} = -R\ddot{\theta} \sin \theta - R\dot{\theta}^2 \cos \theta \\ a_y = \frac{dv_y}{dt} = R\ddot{\theta} \cos \theta - R\dot{\theta}^2 \sin \theta \end{cases}$$

$$a = \sqrt{R^2 \ddot{\theta}^2 + R^2 \dot{\theta}^4}$$

The tangential acceleration:

$$a_T = \frac{dv}{dt} = \frac{dR\dot{\theta}}{dt} \Rightarrow a_T = R\ddot{\theta}$$

The normal acceleration:

$$a_N = \frac{v^2}{R} = R\dot{\theta}^2$$

d- We have  $\alpha = \frac{d\omega}{dt} \Rightarrow d\omega = \alpha dt$  so  $\omega - \omega_0 = \alpha t \Rightarrow \omega = \omega_0 + \alpha t$  (\*)

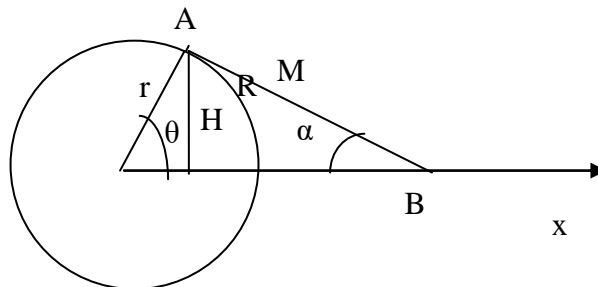
At (t=0),  $\theta_0=0$  and  $\omega=\omega_0$  with  $\omega = \frac{d\theta}{dt} = \omega_0 + \alpha t \Rightarrow d\theta = \omega_0 dt + \alpha t dt$

So  $\int_0^\theta d\theta = \int_0^t \omega_0 dt + \int_0^t \alpha t dt$  Then  $\theta = \alpha \frac{t^2}{2} + \omega_0 t$

$$(*) \Rightarrow t = \frac{\omega - \omega_0}{\alpha} \Rightarrow \theta = \frac{\alpha}{2} \left( \frac{\omega - \omega_0}{\alpha} \right)^2 + \omega_0 \left( \frac{\omega - \omega_0}{\alpha} \right) = \frac{\omega^2 - \omega_0^2}{2\alpha}$$

$$\Rightarrow \theta = \frac{\omega^2 - \omega_0^2}{2\alpha}$$

### Exercise 19



a- The time equations of point M (the middle of AB).

## Chapter III: Kinematics of material point

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With  $OA=r$  and  $OA=R$

Along the Ox axis :

$$x_M = r \cos\theta + \frac{R}{2} \cos\alpha$$

$$AH = r \sin\alpha = R \sin\theta \Rightarrow \sin\alpha = \frac{r \sin\theta}{R}$$

$$\cos^2\alpha + \sin^2\alpha = 1 \Rightarrow \cos^2\alpha = 1 - \sin^2\alpha = 1 - \frac{r^2}{R^2} \sin^2\theta$$

$$\text{So } \cos\alpha = \sqrt{1 - \frac{r^2}{R^2} \sin^2\theta}$$

$$\theta = \omega t \text{ then } x_M = r \cos\omega t + \frac{R}{2} \sqrt{1 - \frac{r^2}{R^2} \sin^2\omega t}$$

$$y_M = \frac{R}{2} \sin\alpha = \frac{Rr \sin\theta}{2R} = \frac{r}{2} \sin\omega t \quad \text{so } y_M = \frac{r}{2} \sin\omega t$$

$$\begin{cases} x_M = r \cos\omega t + \frac{R}{2} \sqrt{1 - \frac{r^2}{R^2} \sin^2\omega t} \\ y_M = \frac{r}{2} \sin\omega t \end{cases}$$

b- The abscissa of B:

$$x_B = r \cos\theta + R \cos\alpha = r \cos\omega t + R \sqrt{1 - \frac{r^2}{R^2} \sin^2\omega t}$$

According to the equation of motion, point B does not have sinusoidal motion

c- The velocity at point M

$$v_M = \frac{d\vec{OM}}{dt} \Rightarrow \begin{cases} v_x = \frac{dx}{dt} = -r(\omega \sin\omega t) - R \left( \frac{r^2 \omega \sin\omega t \cos\omega t}{\sqrt{R^2 - r^2 \sin^2\omega t}} \right) \\ v_y = \frac{dy}{dt} = \frac{r\omega}{2} \cos\omega t \end{cases}$$

d- If  $r=R$  Showing that the movement of B is sinusoidal

$$\begin{cases} x_B = r \cos\omega t + r \sqrt{1 - \frac{r^2}{r^2} \sin^2\omega t} = 2r \cos\omega t \\ y_B = 0 \end{cases}$$

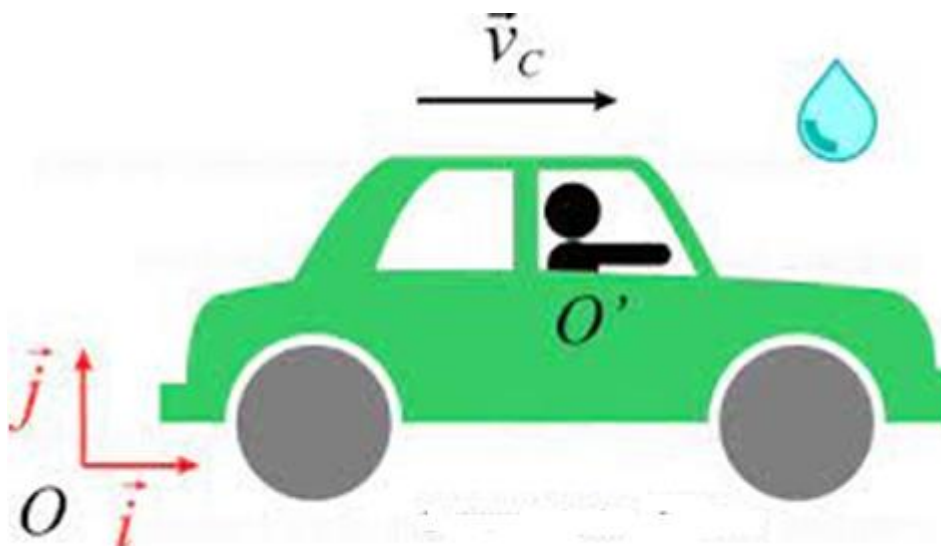
$$\Rightarrow x_B = 2r \cos\omega t$$

So the movement of B in this case is sinusoidal.

**COURSE OF MECHANICS**  
**OF THE MATERIAL POINT**

*Chapter IV: Relative motion*

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**Glossary**

<b>In English</b>	<b>In French</b>	<b>In Arabic</b>
Relative motion	Mouvement relative	حركة نسبية
Absolute reference frame	Un référentiel absolu	مرجع مطلق
Relative reference frame	Un référentiel relatif	مرجع نسبي
Absolute motion	Mouvement absolu	حركة مطلقة
Entrainment motion	Mouvement d'entraînement	حركة المعلم المتحرك بالنسبة للمعلم الثابت
Velocity composition	Composition de la vitesse	تركيب السرعة
Fixed frame of reference	Un référentiel fixe	المعلم الثابت
Moving frame of reference	Un référentiel mobile	المعلم المتحرك
Absolute velocity	Vitesse absolue	السرعة المطلقة
Relative velocity	Vitesse relative	السرعة النسبية
Entrainment velocity	Vitesse d'entraînement	سرعة المعلم المتحرك بالنسبة للمعلم الثابت
Composition of acceleration	Composition des accelerations	تركيب التسارع
Absolute acceleration	Acceleration absolue	التسارع المطلق
Relative acceleration	Acceleration relative	التسارع النسبي
Entrainment acceleration	Acceleration d'entraînement	التسارع المعلم المتحرك بالنسبة للمعلم الثابت
Coriolis acceleration	Acceleration de Coriolis	التسارع Coriolis



### 1. Introduction

State of motion or state of rest are two essentially relative notions, meaning that each of the two states depends on the position of the moving body relative to the body taken as reference frame. All the motions we have studied so far have been in a Galilean frame of reference, i.e. at rest or in uniform rectilinear motion. When two observers linked to two different reference frames are in motion relative to each other, the position, trajectory, velocity and acceleration of the same moving body vary according to the reference frame chosen by the observer.

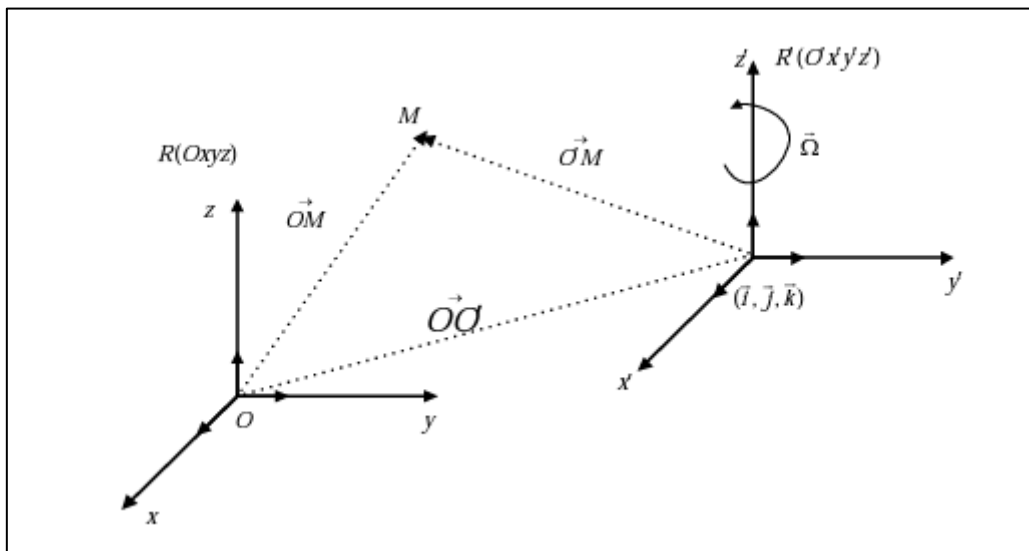
A bus passenger, for example, is in motion relative to an observer seated on the side of the road, whereas he is at rest relative to another observer (a passenger lending the same bus). Clearly, then, the notion of motion or rest is intimately linked to the position of the observer.

To say observer is to say to choose a frame of reference to determine the position, velocity and acceleration of a moving object at each instant.

### 2. Composition of movements مركبات الحركة

Let  $R(O, x, y, z)$  be a fixed or absolute (Galilean) frame of reference and  $R'(O', x', y', z')$  a moving or relative (non-Galilean) frame of reference relative to  $(R)$ .

It's always useful to know how to determine the position, velocity and acceleration of a material point  $M$  in a fixed reference frame if they are known in the other relative reference frame and vice-versa.



Let's associate to the reference frame  $R$  (called **absolute reference frame**) the reference frame  $R(O, \vec{i}, \vec{j}, \vec{k})$  and to the reference frame  $R'$  (called **relative reference frame**) the reference frame  $R'(O', \vec{i}', \vec{j}', \vec{k}')$ .

## Chapter IV: Relative motion

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If M is a movable point in space, defined by coordinates (x, y, z) in the fixed reference frame (R) and by (x', y', z') in the movable reference frame (R').

### We will call:

Relative motion الحركة النسبية : the motion of M relative to (R').

Absolute motion الحركة المطلقة : the motion of M relative to (R).

Entrainment motion: the motion of the moving frame of reference (R') relative to the fixed frame of reference (R).

### 2.1. Velocity composition قانون تركيب السرعات

If we know the relative motion, i.e. the motion of the material point considered in the moving frame of reference (relative frame of reference المعلم المتحرك) and that of the moving frame of reference relative to the fixed frame of reference (fixed frame of reference المعلم الثابت).

We have:  $\overrightarrow{OM} = \overrightarrow{OO'} + \overrightarrow{O'M} = \overrightarrow{OO'} + (x'\vec{i}' + y'\vec{j}' + z'\vec{k}')$

With:  $\overrightarrow{OM} = (x\vec{i} + y\vec{j} + z\vec{k})$  in the fixed reference frame R.

and  $\overrightarrow{O'M} = (x'\vec{i}' + y'\vec{j}' + z'\vec{k}')$  in the moving reference frame R'.

The velocity is then:  $\vec{v} = \frac{d\overrightarrow{OM}}{dt} = \frac{dx}{dt}\vec{i} + \frac{dy}{dt}\vec{j} + \frac{dz}{dt}\vec{k}$

$$\frac{d\overrightarrow{OM}}{dt} = \frac{d\overrightarrow{OO'}}{dt} + \frac{d\overrightarrow{O'M}}{dt}$$

$$= \frac{dx}{dt}\vec{i} + \frac{dy}{dt}\vec{j} + \frac{dz}{dt}\vec{k} + \frac{d\overrightarrow{OO'}}{dt} + x'\frac{d\vec{i}'}{dt} + y'\frac{d\vec{j}'}{dt} + z'\frac{d\vec{k}'}{dt} = \vec{v}_a = \vec{v}_r + \vec{v}_e$$

$$\text{with } \begin{cases} \vec{v}_a = \frac{d\overrightarrow{OM}}{dt} = \frac{dx}{dt}\vec{i} + \frac{dy}{dt}\vec{j} + \frac{dz}{dt}\vec{k} = \overrightarrow{v(M)} / (R) \\ \vec{v}_r = \frac{dx'}{dt}\vec{i}' + \frac{dy'}{dt}\vec{j}' + \frac{dz'}{dt}\vec{k}' = \frac{d\overrightarrow{O'M}}{dt} / (R') = \overrightarrow{v(M)} / (R') \\ \vec{v}_e = \frac{d\overrightarrow{OO'}}{dt} + x'\frac{d\vec{i}'}{dt} + y'\frac{d\vec{j}'}{dt} + z'\frac{d\vec{k}'}{dt} = \overrightarrow{v(R')} / (R) \end{cases}$$

$\vec{v}_a$  represents absolute velocity السرعة المطلقة, the derivative of  $\overrightarrow{OM}$  with respect to time in the fixed reference frame.

$\vec{v}_r$  is the relative velocity السرعة النسبية, i.e. the derivative of  $\overrightarrow{O'M}$  with respect to time in the moving reference frame.

## Chapter IV: Relative motion

$\vec{v}_e$  is entrainment velocity سرعة المعلم المتحرك بالنسبة للمعلم الثابت, this is the derivative of  $\overline{OM}$  with respect to time in the fixed reference frame, considering the moving point M fixed in the moving reference frame ( $x', y'$  and  $z'$  are constant).

It also represents the velocity of the moving frame of reference relative to the fixed frame.

The law of velocity composition is given by:

$$\vec{v}_a = \vec{v}_r + \vec{v}_e$$

### Note :

1. When the entrainment motion is translational, the vectors  $(\vec{i}', \vec{j}' \text{ and } \vec{k}')$  remain parallel to the unit vectors  $(\vec{i}, \vec{j} \text{ and } \vec{k})$  of the fixed reference frame.

So  $\frac{d\vec{i}'}{dt} = \frac{d\vec{j}'}{dt} = \frac{d\vec{k}'}{dt} = \vec{0}$  and  $\vec{v}_e = \left(\frac{d\overline{OO'}}{dt}\right) / R$

2. If the moving frame of reference (R') rotates relative to the fixed frame of reference (R), the entrainment velocity can also be written as:

$$\vec{v}_e = \frac{d\overline{OO'}}{dt} + \vec{\omega} \wedge \overline{O'M}$$

With  $\vec{\omega}$  represents the rotational velocity السرعة الزاوية of R'/R.

3. In the case of translational motion  $\vec{\omega} = \vec{0}$  then  $\vec{v}_e = \frac{d\overline{OO'}}{dt}$ .

### **2.2. Composition of acceleration** قانون تركيب التسارعات

Acceleration is the derivative of velocity with respect to time :

$$\vec{a} = \frac{d\vec{v}}{dt} = \frac{d^2\overline{OM}}{dt^2} = \frac{d^2x}{dt^2}\vec{i} + \frac{d^2y}{dt^2}\vec{j} + \frac{d^2z}{dt^2}\vec{k}$$

$$\vec{v} = \frac{d\overline{OO'}}{dt} + \frac{dx'}{dt}\vec{i}' + \frac{dy'}{dt}\vec{j}' + \frac{dz'}{dt}\vec{k}' + x'\frac{d\vec{i}'}{dt} + y'\frac{d\vec{j}'}{dt} + z'\frac{d\vec{k}'}{dt}$$

$$\frac{d}{dt}\left(\frac{dx'}{dt}\vec{i}'\right) = \frac{d^2x'}{dt^2}\vec{i}' + \frac{dx'}{dt}\frac{d\vec{i}'}{dt}$$

$$\frac{d}{dt}\left(\frac{dy'}{dt}\vec{j}'\right) = \frac{d^2y'}{dt^2}\vec{j}' + \frac{dy'}{dt}\frac{d\vec{j}'}{dt}$$

$$\frac{d}{dt}\left(\frac{dz'}{dt}\vec{k}'\right) = \frac{d^2z'}{dt^2}\vec{k}' + \frac{dz'}{dt}\frac{d\vec{k}'}{dt}$$

and

$$\frac{d}{dt}\left(x'\frac{d\vec{i}'}{dt}\right) = \frac{dx'}{dt}\frac{d\vec{i}'}{dt} + x'\frac{d^2\vec{i}'}{dt^2}$$

$$\frac{d}{dt} \left( y' \frac{d\vec{j}'}{dt} \right) = \frac{dy'}{dt} \frac{d\vec{j}'}{dt} + y' \frac{d^2\vec{j}'}{dt^2}$$

$$\frac{d}{dt} \left( z' \frac{d\vec{k}'}{dt} \right) = \frac{dz'}{dt} \frac{d\vec{k}'}{dt} + z' \frac{d^2\vec{k}'}{dt^2}$$

$$\begin{aligned} \vec{a} = & \frac{d^2x'}{dt^2} \vec{i}' + \frac{d^2y'}{dt^2} \vec{j}' + \frac{d^2z'}{dt^2} \vec{k}' + \frac{d^2\vec{OO}'}{dt^2} + x' \frac{d^2\vec{i}'}{dt^2} + y' \frac{d^2\vec{j}'}{dt^2} + z' \frac{d^2\vec{k}'}{dt^2} + 2 \left( \frac{dx'}{dt} \frac{d\vec{i}'}{dt} + \frac{dy'}{dt} \frac{d\vec{j}'}{dt} \right. \\ & \left. + \frac{dz'}{dt} \frac{d\vec{k}'}{dt} \right) \end{aligned}$$

with

$$\vec{a}_a = \frac{d^2\vec{OM}}{dt^2} = \frac{d^2x}{dt^2} \vec{i} + \frac{d^2y}{dt^2} \vec{j} + \frac{d^2z}{dt^2} \vec{k}$$

$$\vec{a}_r = \frac{d^2\vec{O'M}}{dt^2} = \frac{d^2x'}{dt^2} \vec{i}' + \frac{d^2y'}{dt^2} \vec{j}' + \frac{d^2z'}{dt^2} \vec{k}'$$

$$\vec{a}_e = \frac{d^2\vec{OO}'}{dt^2} + x' \frac{d^2\vec{i}'}{dt^2} + y' \frac{d^2\vec{j}'}{dt^2} + z' \frac{d^2\vec{k}'}{dt^2}$$

$$\Rightarrow \vec{a}_e = \frac{d^2\vec{OO}'}{dt^2} + \vec{\omega} \wedge (\vec{\omega} \wedge \vec{O'M}) + \frac{d\vec{\omega}}{dt} \wedge \vec{O'M}$$

$$\vec{a}_c = 2 \left( \frac{dx'}{dt} \frac{d\vec{i}'}{dt} + \frac{dy'}{dt} \frac{d\vec{j}'}{dt} + \frac{dz'}{dt} \frac{d\vec{k}'}{dt} \right)$$

$$\Rightarrow \vec{a}_c = 2(\vec{\omega} \wedge \vec{v}_r) / R'$$

Then the law of composition of the acceleration will be :

$$\vec{a}_a = \vec{a}_r + \vec{a}_c + \vec{a}_e$$

$\vec{a}_a$  is the absolute acceleration التسارع المطلق representing the second derivative of  $\vec{OM}$  with respect to time in the fixed reference frame. This is the acceleration of M in the fixed frame of reference.

$\vec{a}_r$  is the relative acceleration التسارع النسبي representing the second derivative of  $\vec{O'M}$  with respect to time in the moving frame of reference. This is the acceleration of M in the moving frame of reference.

$\vec{a}_e$  is the entrainment acceleration تسارع المعلم المتحرك بالنسبة للمعلم الثابت, which represents the acceleration of the motion of the moving frame relative to the fixed frame. This is the acceleration of the moving frame R' relative to the fixed frame R.

$\vec{a}_c$  is the Coriolis or complementary acceleration (it has no physical meaning).

## Chapter IV: Relative motion

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### Note:

If the moving frame of reference (R') rotates relative to the fixed frame of reference (R), the drive speed can also be written as:

$$\vec{v}_e = \left(\frac{d\overline{OO'}}{dt}\right)/R + (\vec{\omega}\Lambda\overline{O'M})/R',$$

Entrainment acceleration by:

$$\vec{a}_e = \left(\frac{d^2\overline{OO'}}{dt^2}\right)/R + \left(\frac{d\vec{\omega}}{dt}\Lambda\overline{O'M}\right)/R' + (\vec{\omega}\Lambda\vec{\omega}\Lambda\overline{O'M})/R'$$

and the Coriolis acceleration by:

$$\vec{a}_c = 2. (\vec{\omega}\Lambda\vec{v}_r)/R'$$

With  $\vec{\omega}$  represents the rotational velocity of R'/R.

### Particular cases:

1. When the training motion is translational, the vectors  $\vec{i}'$ ,  $\vec{j}'$  and  $\vec{k}'$  remain parallel to the unit vectors ( $\vec{i}$ ,  $\vec{j}$  and  $\vec{k}$ ) of the fixed reference frame.

$$\text{So } \frac{d\vec{i}'}{dt} = \frac{d\vec{j}'}{dt} = \frac{d\vec{k}'}{dt} = \vec{0} \text{ and } \vec{\omega} = \vec{0}$$

$$\text{Then } \vec{v}_e = \frac{d\overline{OO'}}{dt}, \vec{a}_e = \left(\frac{d^2\overline{OO'}}{dt^2}\right) \text{ and } \vec{a}_c = \vec{0}.$$

2. If R' has a uniform rectilinear motion then  $\vec{\omega} = \vec{0}$ ,  $\vec{a}_c = \vec{0}$ ,  $\frac{d\overline{OO'}}{dt} = cst$  and

$$\frac{d^2\overline{OO'}}{dt^2} = \vec{0} \text{ So } \vec{a}_e = \vec{0} \text{ because } \vec{v}_e = cst.$$

3. If R' has a pure rotation about R (R' and R have the same origin), so we have:

$$\frac{d^2\overline{OO'}}{dt^2} = \vec{0} \text{ and } \frac{d\overline{OO'}}{dt} = \vec{0} \text{ then } \vec{v}_e = (\vec{\omega}\Lambda\overline{O'M})/R'$$

$$\text{And } \vec{a}_e = \left(\frac{d\vec{\omega}}{dt}\Lambda\overline{O'M}\right)/R' + (\vec{\omega}\Lambda\vec{\omega}\Lambda\overline{O'M})/R'$$

### Proposed exercises about chapter IV

#### Exercise 1

A moving train of a speed  $v$  passes through a station. At the moment  $t=0$  a M bulb is detached from the ceiling of one of its compartments. The movement of M is then observed by a passenger of the train and by the station chief motionless on the platform.

Describe the movement of M for each observer.

#### Exercise 2

A man mounting a horse galloping at a constant speed  $v$ , launches an arrow into the air with speed  $v_0$  relative to the horse.

At what angle  $\theta$  should  $v_0$  make with the vertical for the arrow to fall back onto the man?

(air resistance will be neglected).

#### Exercise 3

The coordinates of a moving particle in the reference frame (R) provided with the reference frame  $(O, \vec{i}, \vec{j}, \vec{k})$  are given as a function of time by:

$$x = 2t^3 + 1, \quad y = 4t^2 + t - 1, \quad z = t^2$$

In a second frame of reference (R') with the reference frame  $(O', \vec{i}', \vec{j}', \vec{k}')$  with  $\vec{i} = \vec{i}'$ ,  $\vec{j} = \vec{j}'$ ,  $\vec{k} = \vec{k}'$  the coordinates of a moving particle are given by:

$$x' = 2t^3, \quad y' = 4t^2 - 3t + 2, \quad z' = t^2 - 5$$

1- Calculate the express the velocity  $v$  of m in (R) as a function of its velocity  $v'$  in (R'), and proceed in the same way for the accelerations.

2- Define the entrainment motion of (R') relative to (R).

#### Exercise 4

In the (Oxy) plane, consider a system of moving axes (OXY) with the **same origin O**, rotating with a constant angular velocity  $\omega$  around (OZ). A moving point M moves along axis

## Chapter IV: Relative motion

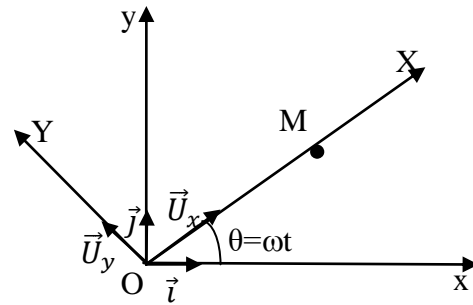
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(OX) with **constant acceleration**  $\gamma$  and no initial velocity. We call relative motion of M its motion with respect to (OXY), and absolute motion with respect to (Oxy).

At time  $t=0$ , axes (Ox) and (OX) are coincident and M is in OA.

Calculate in the moving reference frame :

- 1- The velocity and relative acceleration of M.
- 2- Entrainment velocity and acceleration.
- 3- Coriolis acceleration.
- 4- Deduce its absolute velocity and acceleration.



### Exercise 5

A point M moves with a constant **velocity**  $v_0$  **on the axis (OX)** of a coordinate system (OXYZ) which rotates with a constant angular velocity  $\omega$  around (Oz) in the plane (Oxy).

- 1- What is the expression of  $\overrightarrow{OM}$  in the fixed frame (Oxy)? Calculate the absolute velocity and absolute acceleration.
- 2- Calculate the relative velocity  $\overrightarrow{v_r}$  and the entrainment velocity  $\overrightarrow{v_e}$ , verify that  $\overrightarrow{v_a} = \overrightarrow{v_r} + \overrightarrow{v_e}$ .
- 3- Calculate the relative acceleration  $\overrightarrow{a_r}$ , the entrainment acceleration  $\overrightarrow{a_e}$ , and the coriolis acceleration  $\overrightarrow{a_c}$ , verify that  $\overrightarrow{a_a} = \overrightarrow{a_r} + \overrightarrow{a_e} + \overrightarrow{a_c}$ .

### Exercise 6

In the (Oxy) plane, we consider a system of moving axes (OXY) with the same **origin o** and such that (OX) makes a variable angle  $\theta$  with (Ox). A point M moving **along axis (OX)** is marked by  $\overrightarrow{OM} = \mathbf{r}$ . We call relative motion of M, its motion with respect to (OXY), and absolute motion with respect to (Oxy).

Calculate in the moving frame of reference (**polar coordinates**):

- 1- Relative velocity and acceleration of M.
- 2- The velocity and entrainment acceleration of M.
- 3- Coriolis acceleration.
- 4- Deduce its absolute velocity and acceleration.

### Exercise 7

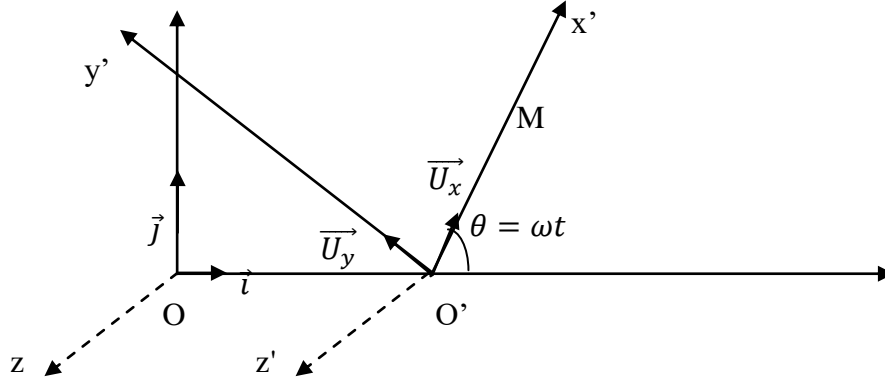
Consider the reference frame R(Oxyz) where point O' moves **along the axis (Ox)** with **constant velocity v**. O' is linked to the reference frame (O'x'y'z'), which rotates **around (oz)**

## Chapter IV: Relative motion

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with **constant angular velocity**  $\omega$ . A moving point M **moves along the axis O'x'** such that  $|\mathbf{O'M}|=t^2$ .

At time  $t=0$ , the axes (Ox) and (O'x') are coincident and M is at O.



1. Calculate the relative velocity  $\vec{v}_r$  and the entrainment velocity  $\vec{v}_e$ , deduce the absolute velocity  $\vec{v}_a$ .
2. Calculate the relative acceleration  $\vec{a}_r$ , the entrainment acceleration  $\vec{a}_e$  and the Coriolis acceleration  $\vec{a}_c$ , deduce the absolute acceleration  $\vec{a}_a$ .

### Exercise 8

Consider the reference frame R(Oxyz) where point **O' moves along axis (Oy)** with constant **acceleration**  $\gamma$ . We link to O' the reference frame (O'XYZ) which **rotates around (Oz)** with a **constant angular velocity**  $\omega$ . The coordinates of a moving body M in the moving frame of reference **are**  $x'=t^2$  and  $y'=t$ .

At time  $t=0$ , the axis (O'x) coincides with (Ox).

Calculate in the moving frame of reference:

- 1- Velocity  $\vec{v}_r$  and  $\vec{v}_e$ , deduce the absolute velocity  $\vec{v}_a$ .
- 2- Relative acceleration  $\vec{a}_r$ , entrainment acceleration  $\vec{a}_e$  and Coriolis acceleration  $\vec{a}_c$ , deduce the absolute acceleration  $\vec{a}_a$ .

### Exercise 9

In a plane (Oxy), consider a system of moving axes (OXY), of the **same origin O**, rotates around (Oz) with a constant angular velocity  $\omega$ . A moving point on axis (OX) is marked by:

$$|\overline{OM}| = r = r_0(1 + \sin\omega t)$$

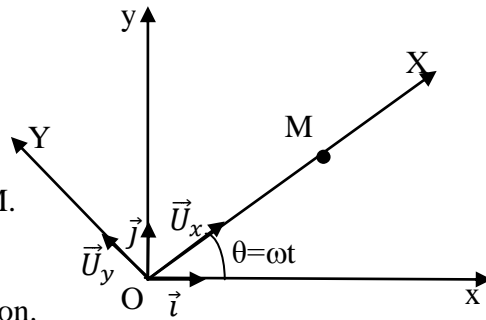


## Chapter IV: Relative motion

Where  $r_0$  is a positive constant:

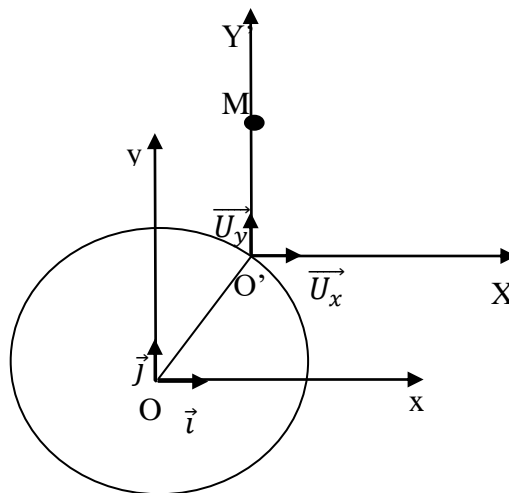
Calculate in the moving reference frame:

- 1) Relative velocity and acceleration of M.
- 2) Entrainment velocity and acceleration of M.
- 3) Coriolis acceleration.
- 4) Deduce its absolute velocity and acceleration.



### Exercise 10

In the plane (Oxy) of a reference frame (Oxyz), a point O', to which the reference frame (O'XYZ) is linked, describes a circle of **center O and radius R**, and rotates **with a constant angular velocity  $\omega$** . A point M moves **along the axis (O'Y) parallel to Oy** with constant **acceleration  $\gamma$**  (at time  $t=0$ , M is merged with  $M_0(R,0,0)$  and its initial velocity is positive).



- 1- Calculate in the (Oxyz) reference frame the position vector  $\overrightarrow{OM}$ , the absolute velocity  $\overrightarrow{v_a}$ .  
And the absolute acceleration  $\overrightarrow{a_a}$ .
- 2- Knowing that  $O'X' // O x$ ,  $O'Y' // O y$  and  $O'Z' // O z$ , calculate:
  - a- Relative speed and drive velocity, check that  $\overrightarrow{v_a} = \overrightarrow{v_r} + \overrightarrow{v_e} + \overrightarrow{v_c}$ .
  - b- A relative acceleration  $\overrightarrow{a_r}$ , entrainment acceleration  $\overrightarrow{a_e}$ , and the Coriolis acceleration  $\overrightarrow{a_c}$ , check that  $\overrightarrow{a_a} = \overrightarrow{a_r} + \overrightarrow{a_e} + \overrightarrow{a_c}$ .

### Exercise 11

Consider a fixed reference frame (Oxyz) and a moving reference frame (Ox'y'z') which rotates **around (Oz)** with a **constant angular velocity  $\omega$** .

## Chapter IV: Relative motion

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A moving point M ( $\mathbf{OM}=\mathbf{r}$ ) moves **along the axis (Ox')** according to the law:

$$\mathbf{r} = r_0 (\cos \omega t + \sin \omega t)$$
 with  $r_0 = \text{constant}$ .

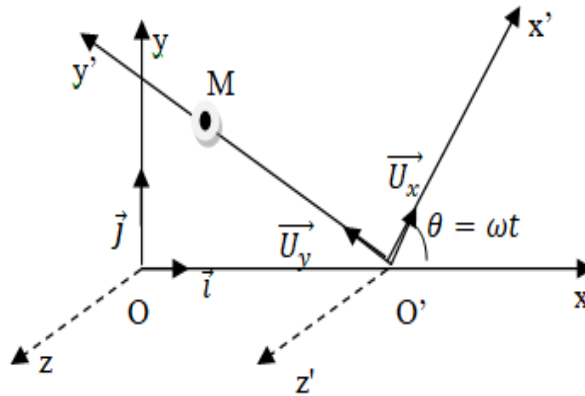
Determine in the **moving reference frame (Ox'y'z')**:

- 1- The velocity  $\vec{v}_r$  and the entrainment velocity  $\vec{v}_e$ , deduce the absolute velocity  $\vec{v}_a$ .
- 2- Relative acceleration  $\vec{a}_r$ , entrainment acceleration  $\vec{a}_e$  and Coriolis acceleration  $\vec{a}_c$ , deduce absolute acceleration  $\vec{a}_a$ .

### Exercise 12

Consider the fixed reference frame R(Oxyz) where point O' moves **along axis (Ox)** with **constant velocity  $\mathbf{v}_0$** . Linked to O' is the moving reference frame (O'x'y'z') which rotates **around (Oz)** with **constant angular velocity  $\omega$** . A moving point M moves **along the (O'y')** axis with constant **acceleration  $\gamma$** .

At time  $t=0$ , the axes (Ox) and (O'x') are coincident and M is at O.



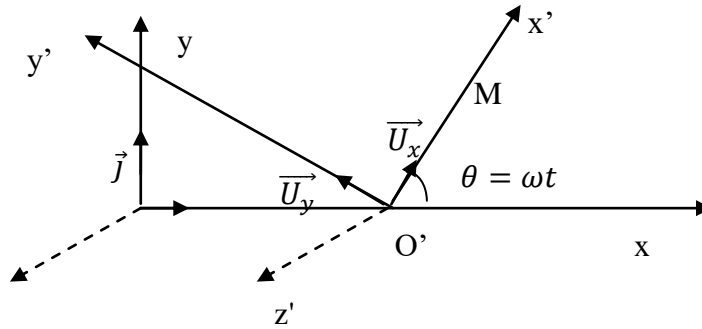
Calculate in the moving frame:

- 1- The relative velocity  $\vec{v}_r$  and the entrainment velocity  $\vec{v}_e$ , deduce the absolute velocity  $\vec{v}_a$ .
- 2- The relative acceleration  $\vec{a}_r$ , the entrainment acceleration  $\vec{a}_e$  and Coriolis acceleration  $\vec{a}_c$ , deduce the absolute acceleration  $\vec{a}_a$ .

### Exercise 13

In the (Oxy) plane, a point O' (the origin of the moving reference frame) moves along the **(Ox) axis** such that  $|\mathbf{OO}'|=t$ . The reference frame (O'X'Y') rotates **around (Oz)** with a constant angular velocity  $\omega$ . A moving point M ( $\mathbf{O'M}=\mathbf{r}$ ) moves along the **axis (O'X')** according to the law:  $\mathbf{r} = r_0 (\cos \omega t + \sin \omega t)$  With;  $r_0 = \text{constant}$ .

## Chapter IV: Relative motion

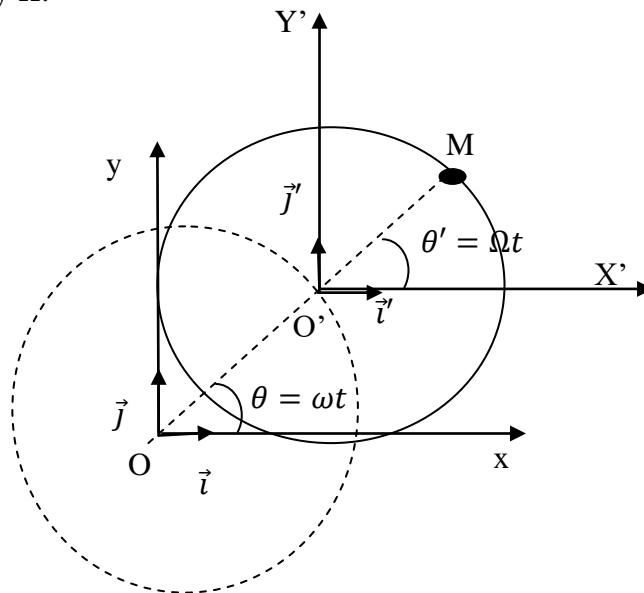


Determine at time  $t$  as a function of  $r_0$  and  $\omega$ :

- 1- The relative velocity  $\vec{v}_r$  and the entrainment velocity  $\vec{v}_e$ , deduce the absolute velocity  $\vec{v}_a$  in the moving frame of reference.
- 2- The relative acceleration  $\vec{a}_r$ , the entrainment acceleration  $\vec{a}_e$  and Coriolis acceleration  $\vec{a}_c$  in the moving frame of reference, deduce the absolute acceleration  $\vec{a}_a$  in this frame of reference.

### Exercise 14

In the frame  $(Oxyz)$ , a point  $O'$  to which the frame  $(O'XYZ)$  is attached, describes a circle with center  $O$  and radius  $R$ ; it rotates with a constant angular velocity  $\omega'$  in the plane  $(Oxy)$ . A point  $M$  describes a circle with center  $O'$  and radius  $d$  in the plane  $(O'XY)$ ; it rotates with a constant angular velocity  $\Omega$ .



Calculate in the fixed frame  $(Oxyz)$  knowing that  $(O'X') \parallel (OX)$ :

- 1- The absolute velocity  $\vec{v}_a$ , relative velocity  $\vec{v}_r$ , and entrainment velocity  $\vec{v}_e$ . Verify that  $\vec{v}_a = \vec{v}_r + \vec{v}_e$ .

## Chapter IV: Relative motion

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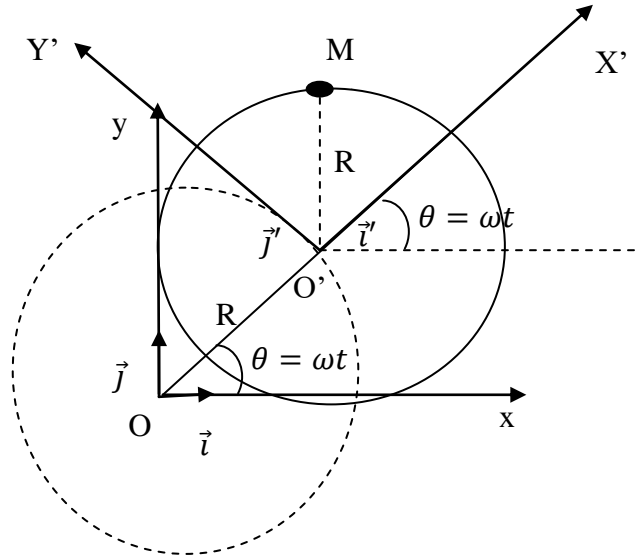
2- The absolute acceleration  $\vec{a}_a$ , relative acceleration  $\vec{a}_r$ , entrainment acceleration  $\vec{a}_e$ , and Coriolis acceleration  $\vec{a}_c$ . Verify that  $\vec{a}_a = \vec{a}_r + \vec{a}_e + \vec{a}_c$ .

### Exercise 15

In the plane (Oxy), a point O' to which the frame (O'x'y'z') is attached, describes a circle with diameter **R** rotates at a constant angular velocity  $\omega$  around point O. A point M initially at O' moves along the circumference in the positive direction with the same angular velocity  $\omega$  and the same radius **R**.

At time t=0, M is on the O'x' axis parallel to O'x':

1. Provide the expression for  $\vec{OM}$  in the fixed reference frame R(Oxy).
2. Calculate the components of the velocity and acceleration vectors of M in the reference frame R(Oxy).
3. Calculate the components of the velocity and acceleration vectors of M in the reference frame R'(O'x'y').
4. Calculate the entrainment velocity  $\vec{v}_e$ , the entrainment acceleration  $\vec{a}_e$ , and the Coriolis acceleration  $\vec{a}_c$ .

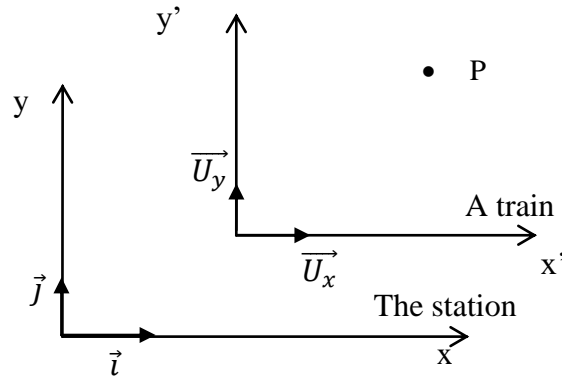


**Correction of exercises about chapter IV**

**Exercise 1**

We consider the station as the fixed reference frame and the train as the moving reference frame. The motion of “P” observed by a passenger (inside the train)

$$\vec{v}_p = -gt\vec{j} = \vec{v}_r$$



The speed of the train relative to the station  $\vec{v} = \vec{v}_e = v\vec{i}$  with  $\vec{U}_x = \vec{i}$  and  $\vec{U}_y = \vec{j}$

The reference frames are parallel because the train undergoes a translational motion relative to the station.

$$v = \frac{dx}{dt} \Rightarrow dx = vdt \quad \text{so } x = vt \quad (\text{at } t=0, x=0)$$

$$x = vt \Rightarrow t = \frac{x}{v}$$

The motion of “P” observed by the stationary stationmaster on the platform:

$$\vec{v}_a = \vec{v}_r + \vec{v}_e = -gt\vec{j} + v\vec{i} \Rightarrow \vec{v}_a = -g\left(\frac{x}{v}\right)\vec{j} + v\vec{i}$$

The position :

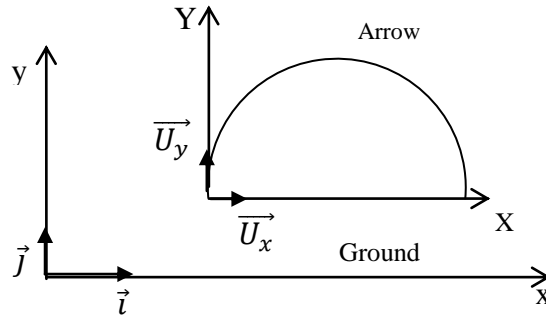
$$\begin{cases} x = vt \\ y = \frac{1}{2}gt^2 + y_0 \end{cases}$$

$$\text{Or } y = \frac{1}{2}g\frac{x^2}{v^2} + Y_0$$

## Chapter IV: Relative motion

### Exercise 2

In this case, the ground is the fixed reference frame, and the horse is the moving reference frame. The horse has a translational motion relative to the ground, so the unit vectors of the two frames are equal.

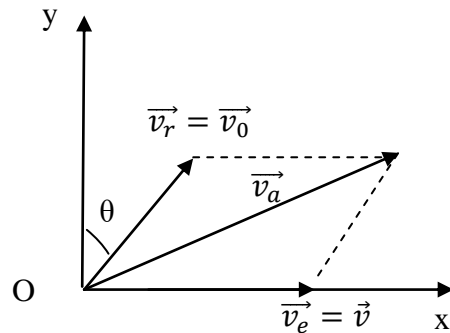


We study the motion of the arrow in both frames.

$$\overrightarrow{v_{Horse}/ground} = \overrightarrow{v_e} = \vec{v} \quad \text{and} \quad \overrightarrow{v_{Arrow}/horse} = \overrightarrow{v_r} = \overrightarrow{v_0}$$

$$\overrightarrow{v_a} = \overrightarrow{v_r} + \overrightarrow{v_e} = v_{ax}\vec{i} + v_{ay}\vec{j} \Rightarrow \begin{cases} v_{ax} = v + v_0 \sin \theta \\ v_{ay} = v_0 \cos \theta \end{cases} (*)$$

The entrainment velocity is along the ox axis (translation motion), whereas the relative velocity makes an angle  $\theta$  with the vertical oy.



$$(*) \Rightarrow \begin{cases} x_{Arrow/ground} = \int_0^r v_{ax} dt = \int_0^r (v + v_0 \sin \theta) dt = (v + v_0 \sin \theta)t \\ x'_{Horse/ground} = \int_0^r v_e dt = \int_0^r v dt = vt \end{cases}$$

For the arrow to fall back on the man, it's necessary that  $(v + v_0 \sin \theta)t = vt \Rightarrow \sin \theta = 0$ .

Thus, it's required that  $\theta=0$ , meaning  $v_0$  should be vertical.

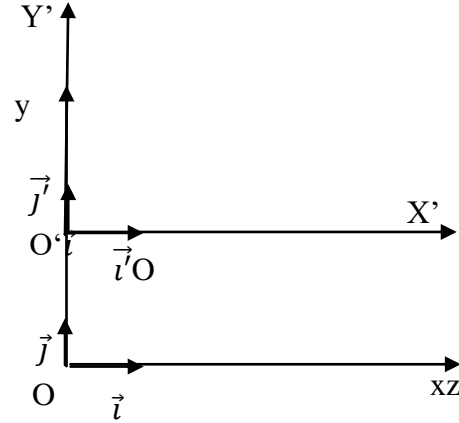
**Exercise 3**

$$\overrightarrow{OM}/(R) \begin{cases} x = 2t^3 + 1 \\ y = 4t^2 + t - 1 \\ z = t^2 \end{cases} \quad \text{and} \quad \overrightarrow{O'M}/(R') \begin{cases} x' = 2t^3 \\ y' = 4t^2 - 3t + 2 \\ z' = t^2 - 5 \end{cases}$$

1. The speed of point m in the fixed reference frame (R) and the moving reference frame (R').

$$\vec{v} = \vec{v}_a = \frac{d\overrightarrow{OM}}{dt} \begin{cases} \frac{dx}{dt} = 6t^2 \\ \frac{dy}{dt} = 8t + 1 \\ \frac{dz}{dt} = 2t \end{cases}$$

$$\vec{v}' = \vec{v}_r = \frac{d\overrightarrow{O'M}}{dt} \begin{cases} \frac{dx'}{dt} = 6t^2 \\ \frac{dy'}{dt} = 8t - 3 \\ \frac{dz'}{dt} = 2t \end{cases}$$



So:

$$\vec{v} = 6t^2\vec{i} + (8t + 1)\vec{j} + 2t\vec{k} \quad \text{And} \quad \vec{v}' = 6t^2\vec{i}' + (8t - 3)\vec{j}' + 2t\vec{k}'$$

$$\text{We have : } \vec{v}_a = \vec{v}_r + \vec{v}_e \Rightarrow \vec{v}_e = \vec{v}_a - \vec{v}_r$$

$$\vec{v}_e = (6t^2\vec{i} + (8t + 1)\vec{j} + 2t\vec{k}) - (6t^2\vec{i}' + (8t - 3)\vec{j}' + 2t\vec{k}') = 4\vec{j}$$

$$\text{So: } \vec{v} = \vec{v}' + 4\vec{j}$$

$$\text{And } \vec{i} = \vec{i}', \vec{j} = \vec{j}', \vec{k} = \vec{k}'$$

2. The acceleration of point m in the two reference frames, fixed (R) and moving (R'), is as follows:

$$\vec{a} = \vec{a}_a = \frac{d\vec{v}}{dt} \begin{cases} \frac{dv_x}{dt} = 12t \\ \frac{dv_y}{dt} = 8 \\ \frac{dv_z}{dt} = 2 \end{cases}$$

$$\text{And } \vec{a}' = \vec{a}_r = \frac{d\vec{v}'}{dt} \begin{cases} \frac{dv'_x}{dt} = 12t \\ \frac{dv'_y}{dt} = 8 \\ \frac{dv'_z}{dt} = 2 \end{cases}$$

$$\text{So: } \vec{a} = \vec{a}' \text{ Or } \vec{a}_a = \vec{a}_r$$

**Conclusion:**

The motion of frame of reference (R') relative to the fixed frame of reference (R) is a uniform translational motion along axis (Oy) with a constant speed of 4m/s.

**Exercise 4**

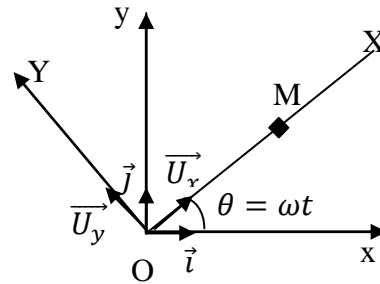
The fixed frame of reference and the moving frame of reference have the same origin, so O' and O are the same.

$$\text{Then, } \overline{OO'} = \vec{0} \text{ and } \overline{OM} = \overline{O'M} = \frac{1}{2}\gamma t^2 \vec{U}_x$$

$$\text{with } \vec{U}_x = \cos \omega t \vec{i} + \sin \omega t \vec{j}$$

$$\text{and } \vec{U}_y = (-\sin \omega t \vec{i} + \cos \omega t \vec{j})$$

$$\text{we have also } \vec{\omega} = \begin{pmatrix} 0 \\ 0 \\ \omega \end{pmatrix} \text{ and } \frac{d\omega}{dt} = 0$$



**1. Absolute velocity**

$$\vec{v}_r = \frac{d\overline{O'M}}{dt} / (OXY) \text{ with } \overline{O'M} = \overline{OM} = \frac{1}{2}\gamma t^2 \vec{U}_x \text{ so } \vec{v}_r = \gamma t \vec{U}_x$$

**Absolute acceleration :**

$$\vec{a}_r = \frac{d\vec{v}_r}{dt} / (O'XY) \text{ with } \vec{v}_r = \gamma t \vec{U}_x$$

$$\text{So } \vec{a}_r = \frac{d^2\overline{O'M}}{dt^2} = \gamma \vec{U}_x$$

**2. Entrainment velocity :**

$$\vec{v}_e = \frac{d\overline{OO'}}{dt} + \vec{\omega} \wedge \overline{O'M} \text{ with } \overline{OO'} = \vec{0} \text{ so } \vec{v}_e = \vec{\omega} \wedge \overline{O'M}$$

$$\vec{\omega} \wedge \overline{O'M} = \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ \frac{1}{2}\gamma t^2 & 0 & 0 \end{vmatrix} = \omega \frac{1}{2}\gamma t^2 \vec{U}_y \text{ so } \vec{v}_e = \omega \frac{1}{2}\gamma t^2 \vec{U}_y$$



**Entrainment acceleration**

$$\vec{a}_e = \frac{d^2 \overline{OO'}}{dt^2} + \vec{\omega} \Lambda (\vec{\omega} \Lambda \overline{O'M}) + \frac{d\vec{\omega}}{dt} \Lambda \overline{O'M}; \frac{d\vec{\omega}}{dt} \Lambda \overline{O'M} = \vec{0} \text{ because } \omega \text{ constant and } \frac{d^2 \overline{OO'}}{dt^2} = \vec{0}$$

$$\text{And } \vec{\omega} \Lambda (\vec{\omega} \Lambda \overline{O'M}) = \vec{\omega} \Lambda \left( \omega \frac{1}{2} \gamma t^2 \vec{U}_y \right) = \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ 0 & \omega \frac{1}{2} \gamma t^2 & 0 \end{vmatrix} = -\omega^2 \frac{1}{2} \gamma t^2 \vec{U}_x$$

$$\text{So } \vec{a}_e = -\omega^2 \frac{1}{2} \gamma t^2 \vec{U}_x$$

**3. Coriolis acceleration**

$$\vec{a}_c = 2\vec{\omega} \Lambda \vec{v}_r = 2 \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ \gamma t & 0 & 0 \end{vmatrix} = 2\omega \gamma t \vec{U}_y \text{ so } \vec{a}_c = 2\omega \gamma t \vec{U}_y$$

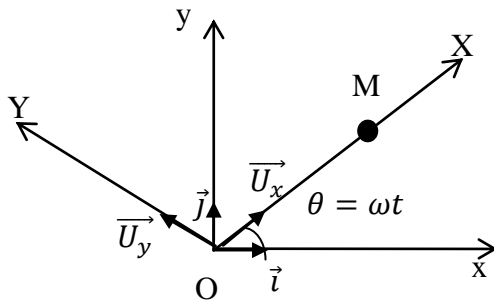
**4. Absolute velocity :**

$$\begin{aligned} \vec{v}_a &= \vec{v}_r + \vec{v}_e = \gamma t \vec{U}_x + \omega \frac{1}{2} \gamma t^2 \vec{U}_y \\ &= \gamma t (\cos \omega t \vec{i} + \sin \omega t \vec{j}) + \omega \frac{1}{2} \gamma t^2 (-\sin \omega t \vec{i} + \cos \omega t \vec{j}) \end{aligned}$$

**Acceleration absolute**

$$\begin{aligned} \vec{a}_a &= \vec{a}_r + \vec{a}_c + \vec{a}_e = (\gamma - \omega^2 \frac{1}{2} \gamma t^2) \vec{U}_x + 2\omega \gamma t \vec{U}_y \\ \vec{a}_a &= (\gamma - \omega^2 \frac{1}{2} \gamma t^2) (\cos \omega t \vec{i} + \sin \omega t \vec{j}) + 2\omega \gamma t (-\sin \omega t \vec{i} + \cos \omega t \vec{j}) \end{aligned}$$

**Exercise 5**



In the frame (OXY)  $\overline{OM} = v_0 t \vec{U}_x$

With  $\vec{U}_X = \cos \omega t \vec{i} + \sin \omega t \vec{j}$  And  $\vec{U}_Y = -\sin \omega t \vec{i} + \cos \omega t \vec{j}$

$\overline{OM}/(R) = v_0 t (\cos \omega t \vec{i} + \sin \omega t \vec{j})$  In the frame (Oxy)

**Absolute velocity**

$$\vec{v}_a = \frac{d\vec{OM}}{dt} / (R) = v_0 (\cos \omega t \vec{I} + \sin \omega t \vec{J}) - v_0 \omega t \sin \omega t \vec{I} + v_0 \omega t \cos \omega t \vec{J}$$

$$= v_0 (\cos \omega t \vec{I} + v_0 \sin \omega t \vec{J}) + v_0 \omega t (-\sin \omega t \vec{I} + \cos \omega t \vec{J}) = v_0 \vec{U}_X + v_0 \omega t \vec{U}_Y$$

(or)  $\vec{v}_a = \frac{d\vec{OM}}{dt} / (R) = \frac{d(v_0 t \vec{U}_x)}{dt} = v_0 \vec{U}_x + v_0 t \frac{d\vec{U}_x}{dt}$ ,

with  $\frac{d\vec{U}_x}{dt} = \frac{d\vec{U}_x}{d\theta} \frac{d\theta}{dt}$  and  $\theta = \omega t$

$$\Rightarrow \frac{d\vec{U}_x}{dt} = \frac{d\vec{U}_x}{d\theta} \frac{d\omega t}{dt} = \omega \vec{U}_y \text{ because } \frac{d\vec{U}_x}{d\theta} = \vec{U}_y$$

So :  $\vec{v}_a = v_0 \vec{U}_X + v_0 \omega t \vec{U}_Y$

**Absolute acceleration**

$$\begin{aligned} \vec{a}_a &= \frac{d\vec{v}_a}{dt} / (R) \\ &= -v_0 \omega \sin \omega t \vec{I} + v_0 \omega \cos \omega t \vec{J} - v_0 \omega \sin \omega t \vec{I} + v_0 \omega \cos \omega t \vec{J} \\ &\quad - v_0 \omega^2 t \cos \omega t \vec{I} + v_0 \omega^2 t \sin \omega t \vec{J} \end{aligned}$$

$$\Rightarrow \vec{a}_a = \frac{d\vec{v}_a}{dt} / (R) = 2v_0 \omega (-\sin \omega t \vec{I} + \cos \omega t \vec{J}) - v_0 \omega^2 t (\cos \omega t \vec{I} + \sin \omega t \vec{J})$$

$$\Rightarrow \vec{a}_a = 2v_0 \omega \vec{U}_Y - v_0 \omega^2 t \vec{U}_X$$

**Relative velocity**

$$\vec{v}_r = \frac{d\vec{O'M}}{dt} / (OXY) \text{ with } \vec{O'M} = \vec{OM} = v_0 t \vec{U}_x \quad \text{so } \vec{v}_r = v_0 \vec{U}_x$$

**Entrainment velocity**

The fixed frame and the mobile frame have the same origin so  $O'$  is confused with  $O$ , then  $\vec{OO'} = \vec{0}$ .

$$\vec{v}_e = \frac{d\vec{OO'}}{dt} + X \frac{d\vec{U}_x}{dt} = \vec{0} + v_0 t \frac{d\vec{U}_x}{dt} \quad \text{With} \quad \frac{d\vec{U}_x}{dt} = \omega \vec{U}_y$$

So ;  $\vec{v}_e = \omega v_0 t \vec{U}_y$

**Or**

$$\vec{v}_e = \frac{d\vec{OO'}}{dt} + \vec{\omega} \wedge \vec{O'M}$$

$$\vec{OO'} = \vec{0} \text{ so } \vec{v}_e = \vec{\omega} \wedge \vec{O'M}$$

$$\vec{\omega} \wedge \vec{O'M} = \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ v_0 t & 0 & 0 \end{vmatrix} = \omega v_0 t \vec{U}_y \text{ So } \vec{v}_e = \omega v_0 t \vec{U}_y$$

With  $\vec{U}_X = \cos \omega t \vec{I} + \sin \omega t \vec{J}$  And  $\vec{U}_Y = -\sin \omega t \vec{I} + \cos \omega t \vec{J}$

So  $\vec{v}_a = \vec{v}_r + \vec{v}_e = v_0 (\cos \omega t \vec{I} + \sin \omega t \vec{J}) + \omega v_0 t (-\sin \omega t \vec{I} + \cos \omega t \vec{J})$

$$\Rightarrow \vec{v}_a = v_0 \vec{U}_x + \omega v_0 t \vec{U}_y$$

So  $\vec{v}_a = \vec{v}_r + \vec{v}_e$  is verified

**Relative acceleration**

$$\vec{a}_r = \frac{d\vec{v}_r}{dt} \text{ with } \vec{v}_r = v_0 \vec{U}_X \text{ Then } \vec{a}_r = \vec{0}$$

**Entrainment acceleration**

$$\vec{a}_e = \frac{d^2 \vec{OO'}}{dt^2} + X \frac{d^2 \vec{U}_x}{dt^2} = v_0 t \frac{d}{dt} \left( \frac{d\vec{U}_x}{dt} \right) \Rightarrow \vec{a}_e = v_0 t \frac{d}{dt} (\omega \vec{U}_y)$$

$$\Rightarrow \vec{a}_e = v_0 t \omega \left( \frac{d\vec{U}_y}{dt} \right) \text{ With } \frac{d\vec{U}_y}{dt} = \frac{d\vec{U}_y}{d\theta} \frac{d\omega t}{dt} = -\omega \vec{U}_x$$

$$\Rightarrow \vec{a}_e = v_0 t \omega (+ (-\omega \vec{U}_x))$$

$$\Rightarrow \vec{a}_e = -v_0 \omega^2 t \vec{U}_x$$

**Or**

$$\vec{a}_e = \frac{d^2 \vec{OO'}}{dt^2} + \vec{\omega} \wedge (\vec{\omega} \wedge \vec{O'M}) + \frac{d\vec{\omega}}{dt} \wedge \vec{O'M}$$

$$\frac{d\vec{\omega}}{dt} \Lambda \overrightarrow{O'M} = \vec{0} \text{ because } \omega \text{ constant and } \frac{d^2 \overrightarrow{OO'}}{dt^2} = \vec{0}$$

$$\text{And } \vec{\omega} \Lambda (\vec{\omega} \Lambda \overrightarrow{O'M}) = \vec{\omega} \Lambda (\omega v_0 t \overrightarrow{U}_y) = \begin{vmatrix} \overrightarrow{U}_x & \overrightarrow{U}_y & \overrightarrow{U}_z \\ 0 & 0 & \omega \\ 0 & \omega v_0 t & 0 \end{vmatrix} = -\omega^2 v_0 t \overrightarrow{U}_x$$

$$\text{So ; } \vec{a}_e = -\omega^2 v_0 t \overrightarrow{U}_x$$

### Coriolis acceleration

$$\vec{a}_c = 2 \frac{dX}{dt} \frac{d\overrightarrow{U}_x}{dt} \text{ with } = v_0 t \text{ et } \frac{d\overrightarrow{U}_x}{dt} = \omega \overrightarrow{U}_y$$

$$\text{So } \vec{a}_c = 2v_0 \omega \overrightarrow{U}_y$$

**Or**

$$\vec{a}_c = 2\vec{\omega} \Lambda \vec{v}_r = 2 \begin{vmatrix} \overrightarrow{U}_x & \overrightarrow{U}_y & \overrightarrow{U}_z \\ 0 & 0 & \omega \\ v_0 & 0 & 0 \end{vmatrix} = 2\omega v_0 \overrightarrow{U}_y$$

$$\text{So } \vec{a}_c = 2\omega v_0 \overrightarrow{U}_y$$

### Absolute acceleration

$$\vec{a}_a = \vec{a}_r + \vec{a}_c + \vec{a}_e = -\omega^2 v_0 t \overrightarrow{U}_x + 2\omega v_0 \overrightarrow{U}_y$$

$$\Rightarrow \vec{a}_a = -\omega^2 v_0 t (\text{Cos } \omega t \vec{I} + \text{Sin } \omega t \vec{J}) + 2\omega v_0 (-\text{sin } \omega t \vec{I} + \text{Cos } \omega t \vec{J})$$

So

$$\vec{a}_a = \vec{a}_r + \vec{a}_c + \vec{a}_e \text{ Is verified}$$

### Exercise 6

$$\text{In the mobile frame (OXY): } \overrightarrow{OM} = r \overrightarrow{U}_x$$

In the mobile frame (polar coordinates) we have:

### Relative velocity

$$\vec{v}_r = \frac{d\overrightarrow{O'M}}{dt} / (OXY) = \frac{dr}{dt} \overrightarrow{U}_x \vec{v}_r = r \cdot \overrightarrow{U}_x$$

### Relative acceleration

$$\vec{a}_r = \frac{d\vec{v}_r}{dt} / (OXY) \quad \text{with } \vec{v}_r = r \cdot \vec{U}_X$$

$$\text{So : } \vec{a}_r = \frac{d^2 r}{dt^2} \vec{U}_X = r'' \vec{U}_X$$

### Entrainment velocity

$\overrightarrow{OO'}$  =  $\vec{0}$  "because both reference frames have the same origin."

$$\vec{v}_e = \frac{d\overrightarrow{OO'}}{dt} + X \frac{d\vec{U}_x}{dt} = \vec{0} + r \frac{d\vec{U}_x}{dt}$$

$$\text{With } \frac{d\vec{U}_x}{dt} = \frac{d\vec{U}_x}{d\theta} \frac{d\theta}{dt} \quad \text{and } \theta = \omega t$$

$$\text{So } \frac{d\vec{U}_x}{dt} = \frac{d\vec{U}_x}{d\theta} \frac{d\omega t}{dt} = \omega \vec{U}_y \quad \text{car } \frac{d\vec{U}_x}{d\theta} = \vec{U}_y$$

$$\text{Then } \vec{v}_e = \omega r \vec{U}_y$$

$$\underline{\text{Or}} \quad \vec{v}_e = \frac{d\overrightarrow{OO'}}{dt} + \vec{\omega} \wedge \overrightarrow{O'M}$$

$$\overrightarrow{OO'} = \vec{0} \quad \text{so } \vec{v}_e = \vec{\omega} \wedge \overrightarrow{O'M}$$

$$\vec{\omega} \wedge \overrightarrow{O'M} = \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ r & 0 & 0 \end{vmatrix} = \omega r \vec{U}_y \quad \text{So } \vec{v}_e = \omega r \vec{U}_y$$

### Entrainment acceleration

$$\vec{a}_e = \frac{d^2 \overrightarrow{OO'}}{dt^2} + X \frac{d^2 \vec{U}_x}{dt^2} = r \frac{d}{dt} \left( \frac{d\vec{U}_x}{dt} \right) \Rightarrow r \frac{d}{dt} (\omega \vec{U}_y)$$

$$\Rightarrow \vec{a}_e = r \omega \left( \frac{d\vec{U}_y}{dt} \right) \quad \text{With } \frac{d\vec{U}_y}{dt} = \frac{d\vec{U}_y}{d\theta} \frac{d\omega t}{dt} = -\omega \vec{U}_x$$

$$\Rightarrow \vec{a}_e = r \omega \left( + (-\omega \vec{U}_x) \right)$$

$$\Rightarrow \vec{a}_e = - r \omega^2 \vec{U}_x$$

**Or**

$$\vec{a}_e = \frac{d^2\overline{OO'}}{dt^2} + \vec{\omega}\Lambda(\vec{\omega}\Lambda\overline{O'M}) + \frac{d\vec{\omega}}{dt}\Lambda\overline{O'M}$$

$$\frac{d\vec{\omega}}{dt}\Lambda\overline{O'M} = \vec{0} \text{ because } \omega \text{ constant and } \frac{d^2\overline{OO'}}{dt^2} = \vec{0}$$

$$\vec{\omega}\Lambda(\vec{\omega}\Lambda\overline{O'M}) = \vec{\omega}\Lambda(\omega r \overline{U}_y) = \begin{vmatrix} \overline{U}_x & \overline{U}_y & \overline{U}_z \\ 0 & 0 & \omega \\ 0 & \omega r & 0 \end{vmatrix} = -\omega^2 r \overline{U}_x$$

So  $\vec{a}_e = -\omega^2 r \overline{U}_x$

**Coriolis acceleration**

$$\vec{a}_c = 2 \frac{dX}{dt} \frac{d\overline{U}_x}{dt} \text{ with } X = r \text{ and } \frac{d\overline{U}_x}{dt} = \omega \overline{U}_y$$

So  $\vec{a}_c = 2r \cdot \omega \overline{U}_y$

**Or**

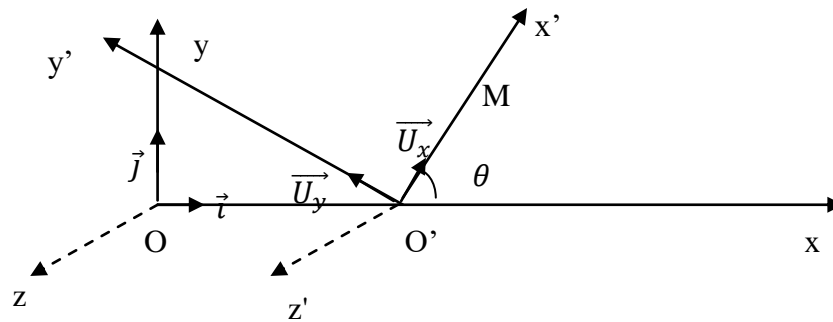
$$\vec{a}_c = 2\vec{\omega}\Lambda\vec{v}_r = 2 \begin{vmatrix} \overline{U}_x & \overline{U}_y & \overline{U}_z \\ 0 & 0 & \omega \\ r & 0 & 0 \end{vmatrix} = 2\omega r \cdot \overline{U}_y$$

So  $\vec{a}_c = 2\omega r \cdot \overline{U}_y$

**Absolute velocity**  $\vec{v}_a = \vec{v}_r + \vec{v}_e = r \cdot \overline{U}_x + \omega r \overline{U}_y$

**Absolute acceleration**  $\vec{a}_a = \vec{a}_r + \vec{a}_c + \vec{a}_e = (r'' - \omega^2 r) \overline{U}_x + 2\omega r \cdot \overline{U}_y$

**Exercise 7**



**Relative velocity**

$$\vec{v}_r = \frac{d\overline{O'M}}{dt} / (R') \text{ with } \overline{O'M} = t^2 \overline{U}_x \text{ so } \vec{v}_r = 2t \overline{U}_x$$

**Entrainment velocity:**

$$\vec{v}_e = \frac{d\overline{OO'}}{dt} + \vec{\omega} \wedge \overline{\Lambda O'M}$$

We are looking for the vector  $\overline{OO'}$ .

The point  $O'$  moves along the axis (Ox) with a velocity  $v$ , thus  $\vec{v}_{O'} = \frac{d\overline{OO'}}{dt} \vec{i} = v \vec{i}$

At  $t=0, x=0$  so  $\frac{d\overline{OO'}}{dt} = v \Rightarrow \overline{OO'} = vt$  so  $\overline{OO'} = vt \vec{i}$

$$\vec{\omega} \wedge \overline{\Lambda O'M} = \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ t^2 & 0 & 0 \end{vmatrix} = \omega t^2 \vec{U}_y \quad \text{and} \quad \frac{d\overline{OO'}}{dt} = v \vec{i}$$

$$\text{So } \vec{v}_e = \omega t^2 \vec{U}_y + v \vec{i}$$

"we need to express  $\vec{v}_e$  in the same coordinate system. To do this, we will express  $\vec{i}$  in terms of  $\vec{U}_x$  and  $\vec{U}_y$ .

$$\text{We have } \begin{cases} \vec{U}_x = \cos \theta \vec{i} + \sin \theta \vec{j} \\ \vec{U}_y = -\sin \theta \vec{i} + \cos \theta \vec{j} \end{cases} \Rightarrow \vec{i} = \cos \theta \vec{U}_x - \sin \theta \vec{U}_y$$

$$\text{So } \vec{v}_e = \omega t^2 \vec{U}_y + v(\cos \theta \vec{U}_x - \sin \theta \vec{U}_y) = v \cos \theta \vec{U}_x + (\omega t^2 - v \sin \theta) \vec{U}_y$$

**Absolute velocity**

$$\begin{aligned} \vec{v}_a &= \vec{v}_r + \vec{v}_e = 2t \vec{U}_x + \omega t^2 \vec{U}_y + v(\cos \theta \vec{U}_x - \sin \theta \vec{U}_y) \\ &\Rightarrow \vec{v}_a = (2t + v \cos \theta) \vec{U}_x + (\omega t^2 - v \sin \theta) \vec{U}_y \end{aligned}$$

**Relative acceleration**

$$\vec{a}_r = \frac{d\vec{v}_r}{dt} / (R') \text{ with } \vec{v}_r = 2t \vec{U}_x \text{ so } \vec{a}_r = 2 \vec{U}_x$$

**Entrainment acceleration**

$$\vec{a}_e = \frac{d^2 \overline{OO'}}{dt^2} + \vec{\omega} \wedge (\vec{\omega} \wedge \overline{\Lambda O'M}) + \frac{d\vec{\omega}}{dt} \wedge \overline{\Lambda O'M}$$

$$\frac{d^2 \overline{OO'}}{dt^2} = \vec{0}, \frac{d\vec{\omega}}{dt} \wedge \overline{\Lambda O'M} = \vec{0} \quad \Omega \text{ constant}$$

$$\text{And } \vec{\omega}\Lambda(\vec{\omega}\Lambda\vec{O'M}) = \vec{\omega}\Lambda\omega t^2\vec{U}_y = \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ 0 & \omega t^2 & 0 \end{vmatrix} = -\omega^2 t^2 \vec{U}_x$$

$$\text{So } \vec{a}_e = -\omega^2 t^2 \vec{U}_x$$

**Coriolis acceleration**

$$\vec{a}_c = 2\vec{\omega}\Lambda\vec{v}_r = 2 \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ 2t & 0 & 0 \end{vmatrix} = 4t\omega\vec{U}_y$$

**Absolute acceleration**

$$\vec{a}_a = \vec{a}_r + \vec{a}_c + \vec{a}_e = 2\vec{U}_x - \omega^2 t^2 \vec{U}_x + 4t\omega\vec{U}_y$$

$$\text{So } \vec{a}_a = (2 - \omega^2 t^2)\vec{U}_x + 4t\omega\vec{U}_y$$

**Exercise 8**

The coordinates of point m in the moving reference frame  $M(t^2, t)/(R')$ .

$$\text{So } \vec{O'M} \text{ is written: } \vec{O'M} = t^2\vec{U}_x + t\vec{U}_y$$

Or O' moves on the axis (Oy) with a constant acceleration  $\gamma$ .

at instant  $t=0$ , the axis (O'x) is confused with (Ox). So  $v_0=0$  and  $y_0=0$

$$\text{then; the acceleration of O' is: } \gamma = \frac{dv}{dt} \Rightarrow dv = \gamma dt$$

After integration  $v = \gamma \cdot t$

$$\text{And } \frac{dy}{dt} = \gamma t \Rightarrow dy = \gamma t dt \quad \text{so } y = \frac{1}{2}\gamma t^2 \text{ And } \vec{OO'} = \frac{1}{2}\gamma t^2 \vec{j}$$

$$\text{Relative velocity: } \vec{v}_r = \frac{d\vec{O'M}}{dt} / (R') \text{ with } \vec{O'M} = t^2\vec{U}_x + t\vec{U}_y \text{ So } \vec{v}_r = 2t\vec{U}_x + \vec{U}_y$$

$$\text{Entrainment velocity: } \vec{v}_e = \frac{d\vec{OO'}}{dt} + \vec{\omega}\Lambda\vec{O'M} \text{ With } \vec{OO'} = \frac{1}{2}\gamma t^2 \vec{j} \Rightarrow \frac{d\vec{OO'}}{dt} = \gamma t \vec{j}$$

$$\vec{\omega}\Lambda\vec{O'M} = \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ t^2 & t & 0 \end{vmatrix} = -\omega t \vec{U}_x + \omega t^2 \vec{U}_y \quad \text{so } \vec{v}_e = \gamma t \vec{j} - \omega t \vec{U}_x + \omega t^2 \vec{U}_y$$



## Chapter IV: Relative motion

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We need to write  $\vec{v}_e$  in the same coordinate system, so we'll write  $\vec{j}$  as a function of  $\vec{U}_x$  and  $\vec{U}_y$  :

$$\text{We have: } \begin{cases} \vec{U}_x = \cos\theta\vec{i} + \sin\theta\vec{j} \\ \vec{U}_y = -\sin\theta\vec{i} + \cos\theta\vec{j} \end{cases} \Rightarrow \vec{j} = \sin\theta\vec{U}_x + \cos\theta\vec{U}_y$$

$$\text{So } \vec{v}_e = \omega t^2\vec{U}_y - \omega t\vec{U}_x + \gamma t(\sin\theta\vec{U}_x + \cos\theta\vec{U}_y)$$

$$\Rightarrow \vec{v}_e = (\gamma t \sin \theta - \omega t)\vec{U}_x + (\omega t^2 + \gamma t \cos \theta)\vec{U}_y$$

**Absolute velocity :**

$$\vec{v}_a = \vec{v}_r + \vec{v}_e = 2t\vec{U}_x + \vec{U}_y + (\gamma t \sin \theta - \omega t)\vec{U}_x + (\omega t^2 + \gamma t \cos \theta)\vec{U}_y$$

$$\Rightarrow \vec{v}_a = (\gamma t \sin \theta - \omega t + 2t)\vec{U}_x + (\omega t^2 + \gamma t \cos \theta + 1)\vec{U}_y$$

**Relative acceleration :**

$$\vec{a}_r = \frac{d\vec{v}_r}{dt} / (R') \text{ with } \vec{v}_r = 2t\vec{U}_x + \vec{U}_y \text{ So } \vec{a}_r = 2\vec{U}_x$$

**Entrainement acceleration :**

$$\vec{a}_e = \frac{d^2\vec{OO'}}{dt^2} + \vec{\omega}\Lambda(\vec{\omega}\Lambda\vec{O'M}) + \frac{d\vec{\omega}}{dt}\Lambda\vec{O'M}$$

$$\frac{d\vec{\omega}}{dt}\Lambda\vec{O'M} = \vec{0} \text{ Because } \omega \text{ constant and } \frac{d^2\vec{OO'}}{dt^2} = \gamma\vec{j}$$

$$\text{and } \vec{\omega}\Lambda(\vec{\omega}\Lambda\vec{O'M}) = \vec{\omega}\Lambda(-\omega t\vec{U}_x + \omega t^2\vec{U}_y) = \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ -\omega t & \omega t^2 & 0 \end{vmatrix} = -\omega^2 t^2 \vec{U}_x - \omega^2 t \vec{U}_y$$

$$\text{Then } \vec{a}_e = \vec{j} - \omega^2 t^2 \vec{U}_x - \omega^2 t \vec{U}_y = \gamma(\sin\theta\vec{U}_x + \cos\theta\vec{U}_y) - \omega^2 t^2 \vec{U}_x - \omega^2 t \vec{U}_y$$

$$\vec{a}_e = (\gamma \sin \theta - \omega^2 t^2)\vec{U}_x + (\gamma \cos \theta - \omega^2 t)\vec{U}_y$$

**Coriolis acceleration :**

$$\vec{a}_c = 2\vec{\omega}\Lambda\vec{v}_r = 2 \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ 2t & 1 & 0 \end{vmatrix} = 4t\omega\vec{U}_y - 2\omega\vec{U}_x$$

## Chapter IV: Relative motion

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### Absolute acceleration :

$$\vec{a}_a = \vec{a}_r + \vec{a}_c + \vec{a}_e$$

$$\Rightarrow \vec{a}_a = 2\vec{U}_x + (\gamma \sin \theta - \omega^2 t^2)\vec{U}_x + (\gamma \cos \theta - \omega^2 t)\vec{U}_y + 4t \omega \vec{U}_y - 2\omega \vec{U}_x$$

$$\text{So } \vec{a}_a = (2 - 2\omega + \gamma \sin \theta - \omega^2 t^2)\vec{U}_x + (\gamma \cos \theta - \omega^2 t + 4t\omega)\vec{U}_y$$

### Exercise 9

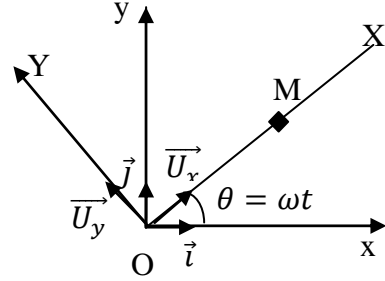
The fixed frame of reference and the moving frame of reference have the same origin, so O' and O are the same.

$$\text{Then } \overrightarrow{OO'} = \vec{0} \text{ and } \overrightarrow{OM} = \overrightarrow{O'M} = |\overrightarrow{OM}| \vec{U}_x = r = r_0(1 + \sin \omega t)\vec{U}_x$$

$$\text{with } \vec{U}_x = \cos \omega t \vec{i} + \sin \omega t \vec{j}$$

$$\text{and } \vec{U}_y = (-\sin \omega t \vec{i} + \cos \omega t \vec{j})$$

$$\text{we have also ; } \vec{\omega} = \begin{pmatrix} 0 \\ 0 \\ \omega \end{pmatrix} \text{ and } \frac{d\omega}{dt} = 0$$



### Relative velocity

$$\vec{v}_r = \frac{d\overrightarrow{O'M}}{dt} / (OXY) \text{ so } \vec{v}_r = r_0(\omega \cos \omega t)\vec{U}_x$$

### Relative acceleration

$$\vec{a}_r = \frac{d\vec{v}_r}{dt} / (O'XY) \text{ so } \vec{a}_r = \frac{d^2\overrightarrow{O'M}}{dt^2} = -r_0(\omega^2 \sin \omega t)\vec{U}_x$$

### Entrainment velocity

$$\vec{v}_e = \frac{d\overrightarrow{OO'}}{dt} + \vec{\omega} \wedge \overrightarrow{O'M} \text{ with } \overrightarrow{OO'} = \vec{0} \text{ so } \vec{v}_e = \vec{\omega} \wedge \overrightarrow{O'M}$$

$$\vec{\omega} \wedge \overrightarrow{O'M} = \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ r_0(1 + \sin \omega t) & 0 & 0 \end{vmatrix} = \omega r_0(1 + \sin \omega t) \vec{U}_y$$

$$\text{so } \vec{v}_e = \omega r_0(1 + \sin \omega t) \vec{U}_y$$

### Entrainment acceleration

$$\vec{a}_e = \frac{d^2\overrightarrow{OO'}}{dt^2} + \vec{\omega} \wedge (\vec{\omega} \wedge \overrightarrow{O'M}) + \frac{d\vec{\omega}}{dt} \wedge \overrightarrow{O'M} ; \frac{d\vec{\omega}}{dt} \wedge \overrightarrow{O'M} = \vec{0} \text{ because } \omega \text{ constant and } \frac{d^2\overrightarrow{OO'}}{dt^2} = \vec{0}$$

$$\text{And } \vec{\omega} \wedge (\vec{\omega} \wedge \overrightarrow{O'M}) = \vec{\omega} \wedge (\omega \frac{1}{2} \gamma t^2 \vec{U}_y) = \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ 0 & \omega r_0(1 + \sin \omega t) & 0 \end{vmatrix}$$

$$= -\omega^2 r_0 (1 + \sin \omega t) \vec{U}_x$$

$$\text{so } \vec{a}_e = -\omega^2 r_0 (1 + \sin \omega t) \vec{U}_x$$

**Coriolis acceleration**

$$\vec{a}_c = 2\vec{\omega} \wedge \vec{v}_r = 2 \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ r_0(\omega \cos \omega t) & 0 & 0 \end{vmatrix} = 2\omega r_0 (\omega \cos \omega t) \vec{U}_y$$

$$\text{so } \vec{a}_c = 2\omega r_0 (\omega \cos \omega t) \vec{U}_y$$

**Absolute velocity**

$$\vec{v}_a = \vec{v}_r + \vec{v}_e = r_0 (\omega \cos \omega t) \vec{U}_x + \omega r_0 (1 + \sin \omega t) \vec{U}_y$$

In the Cartesian coordinate base:

$$\vec{v}_a = \gamma t (\cos \omega t \vec{i} + \sin \omega t \vec{j}) + \omega \frac{1}{2} \gamma t^2 (-\sin \omega t \vec{i} + \cos \omega t \vec{j})$$

**Absolute acceleration**

$$\vec{a}_a = \vec{a}_r + \vec{a}_c + \vec{a}_e$$

$$\vec{a}_a = 2\omega r_0 (\omega \cos \omega t) \vec{U}_y - \omega^2 r_0 (1 + \sin \omega t) \vec{U}_x - r_0 (\omega^2 \sin \omega t) \vec{U}_x$$

In the Cartesian coordinate base:

$$\vec{a}_a = (-r_0 (\omega^2 \sin \omega t) - \omega^2 r_0 (1 + \sin \omega t)) (\cos \omega t \vec{i} + \sin \omega t \vec{j}) + 2\omega r_0 (\omega \cos \omega t) (-\sin \omega t \vec{i} + \cos \omega t \vec{j})$$

**Exercise 10**

At  $t=0$ ,  $y'=0$  and  $v=v_0$

$$\vec{OM} = \vec{OO'} + \vec{O'M} \quad \text{with } \vec{OO'} = R(\cos \omega t \vec{i} + \sin \omega t \vec{j})$$

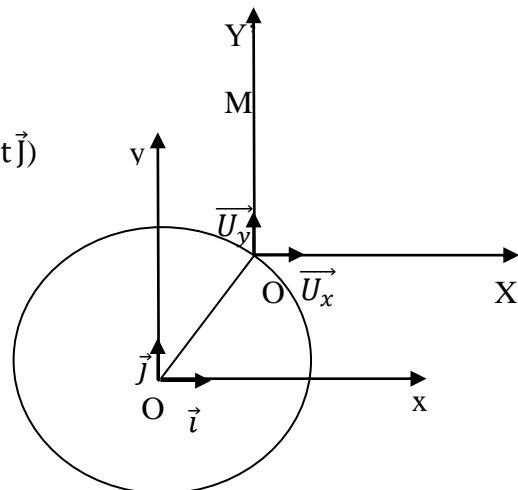
M moves along the axis (O'y) parallel to Oy.

With constant acceleration  $\gamma$ .

$$\vec{OM} = y \vec{U}_y, \quad \gamma = \frac{dv}{dt} \Rightarrow dv = \gamma dt$$

After integration  $v = \gamma t + v_0$

$$\frac{dy}{dt} = \gamma t + v_0 \Rightarrow dy = \gamma t dt + v_0 dt \quad \text{So } y = \frac{1}{2} \gamma t^2 + v_0 t$$



$$\overrightarrow{O'M} = \left(\frac{1}{2}\gamma t^2 + v_0 t\right) \overrightarrow{U}_y$$

Since, (O'y)/(oy) then  $\overrightarrow{U}_y = \vec{j}$  So  $\overrightarrow{O'M} = \left(\frac{1}{2}\gamma t^2 + v_0 t\right) \overrightarrow{U}_y = \left(\frac{1}{2}\gamma t^2 + v_0 t\right) \vec{j}$

Finally  $\overrightarrow{OM} = \overrightarrow{OO'} + \overrightarrow{O'M} = R(\cos \omega t \vec{i} + \sin \omega t \vec{j}) + \left(\frac{1}{2}\gamma t^2 + v_0 t\right) \vec{j}$

$$\Rightarrow \overrightarrow{OM} = R \cos \omega t \vec{i} + \left(r \sin \omega t + \frac{1}{2}\gamma t^2 + v_0 t\right) \vec{j}$$

**Absolute velocity :**

$$\vec{v}_a = \frac{d\overrightarrow{OM}}{dt} / (R) = -r \omega \sin \omega t \vec{i} + (r \omega \cos \omega t + \gamma t + v_0) \vec{j}$$

**Absolute acceleration :**

$$\vec{a}_a = \frac{d\vec{v}_a}{dt} / (R) = -r \omega^2 \cos \omega t \vec{i} + (-r \omega^2 \sin \omega t + \gamma) \vec{j}$$

**Relative velocity :**

$$\vec{v}_r = \frac{d\overrightarrow{O'M}}{dt} / (R') = (\gamma t + v_0) \vec{j}$$

**Entrainement velocity :**

$$\vec{v}_e = \frac{d\overrightarrow{OO'}}{dt} + \vec{\omega} \wedge \overrightarrow{O'M}$$

$\vec{\omega} \wedge \overrightarrow{O'M} = \vec{0}$  Because the unit vectors of the two marks are parallel, so there is no rotational movement.

**There is a translational movement**

$$\vec{v}_e = \frac{d\overrightarrow{OO'}}{dt} = -r \omega \sin \omega t \vec{i} + r \omega \cos \omega t \vec{j}$$

Let's check that :  $\vec{v}_a = \vec{v}_r + \vec{v}_e$

$$\begin{aligned} \vec{v}_a &= \vec{v}_r + \vec{v}_e = (\gamma t + v_0) \vec{j} - R \omega \sin \Omega t \vec{i} + r \omega \cos \omega t \vec{j} \\ &= -r \omega \sin \omega t \vec{i} + (r \omega \cos \omega t + \gamma t + v_0) \vec{j} \end{aligned}$$

So  $\vec{v}_a = \vec{v}_r + \vec{v}_e$  Is verified.

**Relative acceleration :**

$$\vec{a}_r = \left( \frac{d\vec{v}_r}{dt} \right) / R' \text{ with } \vec{v}_r = (\gamma t + v_0)\vec{j} \text{ So } \vec{a}_r = \gamma\vec{j}.$$

**Entrainement acceleration :**

$$\vec{a}_e = \frac{d^2\overrightarrow{OO'}}{dt^2} + \vec{\omega}\Lambda(\vec{\omega}\Lambda\overrightarrow{O'M}) + \frac{d\vec{\omega}}{dt}\Lambda\overrightarrow{O'M}$$

$$\frac{d\vec{\omega}}{dt}\Lambda\overrightarrow{O'M} = \vec{0} \text{ and } \vec{\omega}\Lambda(\vec{\omega}\Lambda\overrightarrow{O'M}) = \vec{0}$$

Because there is a translational movement between the reference marks.

$$\vec{a}_e = \frac{d^2\overrightarrow{OO'}}{dt^2} = -r\omega^2\text{Cos}\omega t\vec{i} + -r\omega^2\text{Sin}\omega t\vec{j} \text{ So } \vec{a}_e = -\omega^2r_0(\text{Cos}\omega t\vec{i} + \text{Sin}\omega t\vec{j})$$

**Coriolis acceleration :**

$$\vec{a}_c = 2\vec{\omega}\Lambda\vec{v}_r = \vec{0}$$

Let's check that :  $\vec{a}_a = \vec{a}_r + \vec{a}_c + \vec{a}_e$

$$\begin{aligned} \vec{a}_r + \vec{a}_c + \vec{a}_e &= \gamma\vec{j} + -r\omega^2\text{Cos}\omega t\vec{i} \pm r\omega^2\text{Sin}\omega t\vec{j} \\ &= -r\omega^2\text{Cos}\omega t\vec{i} + (-r\omega^2\text{Sin}\omega t + \gamma)\vec{j} \end{aligned}$$

So  $\vec{a}_a = \vec{a}_r + \vec{a}_c + \vec{a}_e$  Is verified

**Exercise 11**

$$r = r_0(\text{Cos}\omega t + \text{Sin}\omega t)\vec{U}_x$$

**Relative velocity:**  $\vec{v}_r = \frac{d\overrightarrow{O'M}}{dt} / (R')$

O' is confused with O, then :  $\overrightarrow{OM} = \overrightarrow{O'M} \Rightarrow \vec{v}_r = \frac{d\overrightarrow{OM}}{dt} / (R')$

$$\vec{v}_r = r_0\omega(-\text{sin}\omega t + \text{Cos}\omega t)\vec{U}_x$$

**Entrainment velocity :**  $\vec{v}_e = \frac{d\overrightarrow{OO'}}{dt} + \vec{\omega}\Lambda\overrightarrow{O'M}$  With  $\overrightarrow{OO'} = \vec{0} \Rightarrow \frac{d\overrightarrow{OO'}}{dt} = \vec{0}$

$$\vec{\omega}\Lambda\overrightarrow{O'M} = \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ r & 0 & 0 \end{vmatrix} = \omega r \vec{U}_y \text{ so } \vec{v}_e = \omega r \vec{U}_y = \omega r_0(\text{Cos}\omega t + \text{Sin}\omega t)\vec{U}_y$$

## Chapter IV: Relative motion

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**Absolute velocity :**  $\vec{v}_a = \vec{v}_r + \vec{v}_e = r_0\omega[(-\sin\omega t + \cos\omega t)\vec{U}_x + (\cos\omega t + \sin\omega t)\vec{U}_y]$

$$\Rightarrow |\vec{v}_a| = r_0\omega \sqrt{(-\sin\omega t + \cos\omega t)^2 + (\cos\omega t + \sin\omega t)^2}$$

Then  $|\vec{v}_a| = r_0\omega\sqrt{2}$  So  $|\vec{v}_a|$  Is constant

**Relative acceleration :**  $\vec{a}_r = \frac{d\vec{v}_r}{dt} / (R') \vec{v}_r = r_0\omega(-\sin\omega t + \cos\omega t)\vec{U}_x$

So  $\vec{a}_r = r_0\omega^2(-\cos\omega t - \sin\omega t)\vec{U}_x$

**Entrainment acceleration :**

$$\vec{a}_e = \frac{d^2\vec{OO'}}{dt^2} + \vec{\omega}\Lambda(\vec{\omega}\Lambda\vec{O'M}) + \frac{d\vec{\omega}}{dt}\Lambda\vec{O'M}$$

With  $\frac{d\vec{\omega}}{dt}\Lambda\vec{O'M} = \vec{0}$  Because  $\omega$  constant and  $\frac{d^2\vec{OO'}}{dt^2} = \vec{0}$

$$\text{And } \vec{\omega}\Lambda(\vec{\omega}\Lambda\vec{O'M}) = \vec{\omega}\Lambda(\omega r\vec{U}_y) = \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ 0 & \omega r & 0 \end{vmatrix} = -\omega^2 r\vec{U}_x$$

Then  $\vec{a}_e = -\omega^2 r_0(\cos\omega t + \sin\omega t)\vec{U}_x$

$$\text{Coriolis acceleration: } \vec{a}_c = 2\vec{\omega}\Lambda\vec{v}_r = 2 \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ v_r & 0 & 0 \end{vmatrix} = 2\omega v_r\vec{U}_y$$

So  $\vec{a}_c = 2\omega v_r\vec{U}_y = 2r_0\omega^2(-\sin\omega t + \cos\omega t)\vec{U}_y$

**Absolute acceleration :**  $\vec{a}_a = \vec{a}_r + \vec{a}_c + \vec{a}_e$

$$\vec{a}_a = -r_0\omega^2(\cos\omega t + \sin\omega t)\vec{U}_x - \omega^2 r_0(\cos\omega t + \sin\omega t)\vec{U}_x + 2r_0\omega^2(-\sin\omega t + \cos\omega t)\vec{U}_y$$

$$\Rightarrow \vec{a}_a = -2r_0\omega^2(\cos\omega t + \sin\omega t)\vec{U}_x + 2r_0\omega^2(-\sin\omega t + \cos\omega t)\vec{U}_y$$

$$|\vec{a}_a| = 2r_0\omega^2 \sqrt{(-(\cos\omega t + \sin\omega t))^2 + (-\sin\omega t + \cos\omega t)^2}$$

Then,  $|\vec{a}_a| = 2r_0\omega^2\sqrt{2}$  so  $|\vec{a}_a|$  Is constant.

**Exercise 12**

**1- Speeds:**

M moves along the (Oy') axis with constant acceleration, so:  $\overrightarrow{O'M} = Y \overrightarrow{u}_y$  with  $\gamma = \frac{dv}{dt}$  and at t=0 the point M is at O' :

$$\gamma = \frac{dv}{dt} \Rightarrow \int_0^v dv = \gamma \int_0^t dt \text{ so } v = \gamma t \text{ (at t=0, } v_0(M)=0)$$

$$v = \gamma t = \frac{dY}{dt} \Rightarrow \int_0^Y dY = \gamma \int_0^t t dt \text{ so } Y = \frac{1}{2} \gamma t^2 \text{ (at t=0, } Y_0(M)=0)$$

$$\overrightarrow{O'M} = \frac{1}{2} \gamma t^2 \overrightarrow{u}_y$$

O' moves on Ox with a constant speed  $v_0$  so  $\overrightarrow{OO'} = x\vec{i}$  and  $v_0 = \frac{dx}{dt}$  and à t=0, axis (O'x') is confused with (Ox).

$$v_0 = \frac{dx}{dt} \Rightarrow \int_0^x dx = v_0 \int_0^t dt \text{ so } x = v_0 t \text{ (at t=0, } x_0(O')=0) \text{ then } \overrightarrow{OO'} = v_0 t \vec{i}$$

$$\overrightarrow{v}_r = \frac{d\overrightarrow{O'M}}{dt} = \gamma t \overrightarrow{u}_y$$

$$\overrightarrow{v}_e = \frac{d\overrightarrow{OO'}}{dt} + \overline{\omega} \cdot \overrightarrow{O'M} \text{ with } \overline{\omega} = \begin{pmatrix} 0 \\ 0 \\ \omega \end{pmatrix}$$

$$\overrightarrow{u}_x = \cos\theta \vec{i} + \sin\theta \vec{j} \text{ and } \overrightarrow{u}_y = -\sin\theta \vec{i} + \cos\theta \vec{j}$$

Using the passage table :

$$\text{So } \vec{i} = \cos\theta \overrightarrow{u}_x - \sin\theta \overrightarrow{u}_y$$

	$\overrightarrow{u}_x$	$\overrightarrow{u}_y$
$\vec{i}$	Cos $\theta$	-sin $\theta$
$\vec{j}$	Sin $\theta$	cos $\theta$

$$\frac{d\overrightarrow{OO'}}{dt} = v_0 \vec{i} = v_0 (\cos\omega t \overrightarrow{u}_x - \sin\omega t \overrightarrow{u}_y)$$

$$\overline{\omega} \cdot \overrightarrow{O'M} = \begin{vmatrix} \overrightarrow{u}_x & \overrightarrow{u}_y & \overrightarrow{u}_z \\ 0 & 0 & \omega \\ 0 & \frac{1}{2} \gamma t^2 & 0 \end{vmatrix} = -\frac{1}{2} \gamma t^2 \omega \overrightarrow{u}_x$$

$$\overrightarrow{v}_e = \left( -\frac{1}{2} \gamma t^2 \omega + v_0 \cos\omega t \right) \overrightarrow{u}_x + (-v_0 \sin\omega t) \overrightarrow{u}_y$$

$$\overrightarrow{v}_a = \overrightarrow{v}_r + \overrightarrow{v}_e = \left( -\frac{1}{2} \gamma t^2 \omega + v_0 \cos\omega t \right) \overrightarrow{u}_x + (\gamma t - v_0 \sin\omega t) \overrightarrow{u}_y$$

**2- The accelerations :**

$$\overrightarrow{a}_r = \frac{d\overrightarrow{v}_r}{dt} = \gamma \overrightarrow{u}_y$$

$$\overrightarrow{a}_e = \frac{d^2\overrightarrow{OO'}}{dt^2} + \frac{d\overline{\omega}}{dt} \cdot \overrightarrow{O'M} + \overline{\omega} \cdot \overline{\omega} \cdot \overrightarrow{O'M} \text{ with } \frac{d^2\overrightarrow{OO'}}{dt^2} = \vec{0}$$

$$\vec{\omega} \cdot \dots \cdot \vec{\omega} \cdot \dots \cdot \vec{O'M} = \begin{vmatrix} \vec{u}_x & \vec{u}_y & \vec{u}_z \\ 0 & 0 & \omega \\ -\frac{1}{2}\gamma t^2 \omega & 0 & 0 \end{vmatrix} = -\frac{1}{2}\gamma t^2 \omega^2 \vec{u}_y$$

$$\vec{a}_e = -\frac{1}{2}\gamma t^2 \omega^2 \vec{u}_y$$

$$\vec{a}_c = 2\vec{\omega} \cdot \dots \cdot \vec{v}_r = \begin{vmatrix} \vec{u}_x & \vec{u}_y & \vec{u}_z \\ 0 & 0 & 2\omega \\ 0 & \gamma t & 0 \end{vmatrix} = -2\gamma t \omega \vec{u}_x$$

$$\vec{a}_a = \vec{a}_r + \vec{a}_e + \vec{a}_c$$

$$\text{So } \vec{a}_a = (-2\gamma t \omega) \vec{u}_x + \left(\gamma - \frac{1}{2}\gamma t^2 \omega^2\right) \vec{u}_y$$

### Exercise 13

$$r = r_0(\cos \omega t + \sin \omega t) \vec{U}_x$$

### Relative velocity

$$\vec{v}_r = \frac{d\vec{O'M}}{dt} / (R') \quad \text{with} \quad \vec{O'M} = \vec{O'M} \Rightarrow \vec{v}_r = \frac{d\vec{O'M}}{dt} / (R')$$

$$\vec{v}_r = r_0 \omega (-\sin \omega t + \cos \omega t) \vec{U}_x$$

### Entrainment velocity

$$\vec{v}_e = \frac{d\vec{OO'}}{dt} + \vec{\omega} \wedge \vec{O'M} \quad \text{with} \quad \vec{OO'} = t\vec{i} \Rightarrow \frac{d\vec{OO'}}{dt} = \vec{i}$$

$$\vec{\omega} \wedge \vec{O'M} = \begin{vmatrix} \vec{U}_x & \vec{U}_y & \vec{U}_z \\ 0 & 0 & \omega \\ r & 0 & 0 \end{vmatrix} = \omega r \vec{U}_y$$

$$\text{So } \vec{v}_e = \vec{i} + \omega r \vec{U}_y = (\cos \theta \vec{U}_x - \sin \theta \vec{U}_y) - \omega r_0 (\cos \omega t + \sin \omega t) \vec{U}_y$$

$$\vec{v}_e = \cos \theta \vec{U}_x - (\sin \theta + \omega r_0 (\cos \omega t + \sin \omega t)) \vec{U}_y$$

### Absolute velocity

$$\vec{v}_a = \vec{v}_r + \vec{v}_e = [(\cos \omega t + r_0 \omega (-\sin \omega t + \cos \omega t)) \vec{U}_x - (\sin \theta + r_0 \omega (\cos \omega t + \sin \omega t)) \vec{U}_y]$$

### Absolute acceleration

$$\vec{a}_r = \frac{d\vec{v}_r}{dt} / (R') \quad \text{with} \quad \vec{v}_r = r_0 \omega (-\sin \omega t + \cos \omega t) \vec{U}_x$$

$$\text{So } \vec{a}_r = r_0 \omega^2 (-\cos \omega t - \sin \omega t) \vec{U}_x$$



## Chapter IV: Relative motion

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### Entrainment acceleration

$$\vec{a}_e = \frac{d^2 \overline{OO'}}{dt^2} + \vec{\omega} \Lambda (\vec{\omega} \Lambda \overline{O'M}) + \frac{d\vec{\omega}}{dt} \Lambda \overline{O'M}$$

$$\frac{d\vec{\omega}}{dt} \Lambda \overline{O'M} = \vec{0} \text{ because } \omega \text{ constant and } \frac{d^2 \overline{OO'}}{dt^2} = \vec{0}$$

$$\text{And } \vec{\omega} \Lambda (\vec{\omega} \Lambda \overline{O'M}) = \vec{\omega} \Lambda (\omega r \overline{U}_y) = \begin{vmatrix} \overline{U}_x & \overline{U}_y & \overline{U}_z \\ 0 & 0 & \omega \\ 0 & \omega r & 0 \end{vmatrix} = -\omega^2 r \overline{U}_x$$

$$\text{so } \vec{a}_e = -\omega^2 r_0 (\cos \omega t + \sin \omega t) \overline{U}_x$$

### Coriolis acceleration

$$\vec{a}_c = 2\vec{\omega} \Lambda \vec{v}_r = 2 \begin{vmatrix} \overline{U}_x & \overline{U}_y & \overline{U}_z \\ 0 & 0 & \omega \\ v_r & 0 & 0 \end{vmatrix} = 2 \omega v_r \overline{U}_y$$

$$\text{so } \vec{a}_c = 2 \omega v_r \overline{U}_y = 2r_0 \omega^2 (-\sin \omega t + \cos \omega t) \overline{U}_y$$

### Absolute acceleration

$$\vec{a}_a = \vec{a}_r + \vec{a}_c + \vec{a}_e$$

$$\vec{a}_a = r_0 \omega^2 (-\cos \omega t - \sin \omega t) \overline{U}_x - \omega^2 r_0 (\cos \omega t + \sin \omega t) \overline{U}_x + 2r_0 \omega^2 (-\sin \omega t + \cos \omega t) \overline{U}_y$$

$$\Rightarrow \vec{a}_a = -2r_0 \omega^2 (\cos \omega t + \sin \omega t) \overline{U}_x + 2r_0 \omega^2 (-\sin \omega t + \cos \omega t) \overline{U}_y$$

### Exercise 14

$$\overline{OO'} = R(\cos \omega' t \vec{i} + \sin \omega' t \vec{j}) \text{ and } \overline{O'M} = d(\cos \Omega t \vec{i} + \sin \Omega t \vec{j})$$

$$\overline{OM} = \overline{OO'} + \overline{O'M} = R(\cos \omega' t \vec{i} + \sin \omega' t \vec{j}) + d(\cos \Omega t \vec{i} + \sin \Omega t \vec{j})$$

### Absoltue velocity

$$\vec{v}_a = \frac{d\overline{OM}}{dt} / (R) = \frac{d}{dt} (\overline{OO'} + \overline{O'M})$$

The axes (O'X), (O'Y) (moving reference) are parallel with the axes (Ox), (Oy) (fixed reference), therefore:

$$\vec{i} = \vec{i}' \text{ and } \vec{j} = \vec{j}'$$

## Chapter IV: Relative motion

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$$\vec{v}_a = R \omega' (-\sin \omega' t \vec{i} + \cos \omega' t \vec{j}) + d \cdot \Omega (-\sin \Omega t \vec{i} + \cos \Omega t \vec{j})$$

$$\Rightarrow \vec{v}_a = -(R \omega' \sin \omega' t + d \cdot \Omega \sin \Omega t) \vec{i} + (R \omega' \cos \omega' t + d \cdot \Omega \cos \Omega t) \vec{j}$$

### Relative velocity

$$\vec{v}_r = \frac{d\overline{O'M}}{dt} / (R') = d \cdot \Omega (-\sin \Omega t \vec{i} + \cos \Omega t \vec{j})$$

### Entrainment velocity

$$\vec{v}_e = \frac{d\overline{OO'}}{dt} + \vec{\omega} \wedge \overline{O'M}$$

$\vec{\omega} \wedge \overline{O'M} = \vec{0}$  Because the unit vectors of the two frames of reference are parallel, so there is a translational movement and not a rotational movement of axes of moving reference frame.

$$\vec{v}_e = \frac{d\overline{OO'}}{dt} = -R \omega' \sin \omega' t \vec{i} + R \omega' \cos \omega' t \vec{j}$$

### Absolute velocity

$$\begin{aligned} \vec{v}_a &= \vec{v}_r + \vec{v}_e = d \cdot \Omega (-\sin \Omega t \vec{i} + \cos \Omega t \vec{j}) + R \omega' (-\sin \omega' t \vec{i} + \cos \omega' t \vec{j}) \\ &= -(R \omega' \sin \omega' t + d \Omega \sin \Omega t) \vec{i} + (R \omega' \cos \omega' t + d \Omega \cos \Omega t) \vec{j} \end{aligned}$$

So  $\vec{v}_a = \vec{v}_r + \vec{v}_e$  Est is verified

### Absolute acceleration

$$\vec{a}_a = \frac{d\vec{v}_a}{dt} / (R) = \frac{d}{dt} [-(R \omega' \sin \omega' t + d \Omega \sin \Omega t) \vec{i} + (R \omega' \cos \omega' t + d \Omega \cos \Omega t) \vec{j}]$$

$$\vec{a}_a = -(R \omega'^2 \cos \omega' t + d \Omega^2 \cos \Omega t) \vec{i} - (R \omega'^2 \sin \omega' t + d \cdot \Omega^2 \sin \Omega t) \vec{j}$$

### Relative acceleration

$$\vec{a}_r = \frac{d\vec{v}_r}{dt} / (R') \text{ with } \vec{v}_r = d \cdot \Omega (-\sin \Omega t \vec{i} + \cos \Omega t \vec{j})$$

$$\text{Donc } \vec{a}_r = d \cdot \Omega^2 (-\cos \Omega t \vec{i} - \sin \Omega t \vec{j})$$

### Entrainment acceleration

$$\vec{a}_e = \frac{d^2\overline{OO'}}{dt^2} + \vec{\omega} \wedge (\vec{\omega} \wedge \overline{O'M}) + \frac{d\vec{\omega}}{dt} \wedge \overline{O'M}$$

$$\frac{d\vec{\omega}}{dt} \Lambda \overrightarrow{O'M} = \vec{0} \text{ and } \vec{\omega} \Lambda (\vec{\omega} \Lambda \overrightarrow{O'M}) = \vec{0}$$

Because there is a translation movement between the two references

$$\vec{a}_e = \frac{d^2 \overrightarrow{OO'}}{dt^2} = -R \omega'^2 \text{Cos } \omega't \vec{i} + -R \omega'^2 \text{Sin } \omega't \vec{j}$$

$$\text{So } \vec{a}_e = -\omega'^2 R (\text{Cos } \omega't \vec{i} + \text{Sin } \omega't \vec{j})$$

### Coriolis accélération

$$\vec{a}_c = 2\vec{\omega} \Lambda \vec{v}_r = \vec{0}$$

### Absolute accélération

$$\begin{aligned} \vec{a}_r + \vec{a}_c + \vec{a}_e &= d.\Omega^2 (-\text{Cos } \Omega t \vec{i} - \text{sin } \Omega t \vec{j}) + -R \omega'^2 \text{Cos } \omega't \vec{i} \pm R \omega'^2 \text{Sin } \omega't \vec{j} \\ &= -(R \omega'^2 \text{Cos } \omega't + d.\Omega^2 \text{Cos } \Omega t) \vec{i} - (R \omega'^2 \text{Sin } \omega't + d.\Omega^2 \text{Sin } \Omega t) \vec{j} \end{aligned}$$

So  $\vec{a}_a = \vec{a}_r + \vec{a}_c + \vec{a}_e$  is vérified

### Exercise 15

The axes (Ox) and (O'x') are not parallel, and the rotation is along the (Oz) axis.

The expression of (OM) vector in the fixed frame (Oxy).

$$\overrightarrow{OM} = \overrightarrow{OO'} + \overrightarrow{O'M} \text{ with } R=O'M=OO'$$

$$\overrightarrow{OO'} = R(\text{Cos } \omega t \vec{i} + \text{Sin } \omega t \vec{j}) \text{ and } \overrightarrow{O'M} = R(\text{Cos } \omega t \vec{i}' + \text{Sin } \omega t \vec{j}')$$

$$\text{with } \vec{i}' = \text{Cos } \omega t \vec{i} + \text{Sin } \omega t \vec{j} \text{ and } \vec{j}' = -\text{sin } \omega t \vec{i} + \text{Cos } \omega t \vec{j}$$

$$\text{So } \overrightarrow{O'M} = R(\text{Cos } \omega t (\text{Cos } \omega t \vec{i} + \text{Sin } \omega t \vec{j}) + \text{sin } \omega t (-\text{sin } \omega t \vec{i} + \text{Cos } \omega t \vec{j}))$$

$$\Rightarrow \overrightarrow{O'M} = R(\text{Cos } 2\omega t \vec{i} + \text{sin } 2\omega t \vec{j})$$

$$\text{So } \overrightarrow{OM} = R[(\text{Cos } \omega t \vec{i} + \text{Sin } \omega t \vec{j}) + (\text{Cos } 2\omega t \vec{i} + \text{sin } 2\omega t \vec{j})]$$

$$\Rightarrow \overrightarrow{OM} = R[(\text{Cos } 2\omega t + \text{Cos } \omega t) \vec{i} + (\text{sin } 2\omega t + \text{sin } \omega t) \vec{j}]$$

## Chapter IV: Relative motion

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*"The expression of absolute velocity and absolute acceleration in the fixed frame."*

### Absolute velocity

$$\vec{v}_a = \frac{d\overrightarrow{OM}}{dt} / (R) = \frac{d}{dt} (\overrightarrow{OO'} + \overrightarrow{O'M})$$

$$\Rightarrow \vec{v}_a = R [(-\omega \sin \omega t - 2\omega \sin 2\omega t) \vec{i} + (\omega \cos \omega t + 2\omega \cos \omega t) \vec{j}]$$

### Absolute accélération

$$\vec{a}_a = \frac{d\vec{v}_a}{dt} / (R) = \frac{d}{dt} (R [(-\omega \sin \omega t - 2\omega \sin 2\omega t) \vec{i} + (\omega \cos \omega t + 2\omega \cos \omega t) \vec{j}])$$

$$\Rightarrow \vec{a}_a = R [(-\omega^2 \cos \omega t - 2\omega^2 \cos 2\omega t) \vec{i} - (\omega^2 \sin \omega t + 2\omega^2 \sin \omega t) \vec{j}]$$

*"The expression of relative velocity and relative acceleration in the fixed frame."*

### Relative velocity

$$\vec{v}_r = \frac{d\overrightarrow{O'M}}{dt} / R' = R \omega (-\sin \omega t \vec{i}' + \cos \omega t \vec{j}')$$

### Relative acceleration

$$\vec{a}_r = \frac{d\vec{v}_r}{dt} / R' = -R \omega^2 (\cos \omega t \vec{i}' + \sin \omega t \vec{j}')$$

### Entrainment velocity

$$\vec{v}_e = \frac{d\overrightarrow{OO'}}{dt} + \vec{\omega} \wedge \overrightarrow{O'M} \quad \text{with} \quad \frac{d\overrightarrow{OO'}}{dt} = R \omega (-\sin \omega t \vec{i} + \cos \omega t \vec{j})$$

$$\vec{\omega} \wedge \overrightarrow{O'M} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k}' \\ 0 & 0 & \omega \\ R \cos \omega t & R \sin \omega t & 0 \end{vmatrix} = R \omega (-\sin \omega t \vec{i} + \cos \omega t \vec{j})$$

$$\Rightarrow \vec{v}_e = R \omega (-\sin \omega t \vec{i} + \cos \omega t \vec{j}) + R \omega (-\sin \omega t \vec{i}' + \cos \omega t \vec{j}')$$

$$\Rightarrow \vec{v}_e = R \omega (-\sin \omega t \vec{i} + \cos \omega t \vec{j}) + R \omega (-\sin \omega t (\cos \omega t \vec{i} + \sin \omega t \vec{j}) + \cos \omega t (-\sin \omega t \vec{i} + \cos \omega t \vec{j}))$$

$$\Rightarrow \vec{v}_e = R \omega [(-\sin \omega t \vec{i} + \cos \omega t \vec{j}) + (-\sin \omega t \cos \omega t - \cos^2 \omega t) \vec{i} + (-\sin^2 \omega t + \sin \omega t \cos \omega t) \vec{j}]$$

### Entrainment acceleration

$$\vec{a}_e = \frac{d^2 \overline{OO'}}{dt^2} + \vec{\omega} \Lambda (\vec{\omega} \Lambda \overline{O'M}) + \frac{d\vec{\omega}}{dt} \Lambda \overline{O'M}$$

$$\frac{d\vec{\omega}}{dt} \Lambda \overline{O'M} = \vec{0}$$

With  $\frac{d^2 \overline{OO'}}{dt^2} = -R \omega^2 (\cos \omega t \vec{i} + \sin \omega t \vec{j})$

$$\vec{\omega} \Lambda (\vec{\omega} \Lambda \overline{O'M}) = \vec{\omega} \Lambda R \omega (-\sin \omega t \vec{i} + \cos \omega t \vec{j})$$

$$= \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ 0 & 0 & \omega \\ -R\omega \sin \omega t & R\omega \cos \omega t & 0 \end{vmatrix}$$

$$= -R \omega^2 (\cos \omega t \vec{i} + \sin \omega t \vec{j})$$

$$\Rightarrow \vec{a}_e = -R \omega^2 (\cos \omega t \vec{i} + \sin \omega t \vec{j}) - R \omega^2 (\cos \omega t \vec{i} + \sin \omega t \vec{j})$$

$$\Rightarrow \vec{a}_e = -R \omega^2 (\cos \omega t \vec{i} + \sin \omega t \vec{j}) - R \omega^2 (\cos \omega t (\cos \omega t \vec{i} + \sin \omega t \vec{j}) + \sin \omega t (-\sin \omega t \vec{i} + \cos \omega t \vec{j}))$$

$$\Rightarrow \vec{a}_e = -R \omega^2 [(\cos \omega t \vec{i} + \sin \omega t \vec{j}) - (\cos 2\omega t \vec{i} + \sin 2\omega t \vec{j})]$$

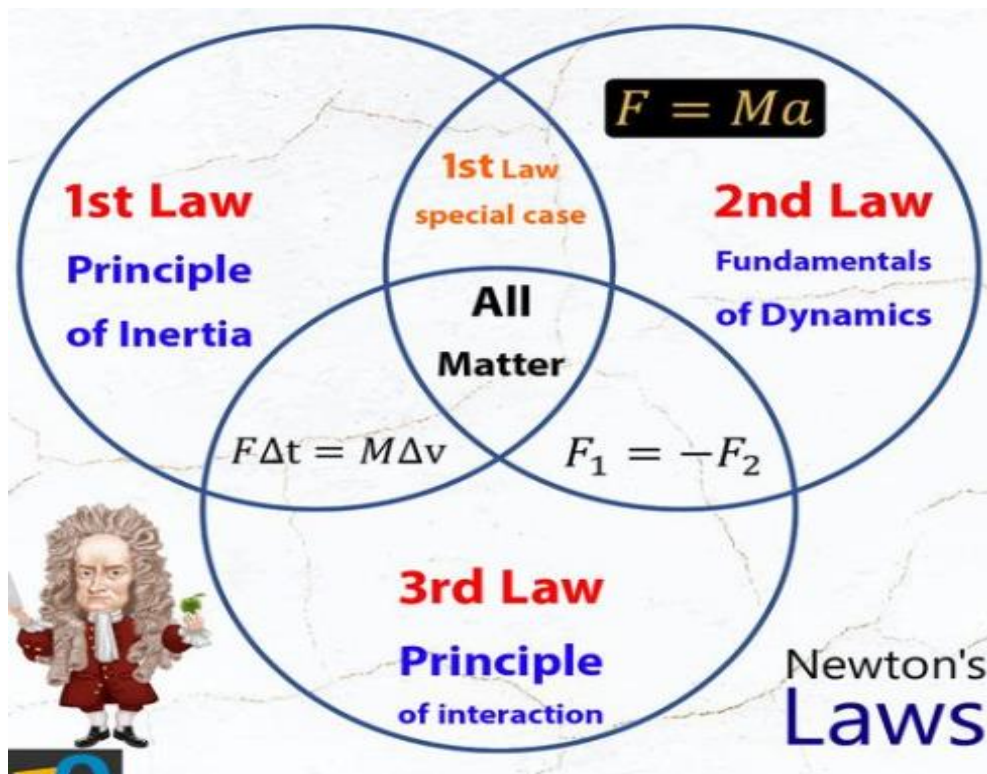
### Coriolis acceleration

$$\vec{a}_c = 2\vec{\omega} \Lambda \vec{v}_r$$

$$\Rightarrow \vec{a}_c = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ 0 & 0 & \omega \\ -R\omega \sin \omega t & R\omega \cos \omega t & 0 \end{vmatrix} = -R \omega^2 (\cos \omega t \vec{i} + \sin \omega t \vec{j})$$

# COURSE OF MECHANICS OF THE MATERIAL POINT

## *Chapter V: Dynamics of a particle*



## Glossary

In English	In French	In Arabic
The momentum	La quantité du mouvement	كمية الحركة
Isolated system	Un système isolé	جملة معزولة
The principle of inertia	Principe d'inertie	مبدأ العطالة
A free and isolated particle	Une particule libre et isolée	جسم معزول او حر
Fundamental Principle of Dynamics	Principe fondamental de la dynamique	المبدأ الأساسي للتحريك
Newtonian mechanics	La mécanique newtonienne	ميكانيك نيوتن
Principle of action and reaction	Principe d'action et de réaction	مبدأ الفعل و رد الفعل
Force of gravity-weight	La force de gravitation	الثقل او قوة الجاذبية
Force at a distance	La force à distance	القوة عن بعد
Force électrique	La force électrique	القوة الكهربائية
Binding or contact forces	La force de réaction	قوة رد الفعل
Equilibrium	Equilibre	حالة التوازن
Friction forces	La force de frottement	قوة الاحتكاك
Static friction force	La force de frottement statique	قوة الاحتكاك في حالة السكون
The coefficient of static friction	Le coefficient de frottement statique	معامل الاحتكاك في حالة السكون
The coefficient of dynamic friction	Le coefficient de frottement dynamique	معامل الاحتكاك في حالة الحركة
Elastic forces	La force élastique	قوة الارجاع او قوة المرونية
The spring stiffness constant	La constante de raideur d'un ressort	ثابت المرونة لناقض

### 1. Introduction

In physics, dynamics is the science that studies the relationship between a body in motion and the causes of that motion. It also predicts the motion of a body located in a given environment. Dynamics, more precisely, is the analysis of the relationship between applied force and changes in body motion.

### 2. Newton's laws of motion

#### 2.1. The momentum كمية الحركة

The momentum of a particle is the product of its mass and its instantaneous velocity vector.

$$\vec{P} = m\vec{v}$$

Experiments have shown that the momentum of a system composed of two particles, subject only to their mutual influences, remains constant.

#### Theorem:

"In an isolated system of two particles, the variation in the momentum of one particle over time is equal to and opposite in direction to the variation in the momentum of the other particle over the same time".

#### 2.2. Newton's three laws

##### 2.2.1. Galilean principle of inertia مبدأ العطالة (Newton's first laws)

Newton's first law, also known as the law of inertia, states that any object continues to move at a constant speed in a straight line, or remains at rest, unless an external force is applied to it. In other words, if the material body is not subjected to any force, it is either in uniform rectilinear motion, or at rest, if it was initially at rest.

For a particle the principle of inertia thus states: "A free and isolated particle moves in rectilinear motion with constant velocity".

**Note:** A free particle always moves with a constant momentum (principle of inertia).

##### 2.2.2. Newton's second law (Fundamental Principle of Dynamics) المبدأ الأساسي

##### للتحرك

In an abstract sense, force represents the effort required to modify a body's state of motion, in particular to modify its speed. Different bodies have different inertia, i.e. different resistance to a change in their state of inertia, and therefore different resistance to a change in their state



## Chapter V: Dynamics of a particle

of motion. This property must therefore be taken into account in the definition of force. To this end, we introduce a new physical quantity called the momentum of a body.

$\vec{P} = m\vec{v}$ . Consequently, force can be defined by the derivative of momentum P. This means that the resultant of the forces applied to a particle is:

$$\vec{F} = \frac{d\vec{p}}{dt}$$

This equation is called the "equation of motion»

$$\vec{F} = \frac{dm\vec{v}}{dt} \Rightarrow \vec{F} = m \frac{d\vec{v}}{dt} + \vec{v} \frac{dm}{dt} \overset{=0}{=} m \frac{d\vec{v}}{dt}$$

So,  $\vec{F} = m\vec{a}$

This is because the mass m of the moving particle is constant (as is often the case in Newtonian mechanics).

In general, Newton's second law for a moving particle can be written as:

$$\sum \vec{F}_{ext} = m\vec{a}$$

In the S.I. system, the unit of force is:  $1\text{Newton}=1\text{N}=1\text{kg.m.s}^{-2}$ .

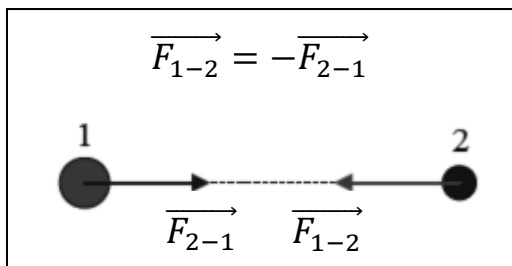
### Statement of the Fundamental Principle of Dynamics (2<sup>nd</sup> Newton law) :

In a Galilean frame of reference, the sum of the external forces applied to a system is equal to the derivative of the momentum vector of the system's center of inertia.

#### 2.2.3. Newton's third law or principle of action and reaction مبدأ الفعل ورد الفعل

Newton's third law, often referred to as the law of action and reaction, states that for every action, there is an equal and opposite reaction. In other words, when two particles are under mutual influence, the force applied by the first particle on the second is equal to, and opposite in sign to, the force applied by the second particle on the first.

This is shown in the following figure, which allows us to write:



$$|\vec{F}_{1-2}| = |\vec{F}_{2-1}|$$

### 3. Notion of force and law of force

The definition of force by the equation  $\vec{F} = m\vec{a}$  allows us to express the force corresponding to the effect studied as a function of physical factors such as distance, mass, electric charge of the bodies.... We will ultimately arrive at deriving "the law of force".

This law clearly shows the expression of the force (the resultant) applied to a material point in a well-defined situation.

#### 3.1. Force of gravity "weight $\vec{p}$ " الثقل او قوة الجاذبية

It's gravitation that makes all the bodies in the universe attract each other. It's an attractive, long-range, low-amplitude force. The gravitational phenomenon is created by the interaction between two bodies. The force of gravity acting on a human being when on Earth is the result of the interaction between the earth and the human body. As the Earth is more imposing, the gravitational force pulls the human body towards the center of the Earth. This is gravity.

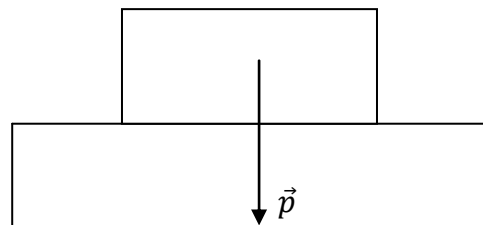
Mass (m) is the total amount of matter that makes up an object, while weight (p) is the result of the force of gravity (g) on mass. The mathematical formula is as follows :

$$p = m \times g.$$

The gravity field is represented at any point on the globe by the vector:  $\vec{p} = m\vec{g}$ .

With  $\vec{g}$  is the gravity acceleration vector, it depends on the altitude and latitude at which the body is located. It is generally considered to be constant, and the value adopted, at mean sea level, is  $9.81 \text{ m.s}^{-2}$ .

Representation of the force of weight:  $\vec{p}$  is always vertical, and directed downwards.



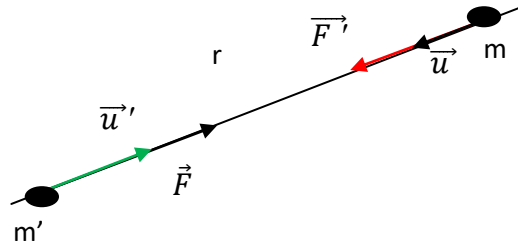
#### 3.2. Force at a distance

Assume two bodies separated by a distance r, of mass m and m' respectively.

The attractive force exerted by m on m' is :  $\vec{F} = \overrightarrow{F_{m'/m}} = -G \frac{mm'}{r^2} \vec{u}$

The attractive force exerted by m' on m is :  $\vec{F}' = \overrightarrow{F_{m/m'}} = G \frac{mm'}{r^2} \vec{u}$

Then:  $\vec{F} = -\vec{F}'$ .



where G is a constant, the value of which is experimentally determined to be:

$$G = 6.673 \times 10^{-11} \text{ N.m}^2.\text{kg}^{-2}.$$

### 3.3. Force électrique القوة الكهربائية

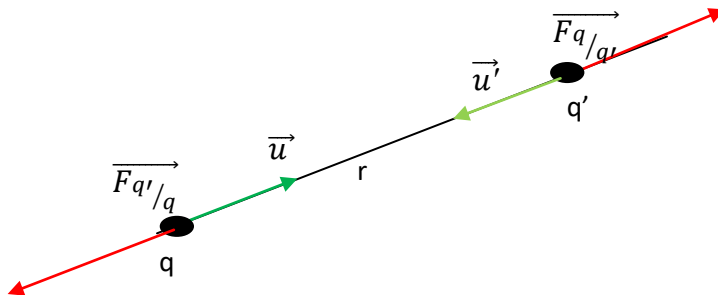
Consider two electric charges q and q' separated by a distance r. The electric force exerted by q on q' is given by:

$$\vec{F}_{q/q'} = k \frac{qq'}{r^2} \vec{u}$$

With : k a constant

The electric force exerted by q' on q is given by :

$$\vec{F}_{q'/q} = k \frac{qq'}{r^2} \vec{u}'$$



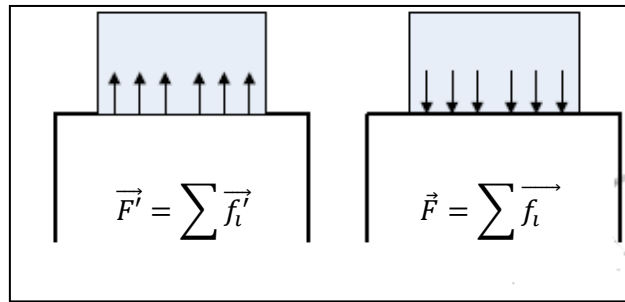
### 3.4. Binding or contact forces قوة رد الفعل

These are the forces acting mutually between bodies in contact.

Consider a solid body placed on a table. The body is in equilibrium on the table, i.e. the acceleration is zero ( $\vec{a} = \vec{0}$ ).

Faced with the force  $\vec{F}$ , representing the resultant of all the interactions of the molecules making up the body, and applied to the table, the latter in turn applies the force  $\vec{F}'$  which is the resultant of all the interactions of the molecules making up the surface of the table that is in contact with the body. The two forces  $\vec{F}$  and  $\vec{F}'$  are called contact or binding forces because of the contact between the two bodies.

With  $\vec{F} = -\vec{F}'$  and  $|\vec{F}| = |\vec{F}'|$ .



**3.5. Friction forces** القوة الاحتكاك

Whenever there is contact between the rough surfaces of two solid bodies, a resistance arises that opposes the relative movement of the two bodies. This resistance is called frictional force.

Friction is influenced by a number of factors. Consider the type of surface in contact. Smooth surfaces generally offer less friction than rough ones. Friction between solid bodies can be both static and dynamic.

**a- Static friction force** القوة الاحتكاك في حالة السكون

Static friction is the force that keeps a body at rest even in the presence of an external force.

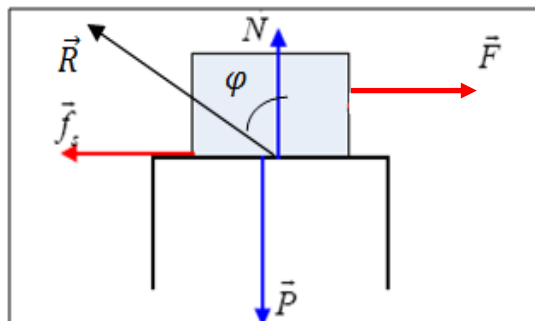
**Example:**

A body resting on a horizontal plane:

Consider the body shown in the figure below. It is subjected to four forces.

Let  $f_s$  , be the static friction force and  $\vec{P}$  and  $\vec{N}$  be the weight and normal reaction force of the support respectively.

For the body on the table to move, a minimum force  $\vec{F}$  must be applied.



The mass remains stationary as long as  $F < f_s$ , there is resistance to movement.

In this case the reaction of the support is the resultant force given by:  $\vec{R} = \vec{N} + \vec{f}_s$

At equilibrium:

$$\sum \vec{F}_{ext} = \vec{0} \Rightarrow \vec{N} + \vec{f}_s + \vec{P} + \vec{F} = \vec{0}$$
$$\Rightarrow \vec{R} + \vec{P} + \vec{F} = \vec{0}$$

By projecting onto the two axes Ox and Oy:

On (Oy) :  $P=N$  and (Ox) :  $F=f_s$

The mass starts moving when  $F > f_s$ .

Experience shows that the ratio ( $f_s/N$ ) is constant.

$$tg\varphi = \frac{f}{N} = k = \mu$$

$\mu$ : is the coefficient of friction and  $\varphi$  is the angle of friction.

The coefficient of friction is called static when the body is stationary. The coefficient of static friction is a ratio between the static frictional force of an object and the normal force, and is written as follows :

$$tg\varphi = \frac{f_s}{N} = \mu_s$$

### **b- Dynamic friction force** القوة الاحتكاك في حالة الحركة

Kinetic or dynamic friction is the frictional force present when an object is in motion on another object.

The dynamic friction coefficient is a ratio between the dynamic friction force of an object and the normal force.

Mass starts moving when  $F > f_d$ .

The coefficient of dynamic friction is written as :

$$tg\varphi = \frac{f_d}{N} = \mu_d$$

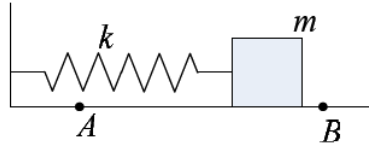
**Note :**

Experience has shown that the coefficient of static friction is greater than the coefficient of dynamic friction  $\mu_s > \mu_d$ .

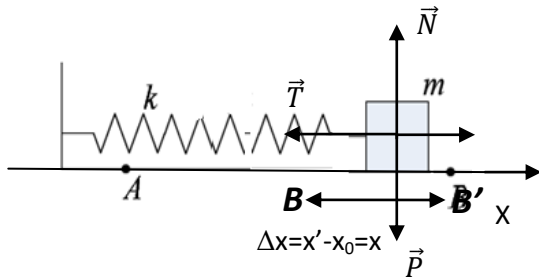
## Chapter V: Dynamics of a particle

### 3.6. Elastic forces قوة الارجاع او قوة المرونية

Elastic force is the force applied to an object that tends to return to its shape after being deformed. Elastic forces cause periodic movements. The most common is sinusoidal motion, as in the case of a spring.



(At equilibrium)



(Motion)

We have: FPD:  $\sum \vec{F}_{ext} = m\vec{a}$

$$\vec{N} + \vec{P} + \vec{T} = m\vec{a}$$

M is the mass of the body

N is the reaction force of the support قوة رد الفعل

P is the force of the body weight قوة الثقل

T is the spring return force قوة الارجاع

By projection:

On (Ox) we have:  $T = -kx = ma$

On (Oy) we have:  $N - P = 0$  So  $N = P = mg$

Where k is the spring stiffness constant ثابت المرونة.

### Proposed exercises about chapter V

#### Exercise 1

Consider a small block of mass  $m$  abandoned without initial velocity at point **A** of an inclined plane at an angle  $\alpha=30^\circ$  to the horizontal. Point **A** is at height  $h$ .

1- What is the value of the coefficient of static friction  $\mu_s$  that keeps the mass in equilibrium at point **A**.

#### Exercise 2

A man pushes a  $20\text{ kg}$  lawnmower with a force of  $80\text{N}$  directed parallel to the handle, which is inclined at  $30^\circ$  to the horizontal.

1. If moving at constant speed, what is the modulus of the friction force due to the ground?
2. What force parallel to the handle would produce an acceleration of  $1\text{m/s}^2$ , given that the friction force is that found in question 1?

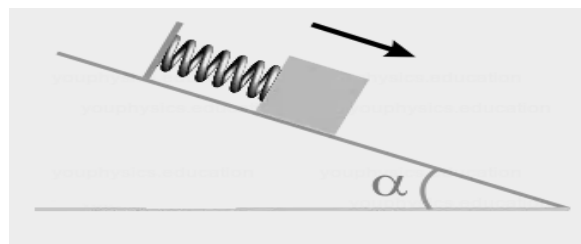
#### Exercise 3

A block of mass  $m$  ascends along a plane inclined by an angle  $\alpha$ , with respect to the horizontal, with initial velocity  $v_0$ , and coefficient of friction  $\mu_d$ .

1. Determine how far the block travels before coming to rest.
2. What is the maximum value that the static friction coefficient  $\mu_s$  can take for the body to remain stationary.
3. For a value of the dynamic friction coefficient  $\mu_d$  lower than the maximum value found in the second question, what is the velocity  $v_1$  of the body when it returns to its starting position.

#### Exercise 4

A mass  $m = 15\text{ kg}$  suspended from a spring of stiffness  $K = 100\text{N/m}$  descends along an inclined plane which makes an angle  $\alpha = 30^\circ$  with the horizontal.



## Chapter V: Dynamics of a particle

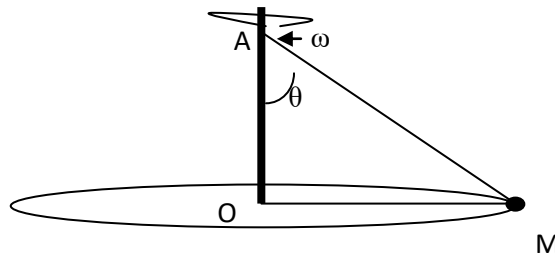
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Assuming there is no friction, determine the normal reaction of the support and the acceleration of the mass when the spring is stretched by a length  $x = 0.02\text{m}$ .

### Exercise 5

A ball of mass  $m$  is attached by two wires ( $Am$  and  $Om$ ) to a vertical pole. The whole system rotates with a constant angular velocity  $\omega$  around the axis of the post (we know  $g$  the acceleration of gravity,  $\theta$  and  $L = |\overline{OM}|$ )

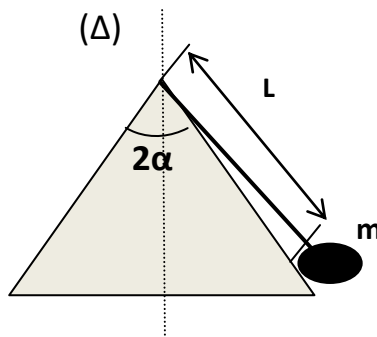
1. Assuming  $\omega$  is large enough to keep both wires taut, find the force (wire tension) each wire exerts on the ball.
2. What is the minimum angular velocity  $\omega_{\min}$  for which the bottom wire remains taut?



### Exercise 6

A body of mass ( $m=1\text{kg}$ ) is attached by a wire of length  $L=30\text{cm}$  to the top of a cone, of axis ( $\Delta$ ) and angle at the top  $2\alpha=60^\circ$ . This body rotates without friction on the surface of the cone with a rotational speed  $\omega=10\text{ rpm.}(10.2\pi/60\text{s})$

1. Calculate the body's linear velocity.
2. Using the fundamental principle of dynamics, determine the reaction ( $R_N$ ) of the cone surface on the body and the thread tension ( $T$ ).



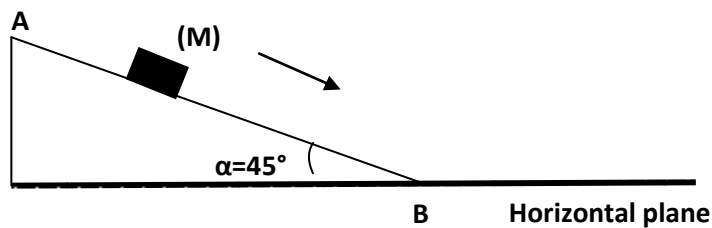


## Chapter V: Dynamics of a particle

### Exercise 7

A block (M) of mass  $m$  is thrown from the top of an inclined plane  $AB=1\text{m}$  at an angle  $\alpha=45^\circ$  to the horizontal, with initial velocity  $v_A=1\text{m/s}$ .

- 1- Knowing that the coefficient of friction  $\mu=0.5$  on AB.
  - Demonstrate, what is the nature of the motion on AB?
  - Calculate the speed of (M) when it reaches point B.
- 2- Friction forces are considered negligible on the horizontal plane:
  - Demonstrate the nature of the motion on the horizontal plane.
  - Will the block (M) stop? Justify your answer.

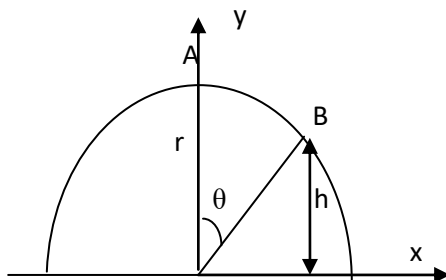


### Exercise 8

A piece of ice M of mass  $m$  slides frictionlessly over the outer surface of an igloo, which is a half-sphere of radius  $r$  with a horizontal base.

At  $t=0$ , it is released from point A without any initial velocity.

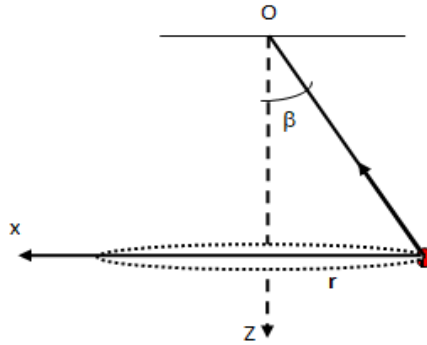
- 1- Find the expression for the velocity at point B, as a function of  $g$ ,  $r$  and  $\theta$ .
- 2- Using the fundamental relation of dynamics, determine the expression of  $|\vec{N}|$  the reaction of the igloo on M at point B as a function of velocity  $v_B$ .
- 3- At what height does M leave the sphere?
- 4- At what speed does M arrive at the axis (Ox)?



### Exercise 9

A material point, of mass  $m$ , is suspended at a fixed point  $O$  by a wire of length  $l$  inextensible and negligible mass. By rotating it around axis  $Oz$ , it acquires a constant angular velocity  $\omega$ . It describes a horizontal circle of radius  $r$ .

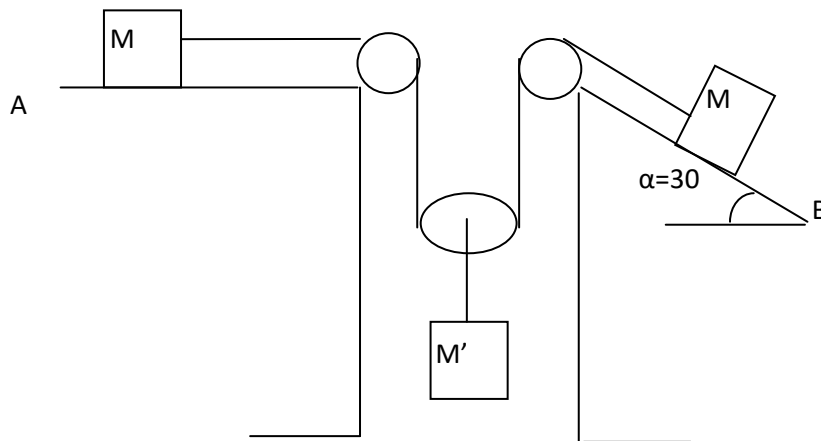
1. Find the expression for the wire tension.
2. Find the expression for the inclination  $\beta$  of the wire with respect to the vertical.



### Exercise 10

Two carriages A and B of the same mass  $M$  are linked by a wire carrying a pulley of negligible mass. The axis of the pulley carries a mass  $M'$ .

- 1- Neglecting all friction, calculate the ratio  $M'/M$  so that cart B remains stationary.
- 2- If  $M'=2M$ , calculate the accelerations to which the masses are subjected.



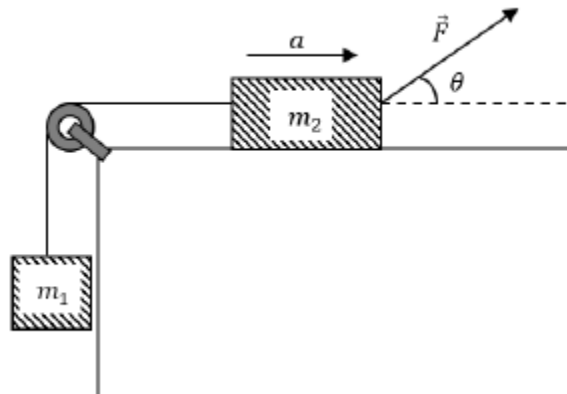
### Exercise 11

A block of mass  $m_1$  assimilated to a material point can slide on a horizontal surface with a coefficient of kinetic friction  $\mu_d$  one of its ends is connected by an inextensible wire of negligible mass passing through a pulley of negligible mass connected to a second mass  $m_2$ . A

## Chapter V: Dynamics of a particle

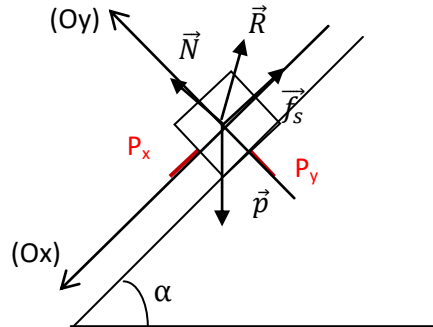
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force of modulus  $F$  is applied to  $m_2$  at an angle  $\theta$  to the horizontal. Find the accelerations of the two masses.



**Exercises correction of about chapter V**

**Exercise 1**



What is the value of the coefficient of static friction  $\mu_s$  that keeps the mass in equilibrium at point A?

-At equilibrium:

$$\sum \vec{F}_{ext} = \vec{0} \Rightarrow \vec{N} + \vec{f}_s + \vec{P} = \vec{0} \Rightarrow \vec{R} + \vec{P} = \vec{0}$$

Following (Ox):  $-f_s + p_x = 0 \Rightarrow f_s = m g \sin\alpha$

Following (Oy):  $N - p_y = 0 \Rightarrow N = m g \cos\alpha$

In order for the body to remain stationary on the plane, the following conditions must be met  $f_s > p_x$ .

$$\text{We have } \tan\varphi = \frac{f_s}{N} = \mu_s = \frac{m g \sin\alpha}{m g \cos\alpha} = \tan\alpha$$

The maximum value that the coefficient of static friction  $\mu_s$  can take is  $\tan\alpha$ .

**Note:** experience shows that:  $\mu_s \geq \mu_d$

**Exercise 2**

Modulus of friction force:

$M=20\text{kg}$ ,  $\alpha = 30^\circ$  and  $F=80\text{N}$ .

1- FPD:  $\Sigma \vec{F} = m\vec{a} \Rightarrow \vec{p} + \vec{R}_N + \vec{F} + \vec{f} = m\vec{a}$

with  $v=\text{cst}$  so  $a=0$ .

Following (Ox):  $F \cos\alpha - f = 0$

Following (Oy):  $R_N - P = 0 \Rightarrow R_N = m g$

## Chapter V: Dynamics of a particle

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Then:  $f=80.\cos.30=69.28N$

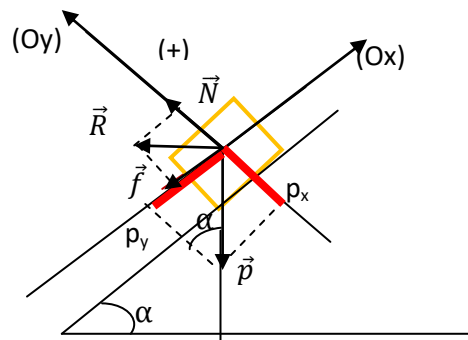
2- Force F for  $a=1m/s^2$

$$\Sigma \vec{F} = m\vec{a} \Rightarrow \vec{p} + \vec{R}_N + \vec{F}' + \vec{f} = m\vec{a}$$

Following (Ox):  $-f + F' \cos\alpha - f = m.a$

$$F' = \frac{m.a + f}{\cos\alpha} = 103.1N$$

### Exercise 3



At:  $t=0, v=v_0$  and  $\mu = \mu_d$

1- Let's find out how far the block can travel before it stops.

According to the fundamental principle of dynamics:

$$\Sigma \vec{F} = m\vec{a} \Rightarrow \vec{p} + \vec{R} = \vec{p} + \vec{N} + \vec{f} = m\vec{a}$$

Initial velocity  $v_i = v_0$  and final velocity  $v_f = 0$  (the body will stop)

We have:

$$v_f^2 - v_i^2 = 2al$$

( $l$  being the distance covered by the body)

$$\text{So ; } a = \frac{v_f^2 - v_i^2}{2l}$$

The reference frame must be chosen so that the axis (Ox) follows the axis of motion, so it is parallel to  $\vec{f}$  and (Oy) is perpendicular to (Ox), so it is parallel to  $\vec{N}$ .

Following (Ox):  $-f - p_x = -f - m g \sin\alpha = ma$

Following (Oy):  $N - p_y = 0 \Rightarrow N = m g \cos\alpha$

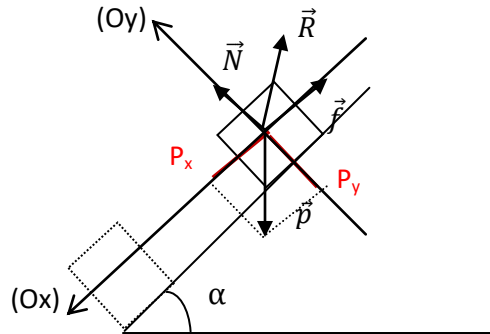
## Chapter V: Dynamics of a particle

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$$\mu_d = \tan \varphi = f/N \Rightarrow f = N \tan \varphi = N \mu_d \text{ so } f = \mu_d m g \cos \alpha$$

$$- \mu_d m g \cos \alpha - m g \sin \alpha = ma \Rightarrow - \mu_d g \cos \alpha - g \sin \alpha = \frac{v_f^2 - v_i^2}{2l}$$

$$\text{Then, } l = \frac{-v_i^2}{2(-\mu_d g \cos \alpha - g \sin \alpha)} = \frac{v_0^2}{2g(\mu_d \cos \alpha + \sin \alpha)}$$



2- The maximum value that the static friction coefficient  $f_s$  can take for the body to sink,

- At equilibrium

$$\text{Following (Ox): } -f + p_x = 0 \Rightarrow f = m g \sin \alpha$$

$$\text{Following (Oy): } N - p_y = 0 \Rightarrow N = m g \cos \alpha$$

For the body to be able to descend, it must:  $p_x > f$

$$p_x \geq f \Rightarrow m g \sin \alpha \geq N \mu_s \quad (*) \quad (\mu_s = f/N)$$

with  $f = N \mu_s$  and  $\mu_s$  is the coefficient of static friction at which the body begins its motion

(with  $\mu_s = f/N \Rightarrow f = N \tan \varphi$  so  $f = \mu_s m g \cos \alpha$ ).

$$(*) \Rightarrow m g \sin \alpha \geq m g \cos \alpha \mu_s \text{ so } \mu_s \leq \tan \alpha$$

The maximum value that can be :  $\mu_s$  is  $\tan \alpha$

3- The velocity  $v_1$  of the body as it returns to its initial position;

$x=l$ ,  $v_i=0$  and we look for  $v_f$ .

$$v_f^2 - v_i^2 = 2al$$

## Chapter V: Dynamics of a particle

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Where 'l' is the distance covered by the body.

$$\text{So, } a = \frac{v_f^2 - v_i^2}{2l}$$

The reference frame must be chosen so that the (Ox) axis follows the axis of motion, i.e. it is parallel to and follows  $p_x$ , and the (Oy) axis is perpendicular to (Ox), i.e. it is parallel to  $\vec{N}$ .

Following (Ox):  $-f + p_x = -f + m g \sin\alpha = m.a$

Following (Oy):  $N - p_y = 0 \Rightarrow N = m g \cos\alpha$

$$\mu_d = \tan\phi = f/N \Rightarrow f = N \tan\phi \quad \text{so } f = \mu_d m g \cos\alpha$$

Hence;  $-\mu_d m g \cos\alpha + m g \sin\alpha = m.a$

$$\Rightarrow -\mu_d g \cos\alpha + g \sin\alpha = \frac{v_f^2 - v_i^2}{2l}$$

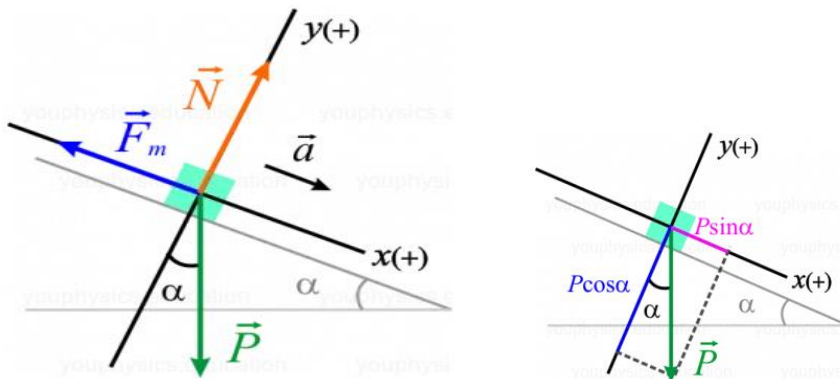
$$v_f^2 = 2gl(\sin\alpha - \mu_d \cos\alpha)$$

(where l is the same distance found in the question 1)

### Exercise 4

To solve this problem, we'll use Newton's second law. The forces acting on the mass are weight (if it's close to the Earth's surface), the normal (because it's resting on the plane) and the force of the spring, which is given by Hooke's law.

Below is a diagram of the forces acting on the mass and the Cartesian axes we'll use to make the projections.



## Chapter V: Dynamics of a particle

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The spring's restoring force acts in the opposite direction to the spring's elongation (and therefore to the direction of mass displacement):

The acceleration of the mass is also shown in the figure. It runs in the positive direction of the x axis. In the diagram below, we've plotted the projections of the weight vector on the axes we've chosen.

Newton's second law applied to mass motion gives:

$$\Sigma \vec{F} = m\vec{a} \Rightarrow \vec{p} + \vec{N} + \vec{F}_r = m\vec{a}$$

By projecting onto the Cartesian axes we obtain:

Following (Ox):  $-F_r + p_x = -F_r + m g \sin\alpha = m.a$  (1)

Following (Oy):  $N - p_y = 0 \Rightarrow N = m g \cos\alpha$  (2)

We obtain the norm of the support reaction from equation (2), and as you can see, it's not equal to the weight.

$$N = m g \cos\alpha$$

On the other hand, the norm of the spring return force is given by:  $F_r = k.x$

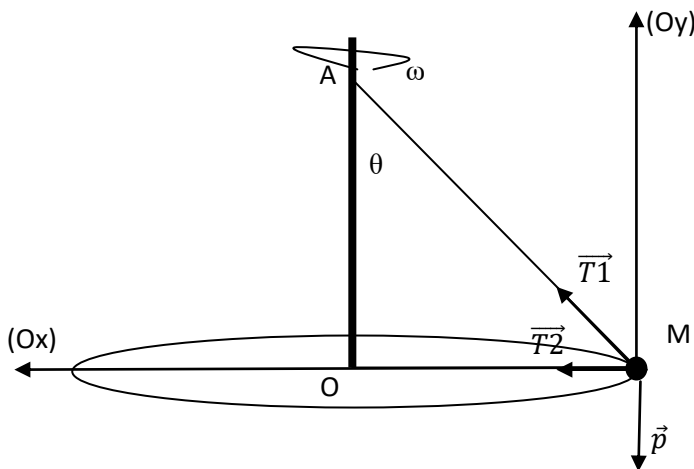
Finally, solving equation (1) yields the acceleration:

$$a = (-k.x + m g \sin\alpha) / m = 4.8 \text{ m/s}^2$$

by taking :  $g = 10 \text{ m/s}^2$ .

### Exercise 5

Calculating the tension T on the wire:



1- Let's find the force (thread tension) that each thread exerts on the ball.



## Chapter V: Dynamics of a particle

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According to the fundamental principle of dynamics FPD:

$$\Sigma \vec{F} = m\vec{a} \Rightarrow \vec{p} + \vec{T}_1 + \vec{T}_2 = m\vec{a}$$

*The ball's motion is circular, so the acceleration in this case is the normal acceleration  $a_N$ , which is directed towards the center of the circle. (with  $a_N=v^2/R$ )*

We choose the reference frame such that :

(Ox) follows the normal acceleration and is directed towards the center of the circle.

By projection onto the axes (Oy) and (Ox) we have :

$$\text{On (Ox) : } T_2 + T_1 \sin \theta = m \cdot a_N \Rightarrow T_2 + T_1 \sin \theta = m \frac{v^2}{R}$$

$$\text{On (Oy) : } p - T_1 \cos \theta = 0 \Rightarrow mg = T_1 \cos \theta$$

$$\text{So } T_1 = \frac{mg}{\cos \theta}$$

$$T_2 = m \frac{v^2}{R} - T_1 \sin \theta = m \frac{v^2}{R} - \frac{mg}{\cos \theta} \sin \theta$$

$$\text{So } T_2 = m \frac{v^2}{R} - m g \operatorname{tg} \theta$$

2- The minimum angular speed  $\omega_{\min}$  at which the bottom wire remains taut.

In order for the lower wire to remain taut, the following conditions must be met  $T_2 > 0$ .

$$T_2 = m \frac{v^2}{R} - m g \operatorname{tg} \theta \geq 0 \Rightarrow \frac{v^2}{R} \geq g \operatorname{tg} \theta$$

$$\text{With, } v = \omega \cdot R \Rightarrow \frac{\omega^2 R^2}{R} \geq g \operatorname{tg} \theta \text{ and } R = OM = L$$

$$\text{So, } \omega^2 L \geq g \operatorname{tg} \theta \Rightarrow \omega^2 \geq \frac{g \operatorname{tg} \theta}{L}$$

$$\text{And, } \omega \geq \sqrt{\frac{g \operatorname{tg} \theta}{L}} \quad \text{Then } \omega_{\min} = \sqrt{\frac{g \operatorname{tg} \theta}{L}}$$

### Exercise 6

1- The linear velocity of the body.

$L=30\text{cm}$ ,  $2\alpha=60^\circ$ . ( $\alpha=30^\circ$ ) and  $\omega=10 \text{ tr/mn}$ .

## Chapter V: Dynamics of a particle

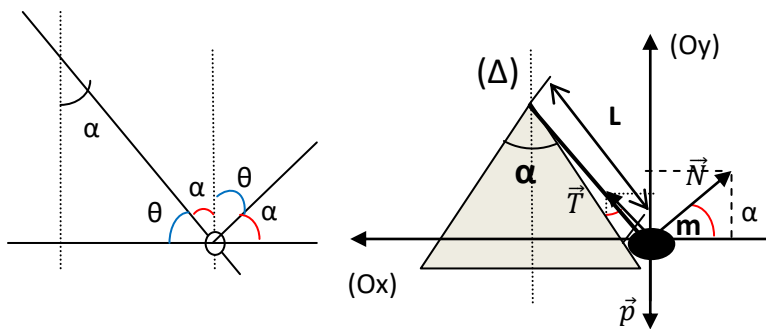
$$\begin{cases} 10 \times 2\pi \longrightarrow 60 \text{ s} \\ \omega \longrightarrow 1 \text{ s} \end{cases} \Rightarrow \omega = \frac{10 \times 2\pi}{60} = \frac{\pi}{3} \text{ rad/s}$$

$$v = \omega R \text{ and } R = l \sin \alpha$$

$$\text{Hence } v = \omega l \sin \alpha = \frac{\pi}{3} \cdot 0,3 \cdot \sin 30 = 0,157 \text{ m/s}$$

Let's determine the reaction (N) of the surface of the cone on the body and the tension of the wire (T).

According to the fundamental principle of dynamics FPD.



$$\Sigma \vec{F} = m\vec{a} \Rightarrow \vec{p} + \vec{N} + \vec{T} = m\vec{a}_N$$

We choose a reference frame such that (Ox) follows the normal acceleration and is directed towards the center of the cone, and axis (Oy) is perpendicular to (N).

$$\text{Following (Ox): } T_x - N_x = m a_N \quad (1)$$

$$\text{Following (Oy): } T_y + N_y - p = m a_T = 0 \quad (a_T = 0, \text{ because the speed is constant})$$

$$\Rightarrow T \cos \alpha + N \sin \alpha - p = 0 \quad (2)$$

$$(1) \Rightarrow T \sin \alpha - N \cos \alpha = m \frac{v^2}{R} = m \frac{\omega^2 R^2}{R}$$

$$\text{So ; } T = \frac{m\omega^2 R}{\sin \alpha} + N \frac{\cos \alpha}{\sin \alpha}$$

We replace it in the second equation:

$$\left( \frac{m\omega^2 R}{\sin \alpha} + N \frac{\cos \alpha}{\sin \alpha} \right) \cos \alpha + N \sin \alpha - p = 0$$

## Chapter V: Dynamics of a particle

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$$\left( \frac{m\omega^2 R}{\sin \alpha} \cos \alpha + N \frac{\cos^2 \alpha}{\sin \alpha} \right) + N \sin \alpha - p = 0$$

$$N \left( \frac{\cos^2 \alpha}{\sin \alpha} + \sin \alpha \right) = p - \frac{m\omega^2 R}{\sin \alpha} \cos \alpha$$

$$\Rightarrow N \left( \frac{1}{\sin \alpha} \right) = \frac{mg \sin \alpha - m\omega^2 R \cos \alpha}{\sin \alpha}$$

So,  $N = m(g \sin \alpha - \omega^2 R \cos \alpha)$ , Replacing  $R$  with  $l \sin \alpha$ , we'll have:

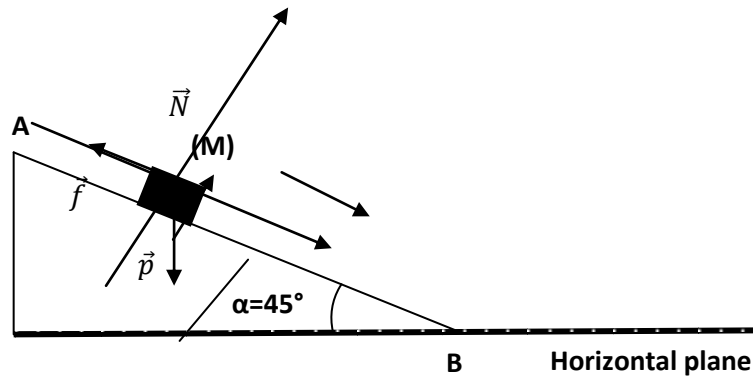
$$N = m(g \sin \alpha - \omega^2 l \sin \alpha \cos \alpha) = 7,92 \text{ N}$$

$$\text{Hence, } T = \frac{m\omega^2 l \sin \alpha}{\sin \alpha} + N \frac{\cos \alpha}{\sin \alpha} = 5,88 \text{ N}$$

If we replace  $N$  by its expression, we find:  $T = m g \cos \alpha + m \cdot \omega^2 l (1 - \cos^2 \alpha)$ .

$$\text{Hence; } T = m (g \cos \alpha + \omega^2 l \sin^2 \alpha)$$

### Exercise 7



$v_A = 1 \text{ m/s}$  and  $\mu = 0,5$  on AB.

The nature of movement on AB: FPD:

$$\Sigma \vec{F} = m\vec{a} \Rightarrow \vec{p} + \vec{N} + \vec{f} = m\vec{a}$$

We choose the reference frame, such that axis (Ox) is along the axis of motion parallel to  $\vec{f}$  and (Oy) is perpendicular to (Ox) therefore along  $\vec{N}$ .

Following (Ox):  $-f + p_x = -f + m g \sin \alpha = ma$

Following (Oy):  $N - p_y = 0 \Rightarrow N = m g \cos \alpha$

## Chapter V: Dynamics of a particle

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$$\mu = \tan \varphi = f/N \Rightarrow f = N \tan \varphi \quad \text{So } f = \mu m g \cos \alpha$$

$$\text{Hence; } -\mu \cdot m g \cdot \cos \alpha + m g \sin \alpha = ma \Rightarrow a = g(\sin \alpha - \mu \cos \alpha)$$

$$a = 10 \left( \frac{\sqrt{2}}{2} - 0,5 \frac{\sqrt{2}}{2} \right) = 3,54 \text{ m/s}^2$$

The acceleration  $a$  is constant and positive, so the motion is uniformly accelerated.

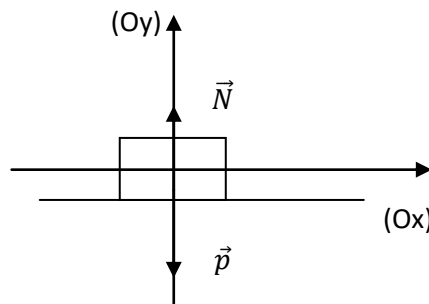
The speed of point M when it reaches point B.

$$v_B^2 - v_A^2 = 2al \Rightarrow v_B^2 = v_i^2 + 2al$$

With ;  $l = AB = 1$

$$v_B = \sqrt{1 + 2a} = 2,84 \text{ m/s}$$

The nature of movement on the horizontal plane: Friction forces are negligible.



$$\Sigma \vec{F} = m\vec{a} \Rightarrow \vec{p} + \vec{N} = m\vec{a}$$

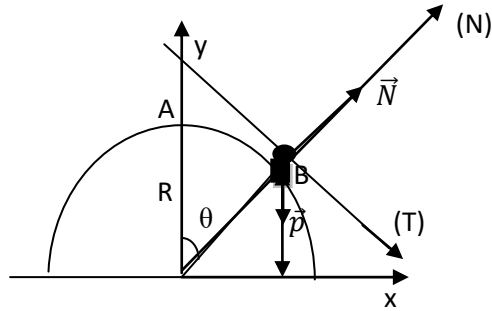
Following (Ox):  $0 = ma'$

Following (Oy):  $N - p = 0 \Rightarrow N = p = mg$

So  $a' = 0$  then the motion is uniformly rectilinear.

Motion is uniform, so speed is constant  $v = v_B$  the block will not stop.

**Exercise 8**



The point at which the point leaves the sphere

$$\Sigma \vec{F} = m\vec{a} \Rightarrow \vec{p} + \vec{N} + \vec{f} = m\vec{a}$$

We choose the reference frame, such that axis (T) is tangent to the sphere and axis (N) is perpendicular to (T).

Following (N):  $N - p_N = -m \cdot a_N \Rightarrow N - m g \cos \theta = -m v^2 / R$  (\*\*)

Following (T):  $p_T = m \cdot a_T \Rightarrow m g \sin \theta = m (dv/dt)$  (\*)

$$v = R\omega \Rightarrow \frac{dv}{dt} = R \frac{d\omega}{dt} = R \frac{d^2\theta}{dt^2} \quad \text{with } \omega = \frac{d\theta}{dt} = \dot{\theta} \text{ and } a_T = R\ddot{\theta} = R \frac{d\dot{\theta}}{dt}$$

Equation (\*) is multiplied by  $\dot{\theta}$ :

$$\Rightarrow \dot{\theta} g \sin \theta = \dot{\theta} R \frac{d\dot{\theta}}{dt} \quad \text{so} \quad \frac{d\dot{\theta}}{dt} g \sin \theta = R \frac{d\dot{\theta}}{dt}$$

$$\text{then } d\dot{\theta} g \sin \theta = \dot{\theta} R d\dot{\theta} \Rightarrow g \int_0^{\dot{\theta}} \sin \theta d\theta = R \int_0^{\dot{\theta}} \dot{\theta} d\dot{\theta}$$

$$g(1 - \cos \theta) = R \frac{\dot{\theta}^2}{2} \Rightarrow 2g(1 - \cos \theta) = \frac{v^2}{R} \quad \text{because } v^2 = R^2 \dot{\theta}^2$$

$$\Rightarrow v^2 = 2(gR - gR \cos \theta)$$

The speed at point B is :  $v_B = \sqrt{2gR(1 - \cos \theta)}$

1- The expression for the reaction of the igloo on M

$$(**) \Rightarrow N = m g \cos \theta - m a_N = m (g \cos \theta - 2g(1 - \cos \theta))$$

$$\Rightarrow N = m (3 g \cos \theta - 2g)$$

## Chapter V: Dynamics of a particle

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1- The material point leaves the half-sphere at point p, so  $N=0$

$$N=0 \Rightarrow mg(3 \cos \theta_0 - 2) = 0 \text{ so } \cos \theta_0 = \frac{2}{3}$$

Then,  $\theta_0 = 48^\circ$

The angle relative to the horizontal at which the point leaves the half-sphere is  $90-48=52$

The height  $h$  at which the material point leaves the half-sphere is:

$$h_p = R \cos \theta = \frac{2}{3} R$$

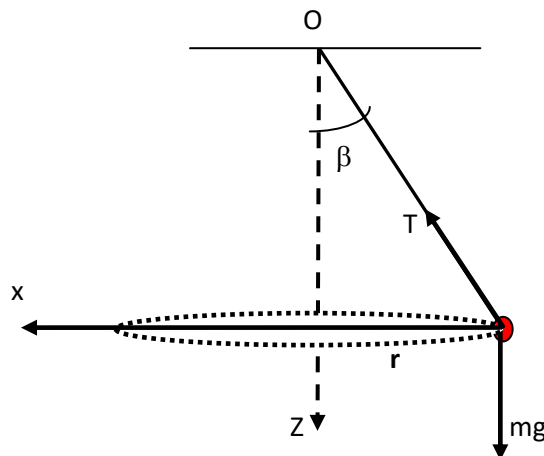
1- The velocity of the material point at this point :

$$2g(1 - \cos \theta_0) = \frac{v_p^2}{R} \Rightarrow v_p^2 = 2gR(1 - \cos \theta_0)$$

$$\text{then } v_p = \sqrt{2gR(1 - \cos \theta_0)} = \sqrt{\frac{2}{3} Rg}$$

(We'll solve the same exercise in the next chapter using the principle of conservation of energy).

### Exercise 9



1. Calculating the tension  $T$  on the wire:

The material point is subjected to two mechanical forces; weight  $mg$  and tension  $T$ .

Expressing Newton's second law, we write:

## Chapter V: Dynamics of a particle

---

$$\Sigma \vec{F} = m\vec{a} \Rightarrow \vec{p} + \vec{T} = m\vec{a}$$

In such a rotational movement of  $m$  around  $Oz$  at the angular velocity  $\omega = \text{constant}$ , the acceleration admits a tangential component and a component normal to the circular trajectory of radius  $r$  around  $OZ$ . Thus, in the radial direction ( $ox$  axis), the projection of Newton's second law gives:

$$T \sin\beta = m \cdot a_N = m(v^2/r)$$

$$\Rightarrow m \cdot a_N = m \cdot \omega^2 r = m \cdot \omega^2 L \sin\beta$$

$$\text{Since } T = m\omega^2 L$$

2. Calculation of wire inclination to the vertical.

Newton's second law is always written:

$$\Sigma \vec{F} = m\vec{a} \Rightarrow \vec{p} + \vec{T} = m\vec{a}.$$

The projection of this relationship along the vertical translates into :

$$- T \cos\beta + mg = 0$$

(because the mass rotates and does not move along  $Oz$ ).

Taking into account the expression for  $T$ , we obtain:

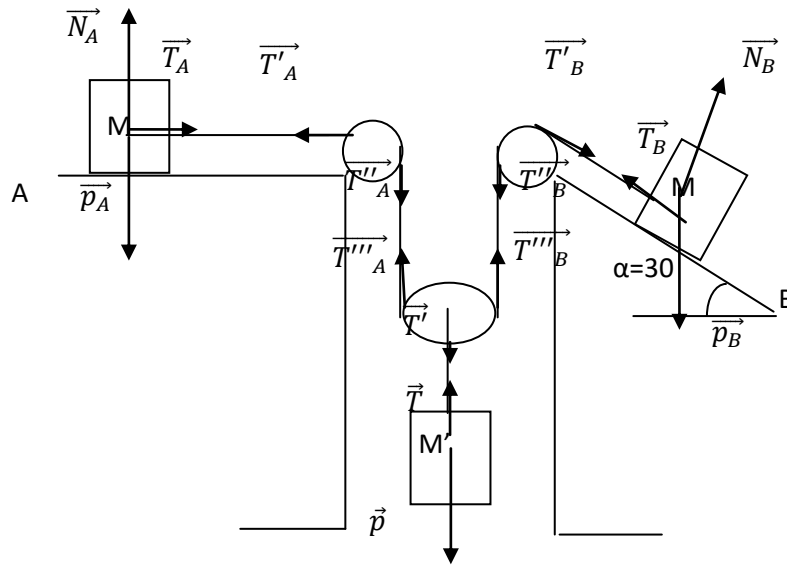
$$\cos\beta = \frac{g}{\omega^2 L}$$

Since this angle is inversely proportional to the angular velocity, which must be minimal, and since :  $0 < \beta < 90^\circ$  :

i.e.  $0 < \cos\beta < 1$ , the limit on angular velocity is defined

$$\omega \geq \sqrt{g/L}$$

Exercise 10



1- Neglecting friction, let's calculate the ratio  $M'/M$  so that cart B remains stationary?

For system A (carriage of mass M):

$$\Sigma \vec{F} = M \vec{a} \Rightarrow \vec{p}_A + \vec{N}_A + \vec{T}_A = M \vec{a}_A$$

Projection on the axis of motion

$$T_A = M a_A \quad (*)$$

For system B (carriage of mass M):

$$\Sigma \vec{F} = M \vec{a} \Rightarrow \vec{p}_B + \vec{N}_B + \vec{T}_B = \vec{0} \quad (\text{because carriage B is stationary})$$

Projection on the axis of motion.

$$-M g \sin \alpha + T_B = 0 \quad (*')$$

It's the same wire, so  $T_B = T_A = T$

$M'$  system ( $M'$  mass carriage):

$$\Sigma \vec{F} = M' \vec{a}' \Rightarrow \vec{p}' + \vec{T}' = M' \vec{a}'$$

Projection on the axis of motion, with  $T' = T_A + T_B = 2T$

$$M' g - 2T = M' a' \quad (*'')$$

When carriage A moves a distance  $x_A$ , carriage  $M'$  moves back a distance  $x'$  with  $x' = x_A/2$ .



$$x' = \frac{x_A}{2} \Rightarrow v' = \frac{v_A}{2} \quad \text{so} \quad a' = \frac{a_A}{2}$$

$$(**) \Rightarrow M'g - 2T = M' \frac{a_A}{2}$$

$$(*) \Rightarrow T = M a_A \Rightarrow a_A' = \left( \frac{2gM'}{4M+M'} \right)$$

$$(***) \Rightarrow T = M g \sin 30 = M g/2 = M a_A$$

$$\text{So } \left( \frac{2gM'}{4M+M'} \right) = \frac{g}{2} \Rightarrow 4M' = 4M + M'$$

$$\text{Then } \frac{M'}{M} = \frac{4}{3}$$

1- Let's calculate the accelerations for  $M'=2M$

The cart A :  $T = M a_A$  (\*)

The cart B :  $T - Mg \sin 30 = M a_B \Rightarrow T - \frac{Mg}{2} = M a_B$  (\*\*)

The mass cart  $M'$  ( $M'=2M$ ):  $2 M g - 2 T = 2 M a'$

When carriage A moves a distance  $x_A$  and carriage B moves a distance  $x_B$ , the mass  $M'$  moves down a distance  $x' = (x_A + x_B)/2$

so  $2 x' = (x_A + x_B) \Rightarrow 2 v' = (v_A + v_B)$  then  $2 a' = (a_A + a_B)$  (\*\*\*)

and  $2 M g - 2 T = M (a_A + a_B)$  (\*\*\*\*)

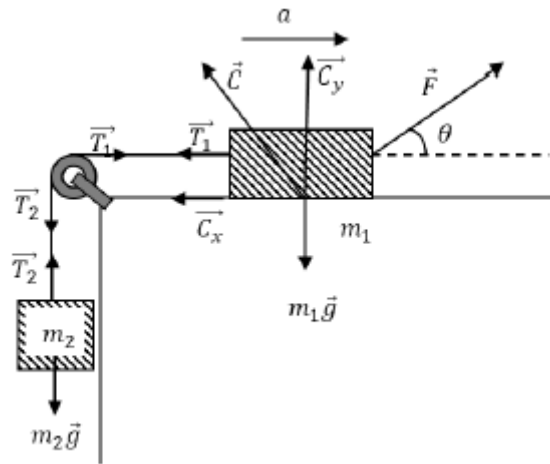
$$(*) \text{ and } (**) \Rightarrow a_A - \frac{g}{2} = a_B$$

$$(*) \text{ and } (***) \Rightarrow 2 g - 2 a_A = (a_A + a_B) \text{ so } 4 a_A = g(5/2)$$

$$\text{So } \begin{cases} a_A = g \left( \frac{5}{8} \right) \\ a_B = g \left( \frac{1}{8} \right) \\ a' = g \left( \frac{3}{8} \right) \end{cases}$$

### Exercise 11

The fundamental principle of dynamics is applied to the masses  $m_1$  and  $m_2$  :



For system  $m_1$  :

$$\Sigma \vec{F} = m_1 \vec{a} \Rightarrow \vec{p}_1 + \vec{C}_1 + \vec{T}_1 + \vec{F} = m_1 \vec{a} \quad (1)$$

For system  $m_2$  :

$$\Sigma \vec{F} = m_2 \vec{a} \Rightarrow \vec{p}_2 + \vec{T}_2 = m_2 \vec{a} \quad (2)$$

The pulley with negligible mass, so  $T_1 = T_2 = T$

Projecting equations (1) and (2) onto the direction of motion gives:

$$-C_x + F \cos \theta - T_1 = m_1 \cdot a \quad (3)$$

and

$$T_2 - m_2 g = m_2 \cdot a \quad (4)$$

Projecting equation (1) onto the direction perpendicular to the motion gives:

$$C_y = m_1 g$$

And

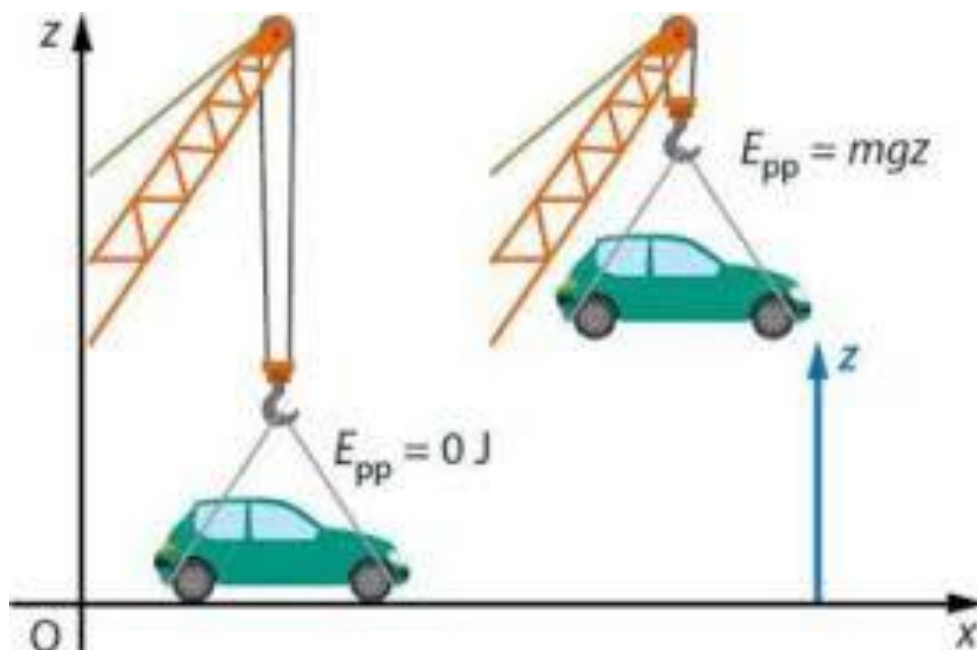
$$\mu_d = C_x / C_y = F_f / R_N \Rightarrow C_x = \mu_d m_1 g \quad (5)$$

Summing equations (3) and (4), assuming  $T_1 = T_2$  and equation (5), we obtain:

$$a = \frac{F \cos \theta - g(m_2 - \mu_d m_1)}{m_1 + m_2}$$

# COURSE OF MECHANICS OF THE MATERIAL POINT

## *Chapter VI: Work and Energy*



**Glossary**

<b>In English</b>	<b>In French</b>	<b>In Arabic</b>
The work	Le travail	العمل
External forces	Forces extérieures	القوى الخارجية
The elementary work	Le travail élémentaire	العمل الجزئي
The elementary displacement	Le déplacement élémentaire	التنقل الجزئي
The power	La puissance	الاستطاعة
The average power	La puissance moyenne	الاستطاعة المتوسطة
The instantaneous power	La puissance instantanée	الاستطاعة اللحظية
Energy	L'énergie	الطاقة
Driving work	Le travail moteur	العمل المحرك
Resistive work	Le travail résistant	العمل المقاوم
Kinetic energy	L'énergie cinétique	الطاقة الحركية
Conservatives forces	La force conservative	القوة المنخفضة
Potential energy	L'énergie potentielle	الطاقة الكامنة
Wight force	Force du poids	قوة الثقل
Spring return force	La force de rappel du ressort	قوة الأرجاع لنابض
Mechanic energy (TotaleEnergie)	L'énergie mécanique ou l'énergie totale	الطاقة الكلية
Friction force work	Le travail de la force de frottement	عمل قوة الاحتكاك

### 1. Introduction

The aim of this chapter is to present the energy tools used in mechanics to solve problems. Indeed, sometimes the fundamental principle of dynamics is not enough to solve a problem. Newton's laws can be used to solve all the problems of classical mechanics. If we know the position and initial velocity of the particles in a system, as well as all the forces acting on them. But in practice, we don't always know all the forces at play, and even if we do, the equations to be solved are too complex. In this case, other concepts such as work and energy must be used. Before describing the different types of energy (kinetic, potential and mechanical) and using them in energy theorems, we'll introduce the notions of power and work of a force.

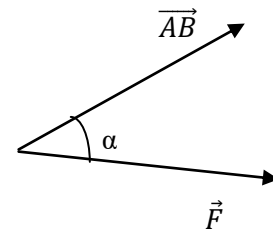
### 2. The work العمل

All motion under the action of external forces  $\vec{F}$ , implies work by these forces. In other words; work supplied by a force moves a body in its own direction and creates motion.

#### 2.1. Work performed by a constant force

Let a particle subjected to a constant force  $\vec{F}$  move this body a distance  $d=AB$ , the mechanical work  $W$  performed by the force  $\vec{F}$  is defined as:

$$W_{AB} = \vec{F} \cdot \vec{AB} = |\vec{F}| \cdot |\vec{AB}| \cdot \cos\alpha$$



$\alpha$  is the angle between the two vectors  $\vec{F}$  and  $\vec{AB}$ .

- For  $\alpha=0$   $W = |\vec{F}| \cdot |\vec{AB}|$  because  $\cos 0 = 1$
- For  $\alpha < \frac{\pi}{2}$  with have  $W > 0$  It's a driving work.
- For  $\alpha = \frac{\pi}{2}$  with have  $W = 0$  because  $\cos \frac{\pi}{2} = 0$ .
- For  $\frac{\pi}{2} < \alpha < \pi$   $W < 0$  It's a resistive work.

Unity of work in the system MKSA is « **Joule** ».

#### Note:

Note that work is a scalar quantity, unlike force and displacement, which are vectors.

#### Example 1:

The muscular effort required to lift an object depends on both its weight (the force of gravity exerted on it), and the height  $h$  from which it is lifted.

## Chapter VI: Work and Energy

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In this case, the force of the weight is directed downwards, the displacement upwards and  $\theta$  is  $180^\circ$ .

$$W = - P \cdot h = -mgh.$$

The force of the weight is negative, since muscular work must be done against the force of gravity.

### Example 2:

To lift a car with a mass of one and a half tons, a force  $F$  of 15,000N vertical to the car is required.

Calculate the work done by this force to move the car by a height (AB) of 3 meters.

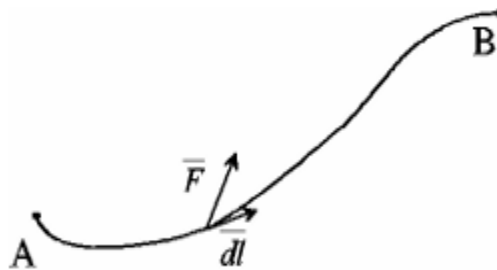
$$W_{AB}(\vec{F}) = |\vec{F}| \cdot |\overline{AB}| \cdot \cos\alpha = F \cdot d \cdot \cos\alpha = 1.5 \cdot 10^4 \cdot 3 = 4.5 \cdot 10^4 \text{ J}$$

### 2.2. The work performed by a variable force

If the force varies in intensity and/or direction during displacement, and if the displacement has any form whatsoever, we need to use integral calculus to generalize the definition of work. Generally speaking, the work of a force depends on the path followed, which is why this elementary work is necessary.

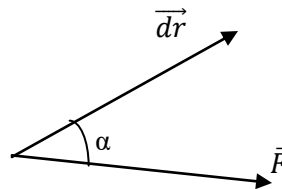
$$dW = \vec{F} \cdot \overline{dr} = \vec{F} \cdot \overline{dl}$$

where  $dl$  is an infinitesimal displacement along the trajectory, tangential to it.



The elementary work  $dW$  performed by a force  $\vec{F}$  on a point mass  $m$  during an elementary displacement  $dr = dl$  is given by:

$$dW = \vec{F} \cdot \overline{dr} = |\vec{F}| \cdot |\overline{dr}| \cos(\vec{F}, \overline{dr})$$



To obtain the work on an AB displacement, we integrate this elementary work:

$$W = \int dW = \int_A^B \vec{F} \cdot \overline{dr} = \int F \cdot dr \cdot \cos\alpha$$

$\alpha$  is the angle between the two vectors  $\vec{F}$  and  $\overline{dr}$ ;  $\alpha = (\vec{F}, \overline{dr})$

### 2.3. The power الاستطاعة

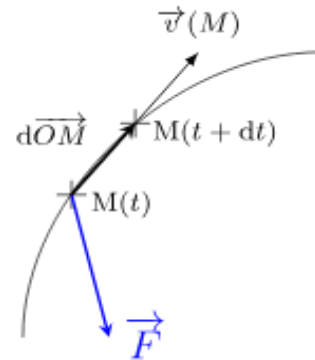
Let a point M move along its trajectory at a velocity  $\vec{v}(M)$  relative to the reference frame of study, It experiences a force  $\vec{F}(M)$  as shown in the figure opposite:

The power of a force  $\vec{F}$  is the work per unit time.

We have two types:

- The average power  $P_{avr} = \frac{\Delta W}{\Delta t}$
- The instantaneous power  $P = \frac{dW}{dt}$

Then the instantaneous power of the  $\vec{F}$  is:



$$P(\vec{F}) = \frac{dW}{dt} = \frac{|\vec{F}| \cdot |d\vec{r}|}{dt} = \vec{F} \cdot \vec{v}(M) = \|\vec{F}\| \times \|\vec{v}(M)\| \times \cos\alpha$$

**Note :**

- ✓ The unit of power is the « Watt ».
- ✓ This force can be classified into three types:
  - It is driving, if its power is positive which corresponds to an angle  $\alpha < \pi/2$ .
  - It is resistive, if its power is negative which corresponds to an angle  $\alpha > \pi/2$ .
  - Finally, it can be of zero power, in which case  $\alpha = \pi/2$ .

### 3. Energy الطاقة

In physics, energy is defined as the capacity of a system to produce work. Energy is not a material substance: it is a physical quantity that characterizes the state of a system; it can be stored and exists in many forms.

#### 3.1. Kinetic energy الطاقة الحركية

In order to accelerate a point mass to a defined speed, work must be done. This work is then stored in the point mass in the form of kinetic energy.

Suppose the object's initial velocity is  $v_0$  and the force F is applied in the direction of  $v_0$ , producing a displacement  $d = dr$ .

We have:  $dW = F \cdot dr$  and  $F = ma = m \frac{dv}{dt}$

From this expression we can deduce the following:

$$dW = F dr = m \frac{dv}{dt} dr$$

$$\Rightarrow dW = m \frac{dr}{dt} dv \quad \text{Then } dW = mv dv$$

Let's integrate the expression of elementary work, and derive the definition of kinetic energy:

$$W = m \int_A^B v dv \Rightarrow W = \frac{1}{2} m (v_B^2 - v_A^2) = \frac{1}{2} m v_B^2 - \frac{1}{2} m v_A^2$$

Where  $v_A$  is the velocity of the moving body at point A and  $v_B$  its velocity at point B.

The kinetic energy of a material point of mass  $m$  and instantaneous velocity  $\vec{v}$  is given by the expression:

$$E_c = \frac{1}{2} m v^2$$

$$\text{So: } W_{\vec{F}(A \rightarrow B)} = E_{cB} - E_{cA} = \Delta E_c$$

**Note :**

- ✓ The unit of energy is the « Joule ».
- ✓ And since  $p=mv$ , we can also write:

$$E_c = \frac{p^2}{2m}$$

### Theorem of the Kinetic Energy Theorem : نظرية الطاقة الحركية

The variation in kinetic energy of a material point subjected to a set of external forces between two positions A and B is equal to the sum of the work of these forces between these two points.

$$W_{\vec{F}(A \rightarrow B)} = E_{cB} - E_{cA} = \Delta E_c \Rightarrow \sum_i W_i = \Delta E_c$$

### 3.2. Conservative forces القوة المنخفضة

A force is said to be conservative, or to derive from a potential, if its work is independent of the path taken, whatever the probable displacement between the starting point and the end point.

Conservative forces include the force of gravity, spring return force and the tension force of a wire.

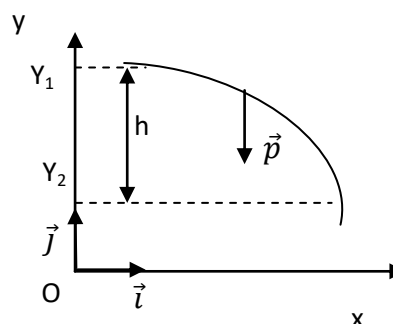
**Example:**

Let's calculate the work of the force of gravity.

$$dW = \vec{p} \cdot d\vec{l} \quad \text{with } p = -mg\vec{j}$$

$$d\vec{l} = dx\vec{i} + dy\vec{j}$$

$$\text{so } dW = -mgdy$$





$$W = -mg \int_{y_1}^{y_2} dy = -mg(y_2 - y_1)$$

$$\Rightarrow W = mg(y_1 - y_2) = mgh$$

So the force of gravity  $\vec{p}$  is a conservative force because its work does not depend on the path followed, and it is said to derive from a potential.

Spring return force is also a conservative force.

**Note:**

A force is said to be non-conservative if its work depends on the path followed, as in the case of *friction force*.

### 3.3. Potential energy الطاقة الكامنة

Potential energy is a function of coordinates, such as the integration between its two values at start and finish. It represents the work done by the particle to move it from its initial position to its final position.

If the force  $\vec{F}$  is a force deriving from a potential (conservative), then:

$$W = \int_A^B \vec{F}_C \cdot d\vec{r} = E_{P_A} - E_{P_B} \Rightarrow dW = -dE_p$$

$$\text{Hence } W_{A \rightarrow B}(\vec{F}_C) = -\Delta E_p$$

Potential energy is always calculated relative to a reference frame ( $E_p=0$ ).

The potential energy function  $E_p$  is determined to within one constant.

By identifying the two expressions  $dE_p$  and  $dW$ , we arrive at the following result: The differential of potential energy is equal to and opposite in direction to the differential of work.

#### Example 1: Wight force قوة الثقل

The force of weight is a conservative force, hence:

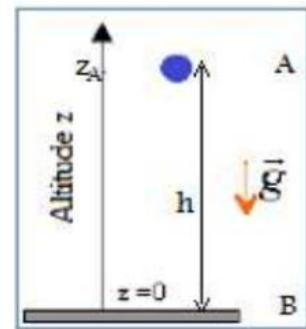
$$W_{A \rightarrow B}(\vec{F}_C) = -\Delta E_p$$

$$\text{And } W = mg(y_1 - y_2) = mgh$$

$$\text{So } W_{\vec{p}} = -\Delta E_p = -(E_{pf} - E_{pi}) = E_{pi} = mg(z_A - z_B)$$

Because  $E_{pf}$  is the reference potential energy.

$$\text{So } E_p = mgH$$

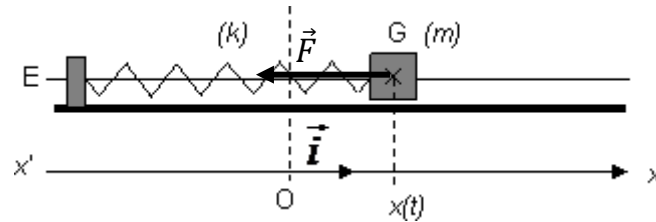


Note:

If  $Z_A > Z_B$  we have  $E_p > 0$

If  $Z_A < Z_B$  we have  $E_p < 0$

### Example 2 : Spring return force قوة الارجاع ل نابض



$$\vec{F} = -kx\vec{i}, \vec{dl} = dx.\vec{i} \text{ et } dW = \vec{F} \cdot \vec{dl}$$

$$dW = -dE_p = -kx \cdot dx \Rightarrow dE_p = kx dx$$

$$\Rightarrow \int dE_p = k \int_{x_i}^{x_f} x dx$$

$$\Rightarrow E_p = \frac{1}{2}k(x_f^2 - x_i^2) = \frac{1}{2}kx^2$$

### 3.4. Mechanic energy (Totale Energie) الطاقة الكلية

The mechanical energy of a material point at a given instant is equal to the sum of kinetic energy and potential energy:

$$E_M = E_C + E_p \Rightarrow E_M = E_C + E_p$$

- **Principle of conservation of mechanical energy** مبدأ انحفاظ الطاقة الميكانيكية

In a conservative (or potential-derived) force field, mechanical energy is conserved over time (**no friction**).

$$E_M = E_C + E_p = Cte$$

This means that the variation in mechanical energy is zero  $\Delta E_M = 0$ , it also means that the variation in kinetic energy is equal to the opposite of the variation in potential energy:

$$\Delta E_c = - \Delta E_p$$

In other words, if the system is isolated or free, mechanical energy is conserved.

## Chapter VI: Work and Energy

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### Note:

In the presence of **frictional forces**, the variation in mechanical energy can't be stored, is equal to the sum of the work of the frictional forces.  $W(F_{frott})$ :

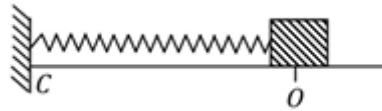
$$\Delta E_M = E_{Mf} - E_{Mi} = \sum W_{A \rightarrow B}(\vec{F}_{NC}) = W_{A \rightarrow B}(\vec{F}_{frott})$$

- **Friction force work** : عمل قوة الاحتكاك

$$W_{A \rightarrow B}(\vec{F}_{frott}) = -F_f \cdot AB$$

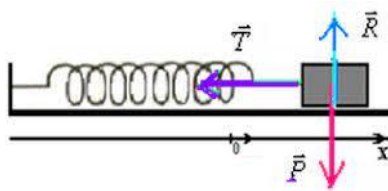
### Example:

A mass  $m$  is attached to a spring of stiffness  $k$ , and the other end of the spring is attached to point C. The mass  $m$  can slide on the horizontal surface. Initially, the mass is at rest at point O of equilibrium.



- 1) Assuming no friction, move mass  $m$  from point O to point A, such that  $OA=a$ . Determine the work of the Spring return force as  $m$  moves from O to A. Then determine the speed of  $m$  at point O.
- 2) Same questions as question 1, but now we assume that friction exists, and give the dynamic friction coefficient  $\mu_c$ .

### Answers:



1- We have,  $\vec{F} = -kx\vec{i}$  and  $d\vec{l} = dx\vec{i}$

$$\Rightarrow W_{\vec{F}} = \int dW_{\vec{F}} = \int \vec{F} \cdot d\vec{l} = -k \int_a^0 x dx = \frac{1}{2}ka^2$$

We also have:  $\sum_i W_i = \Delta E_C = W_{\vec{p}} + W_{\vec{R}} + W_{\vec{F}}$

With,  $W_{\vec{p}} = W_{\vec{R}} = 0$  because  $\vec{R}$  and  $\vec{p} \perp \vec{Ox}$

## Chapter VI: Work and Energy

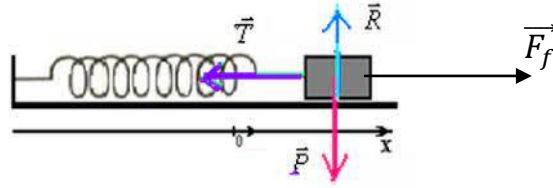
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So  $\Delta E_C = W_{\vec{F}} = \frac{1}{2}ka^2 = \frac{1}{2}mv_o^2 - \frac{1}{2}mv_A^2$  with  $v_A=0$

Hence,  $v_o = a\sqrt{\frac{k}{m}}$

### 2- Case of friction

We also have:  $\sum_i W_i = \Delta E_C = W_{\vec{p}} + W_{\vec{R}} + W_{\vec{F}_f} + W_{\vec{F}}$



with:  $W_{\vec{p}} = W_{\vec{R}} = \vec{0}$

so  $\Delta E_C = W_{\vec{F}} + W_{\vec{F}_f} = \frac{1}{2}ka^2 - a.F_f = \frac{1}{2}ka^2 - a.\mu_c.mg = \frac{1}{2}mv_o^2$  because  $v_A=0$

Hence,  $v_o = \sqrt{\frac{ka^2}{m} - 2\mu_c.a.g} = a\sqrt{\frac{k}{m} - \frac{2\mu_c.g}{a}}$

### Proposed exercises about chapter VI

#### Exercise 1

A body is subjected to a force  $\vec{F}$  such that:  $\vec{F} = (y^2 - x^2)\vec{i} + (3xy)\vec{j}$

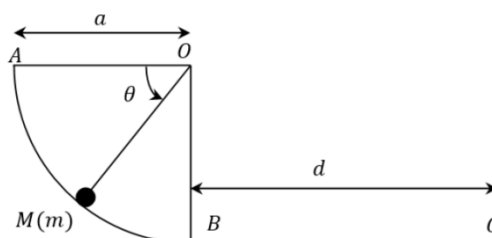
Find the work of force  $\vec{F}$  if the body moves from point A(0,0) to point B(1,3) Following the following trajectories:

1. On the (Ox) axis from A to C(1,0) then parallel to (Oy) from C to B.
2. On the (Oy) axis from A to D(0,3), then parallel to (Ox) from D to B.
3. On the straight line [AB].
4. On trajectory  $y=x^2$ .

#### Exercise 2

A particle of mass  $m$ , initially at rest in A, slides without friction on the circular surface AOB of radius  $a$ .

- 1) Determine the work of weight from A to M.
- 2) Determine the work of the surface-particle contact force  $m$ .
- 3) Determine the potential energy  $E_p$  of  $m$  at the point M ( $E_p(B) = 0$ ).



- 4) Use the kinetic energy theorem to determine the speed of  $m$  at point M, deduce its kinetic energy  $E_c$ .
- 5) Calculate the mechanical energy  $E_m$ .
- 6) Show  $E_c$ ,  $E_p$  and  $E_m$  ( $0 < \theta < \pi/2$ ). Discuss.
- 7) The circular surface AOB is connected to a horizontal part BC, there is friction between B and C, the particle stops at a distance  $d$  from B. Determine the coefficient of kinetic friction.

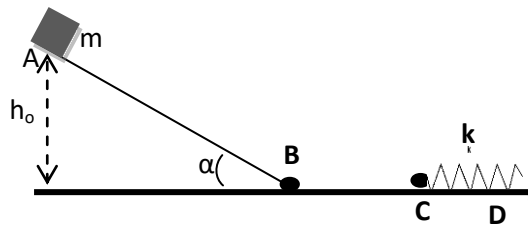
Given  $d = 3a = 3m$ .

#### Exercise 3

Consider a small block of mass  $m = 5\text{kg}$  dropped without initial velocity at point A of an inclined plane at an angle  $\alpha = 30^\circ$  to the horizontal. Point A is at a height  $h_0 = 5\text{m}$  from the horizontal.

- 1- Knowing that the coefficient of dynamic friction on plane AB is  $\mu_d = 0.2$ , applying the fundamental principle of dynamics:

## Chapter VI: Work and Energy



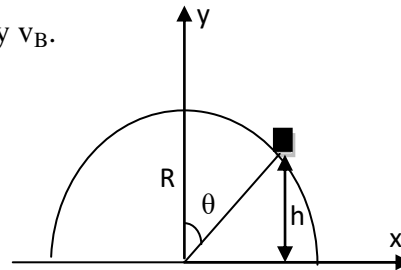
- What is the nature of the motion on plane AB?
  - Calculate the speed of the block when it reaches point B.
- 2- After passing through point B at speed  $V_B$ , the mass arrives at point C. Knowing that the coefficient of friction is negligible on plane BC :
- Deduce the speed at point C?
  - Calculate the maximum compression of the spring, given a stiffness constant equal to  $k=100\text{N/m}$ ? ( $g = 10 \text{ m/s}^2$ ).

### Exercise 4

A piece of ice M of mass m slides without friction over the outer surface of an igloo, which is a half-sphere of radius r with a horizontal base.

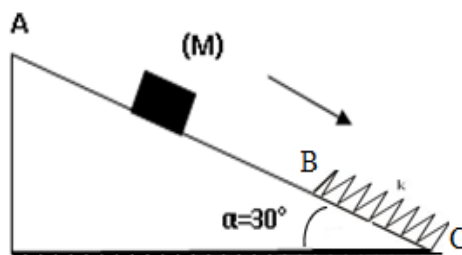
At  $t=0$ , it is released from point A without any initial velocity.

1. Find the expression for the velocity at point B, as a function of g, r and  $\theta$ .
2. Using the fundamental relation of dynamics, determine the expression of  $|\vec{N}|$  the reaction of the igloo on M at point B as a function of velocity  $v_B$ .
3. At what height does M leave the sphere?
4. At what speed does M arrive at the axis (Ox)?



### Exercise 5

Consider a small block of mass  $m = 2\text{kg}$  dropped without initial velocity at point A of an inclined plane at an angle  $\alpha=30^\circ$  to the horizontal. Point A is at a height  $h_A=5\text{m}$  from the horizontal.



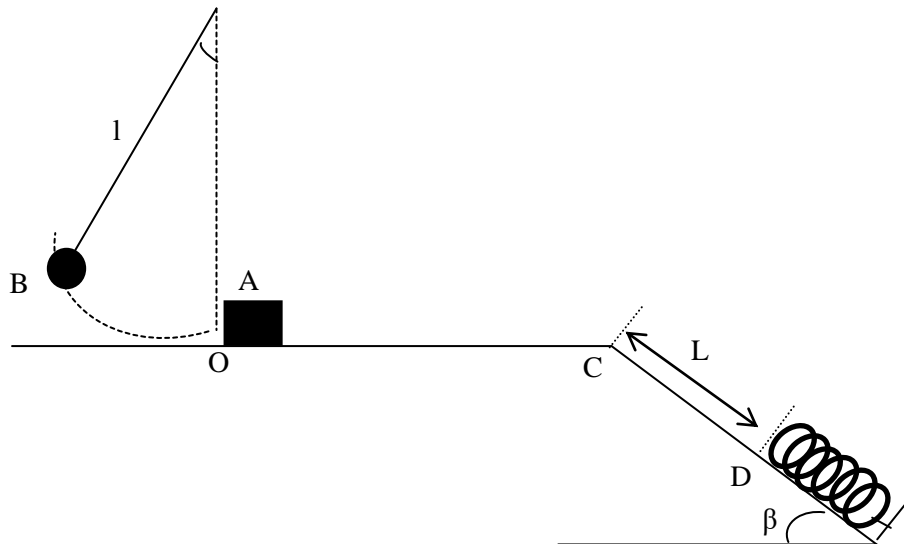
## Chapter VI: Work and Energy

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1. Knowing that the coefficient of dynamic friction on plane AB is  $\mu_d=0.2$ , applying the fundamental principle of dynamics, what is the acceleration of the block on plane AB=8m?
2. Calculate the speed of the block when it reaches point B.
3. Using the kinetic energy theorem, find the speed of the block at point B.
4. At point B, the block hits a spring with stiffness constant  $k=100\text{N/m}$  at speed  $V_B$ . Calculate the maximum compression ( $x$ ) of the spring (given  $g = 10 \text{ m/s}^2$ ).

### Exercise 6

- A ball B of mass  $m$ , attached to an inextensible wire of length  $l$ , is moved away from its equilibrium position by an angle  $\alpha$ . It is dropped without initial velocity. Passing through the vertical position, the ball strikes (touches) a body A of the same mass and stops, body A passes from point O to point C ( $OC=d$ ) on a rough horizontal plane of friction coefficient  $\mu$ .



1. Show the forces exerted on body A.
2. What is the nature of the motion on the horizontal plane?
3. Express the velocity of ball B just before touching body A.
4. Using the principle of conservation of momentum, determine the velocity of body A after the interaction.
5. If  $v_A = v_B$  at point O, give the velocity of body A at point C as a function of  $g$ ,  $l$ ,  $d$ ,  $\alpha$  and  $\mu$ .
6. By what angle must ball B be moved away for body A to arrive at point C with zero velocity.

## Chapter VI: Work and Energy

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- From point C, body A approaches the perfectly smooth (no friction) path  $CD=L$ , inclined at an angle  $\beta$  to the horizontal. It arrives, without initial velocity, on a perfect spring of length  $l_0$  and stiffness constant  $k$ .

- Show the forces exerted on A as the spring compresses.
- What is the value of the spring's maximum compression?

We give  $m=200\text{g}$ ,  $d=OC=1\text{m}$ ,  $l=10\text{ cm}$ ,  $L=1\text{m}$ ,  $\mu=0.1$ ,  $g=10\text{m/s}^2$ ,  $k=140\text{N/m}$ ,  $\beta=30^\circ$ .

### Exercise 7

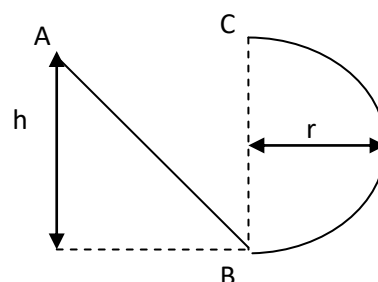
A solid body S of mass  $m$  is linked on one side to a spring of stiffness  $K$ , while the other side of the spring is fixed. The body is moved horizontally from its equilibrium position by a distance  $x$  and then released ( $\mu = \text{tg } \varphi$  : coefficient of friction).

- Show the forces applied to body S.
- Calculate the speed  $V_B$  corresponding to the movement of S from its equilibrium position.

### Exercise 8

A ball slides without friction inside a gutter.

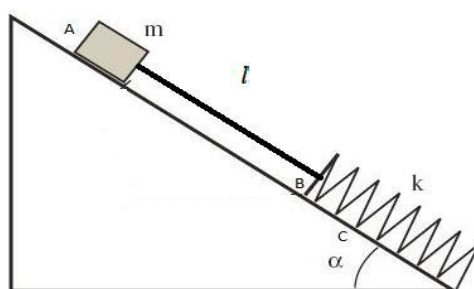
Find the smallest height  $h_{\min}$  from which the ball is launched to reach point C, without leaving the gutter.



### Exercise 9

A block of mass  $m$  is dropped without initial velocity onto an inclined plane making an angle  $\alpha$  with the horizontal, at a distance  $l$  above a light, uncompressed spring of stiffness  $k$ . The motion of the block is frictionless.

- Using the principle of conservation of mechanical energy, find the expression for the velocity of the block when it first touches the spring.
- Using the fundamental principle of dynamics, find the expression for the maximum compression of the spring as a function of  $m$ ,  $g$ ,  $\alpha$  and  $k$ .





### Correction of exercises about chapter VI

#### Exercise 1

$$\vec{F} = (y^2 - x^2)\vec{i} + (3xy)\vec{j}$$

The work of force  $\vec{F}$  when the body moves from point A(0,0) to point B(2,4) along the trajectories:

$$dW = \vec{F} \cdot d\vec{r} \text{ with } d\vec{r} = dx\vec{i} + dy\vec{j}$$

Then

$$dW = \vec{F} \cdot d\vec{r} = \begin{pmatrix} y^2 - x^2 \\ 3xy \end{pmatrix} \cdot \begin{pmatrix} dx \\ dy \end{pmatrix} \Rightarrow dW = (y^2 - x^2)dx + 3xydy$$

$$\text{So } W = \int \vec{F} \cdot d\vec{r} = \int F_x dx + \int F_y dy \Rightarrow W = \int (y^2 - x^2)dx + \int 3xydy$$

1- Following axis (Ox) from A(0,0) to C(2,0):

The variation is on the Ax axis, so  $y=0$ ; therefore  $dy=0$  and  $x$  varies from 0 to 2.

$$W = \int (y^2 - x^2)dx + \int 3xydy = \int_0^2 -x^2 dx \Rightarrow W = -\frac{x^3}{3} = -\frac{8}{3}j$$

2- Following axis (Oy) from C(2,0) to B(2,4):

The variation is parallel to Oy so  $x$  is constant ( $x=2$ ) then  $dx=0$  and  $y$  varies from 0 to 4:

$$W = \int_0^4 3xy dy = \int_0^4 6y dy \Rightarrow W = 6\frac{y^2}{2} = 48j$$

3- On line AB:

The equation of a straight line is generally of the form:  $y=a.x+b$

If the line passes through the two points A(0,0) and B(2,4) then  $b=0$  and  $a = \frac{y_B - y_A}{x_B - x_A} = 2$

Then the equation of the line is of the form  $y=2.x$ , so  $dy= 2.dx$ .

In this case, the expression for work becomes:

( $x$  varies from 0 to 2);

$$W = \int ((2x)^2 - x^2)dx + \int 3x(2x)2dx = 15 \int_0^2 x^2 dx \Rightarrow W = 15\frac{x^3}{3} = 40j$$

## Chapter VI: Work and Energy

4- On the trajectory  $y=x^2$ :

$$y=x^2 \Rightarrow dy = 2x \cdot dx$$

Then the work formula becomes:

$$W = \int ((x^2)^2 - x^2)dx + \int 3x(x^2)2xdx = \int 7x^4 - x^2 dx$$

The body moves from A(0,0) and B(2,4) then x varies from 0 to 2; then,

$$W = \frac{7}{5}x^5 - \frac{x^3}{3} = 42,13j$$

### Exercise 2

1) The work of  $\vec{p}$  from A to M is:

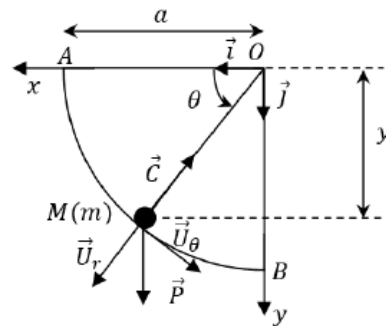
$$dW = \vec{p} \cdot d\vec{l} \quad \text{with } p = mg\vec{j}$$

$$d\vec{l} = dx\vec{i} + dy\vec{j} \quad \text{so } dW = mgdy$$

$$W = mg \int_0^y dy = mgy = mga \sin\theta$$

2) The work of  $R_N$  force is:

$$W_R = \int_0^y \vec{R}_N \cdot d\vec{l} = 0 \quad \text{Because } \vec{R}_N \perp d\vec{l}$$



3) Potential energy:

$$dE_p = -dW \Rightarrow E_p = -mga \sin\theta + c$$

$$E_p(B) = 0, \theta = \pi/2 \text{ so } c = mga$$

$$\Rightarrow E_p = mga(1 - \sin\theta)$$

$$4) \Delta E_C = \sum W \Rightarrow \frac{1}{2}mv_M^2 = mga \sin\theta$$

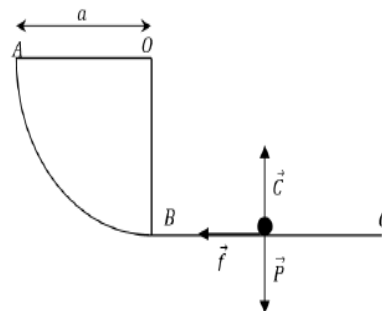
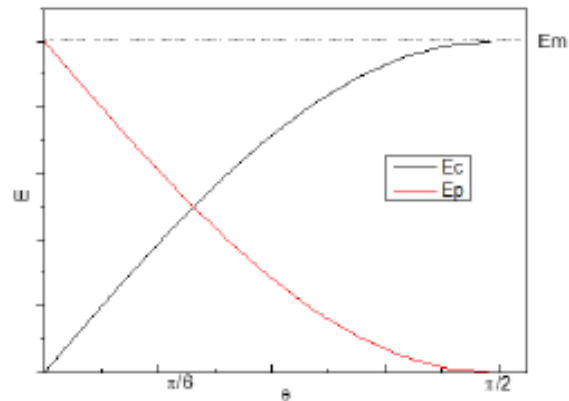
$$V_M = \sqrt{2g a \sin\theta}$$

$$5) E_m = E_C + E_p = mga = \text{cste}$$

6) When  $E_p$  decreases  $E_C$  increases while  $E_m$  remains constant.

$$7) \mu = \frac{f}{R_N} = \frac{f}{p} \Rightarrow f = \mu mg$$

$$\text{So } \Delta E_C = W_f = \int_B^C \vec{f} \cdot d\vec{l} = -\mu mgd \Rightarrow \frac{1}{2}mv_B^2 = -\mu mgd$$



## Chapter VI: Work and Energy

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Then  $v_B = \sqrt{2ag}$

**Note :** Replace  $\theta = \pi/2$  in the formula for  $v_M$ , we find :

$$v_B = \sqrt{2ag}$$

We can use also:  $E_{m_B} = E_{m_A} \Rightarrow E_{C_A} + E_{P_A} = E_{C_B} + E_{P_B}$

**Calculation of  $\mu$ :**

$$\sum \vec{F} = m\vec{\gamma} = \vec{f} + \vec{P} + \vec{R}_N$$

We have  $\mu = \frac{f}{R_N} = \frac{f}{mg}$  because  $R_N = mg$  (with projection on (oy))

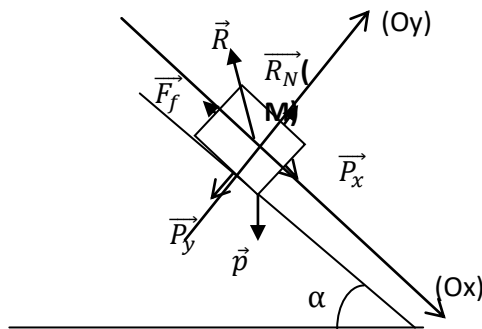
Projection on (ox):  $-f = m\gamma$

We have also:  $v_C^2 - v_B^2 = 2\gamma \cdot d$  ( $v_C=0$ )

$$-v_B^2 = 2\gamma \cdot d = -2ag \quad \text{so: } \gamma = \frac{-ag}{d} \quad \text{with } -f = m\gamma = \frac{-mag}{d} \quad \text{so } f = \frac{m \cdot a \cdot g}{d}$$

$$\text{Then } \mu = \frac{f}{R} = \frac{f}{mg} = \frac{mag}{mg \cdot d} = \frac{a}{d} = \frac{1}{3}$$

### Exercise 3



1. Knowing that the coefficient of dynamic friction on plane AB is  $\mu_d=0.2$ , apply the fundamental principle of dynamics:

- What is the nature of the motion on AB?  $a= ?$

$$\Sigma \vec{F} = m\vec{a} = \vec{p} + \vec{R}_N + \vec{F}_f$$

Following (Ox):  $-F_f + p_x = -F_f + m g \sin\alpha = ma$

Following (Oy):  $R_N - p_y = 0 \Rightarrow R_N = m g \cos\alpha$

$$\mu_d = \tan\varphi = F_f/R_N \Rightarrow F_f = R_N \tan\varphi = \mu_d m g \cos\alpha$$

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$$-\mu_d m g \cos\alpha + m g \sin\alpha = m \cdot a$$

$$\Rightarrow a = g \cdot (\sin\alpha - \mu_d \cos\alpha)$$

$$\text{So: } a = 3.26 \text{ m/s}^2$$

- Calculate the speed of the block when it reaches point B.

$$v_B^2 - v_A^2 = 2al \Rightarrow v_B^2 = 2al = 2a\left(\frac{h}{\sin\alpha}\right)$$

$$v_B = \sqrt{2a\left(\frac{h}{\sin\alpha}\right)} = 8.074 \text{ m/s}$$

2.  $V_c = V_B$  because we have an MRU (principle of inertia or Newton's 1st law)

We calculate the compression distance of the spring:

$$\Delta E_c = \Sigma W_{f_{ext}} \Rightarrow E_{cD} - E_{cC} = W_p + W_{Fr} + W_{RN}$$

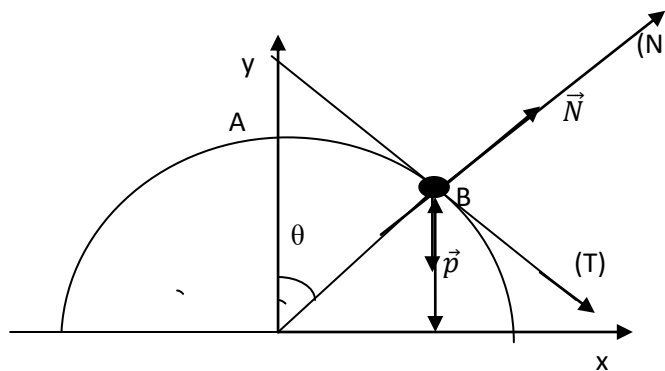
$$-\frac{1}{2}kx^2 = -\frac{1}{2}mv_c^2 \text{ so: } x = \sqrt{\frac{mv_c^2}{k}} = 1.8 \text{ m}$$

**2<sup>nd</sup> Method:** Between points C and D

$$E_{M_C} = E_{M_D} \Rightarrow E_{C_C} + E_{P_C} = E_{C_D} + E_{P_D} \Rightarrow \frac{1}{2}kx^2 = \frac{1}{2}mv_c^2$$

$$\text{So: } x = \sqrt{\frac{mv_c^2}{k}} = 1.8 \text{ m}$$

**Exercise 4**



1- According to the principle of conservation of mechanical energy between two points A and B:

$$E_{M_A} = E_{M_B} \Rightarrow E_{C_A} + E_{P_A} = E_{C_B} + E_{P_B}$$

$$\text{So: } E_{C_A} = E_{C_B} + E_{P_B} (*)$$

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Because  $E_{C_A} = 0$  ( $v_A = 0$ ) because the material point is launched without initial velocity

With  $h_B = R \cos \theta$

So; (\*)  $\Rightarrow mgR = \frac{1}{2}mv_B^2 + mg R \cos \theta$

Then:  $gR = \frac{1}{2}v_B^2 + g R \cos \theta \Rightarrow v_B^2 = 2(gR - gR \cos \theta)$

$$\Rightarrow v_B = \sqrt{2(gR - gR \cos \theta)}$$

2- According to the fundamental principle of dynamics:

$$\Sigma \vec{F} = m\vec{a} \Rightarrow \vec{N} + \vec{p} = m\vec{a}$$

We choose a reference frame consisting of the axis (OT) tangent to the half-sphere and the axis (ON) following the radius and in the direction of  $\vec{N}$  :

Projecting on (ON):

$$N \cdot \cos \theta = m a_N \Rightarrow N - mg \cos \theta = -m \frac{v^2}{R}$$

3- When point P leaves the sphere  $N=0$  so:

$$mg \cdot \cos \theta = m \frac{v_p^2}{R} \Rightarrow v_p^2 = Rg \cdot \cos \theta$$

$$(*) \Rightarrow R = \frac{1}{2}R g \cos \theta + g R \cos \theta \Rightarrow \cos \theta = \frac{2}{3} \quad \text{so } \theta_0 = 48^\circ$$

The material point P leaves the sphere at height:  $h_p = \frac{2}{3}R$

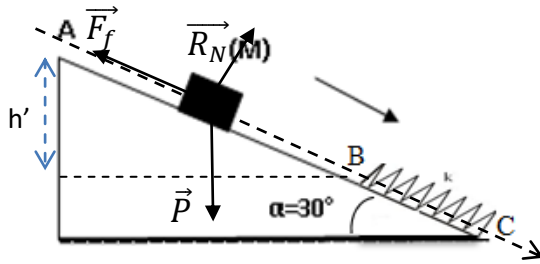
The angle relative to the horizontal at which the point leaves the half-sphere is:  $90-48=52$

4- The velocity of the material point at this point:

$$v_p^2 = Rg \cdot \cos \theta$$

$$\Rightarrow v_p = \sqrt{\frac{2}{3}Rg}$$

Exercise 5



1- The acceleration of mass m on AB:

$$\text{By applying the FPD: } \Sigma \vec{F} = m\vec{a} \Rightarrow \vec{p} + \vec{R}_N + \vec{f} = m\vec{a}$$

$$\text{Following (Ox): } -f + p_x = -f + m g \sin\alpha = ma \dots (1)$$

$$\text{Following (Oy): } R_N - p_y = 0 \Rightarrow R_N = m g \cos\alpha \dots (2)$$

$$\mu_d = \tan\phi = F_f / R_N \Rightarrow F_f = N \tan\phi \quad \text{so} \quad F_f = \mu_d m g \cos\alpha$$

$$(1): -\mu_d m g \cos\alpha + m g \sin\alpha = m \cdot a \Rightarrow a = g (\sin\alpha - \mu_d \cos\alpha) = 3.27 \text{m/s}^2$$

2- The velocity at point B:

we have  $v_A = 0$

$$\text{and } v_B^2 - v_A^2 = 2a(AB) \Rightarrow v_B^2 = 2a(AB)$$

with  $(AB) = 8\text{m}$

$$\Rightarrow v_B = \sqrt{2(3.27)(8)} = 7.23 \text{m/s}^{-1}$$

3- Applying the kinetic energy theorem, find the speed of the block when it reaches point B.

$$\Delta E_c = \Sigma W_{f_{ext}} \Rightarrow E_{c_B} - E_{c_A} = W_p + W_{F_f} + W_{R_N}$$

$$\frac{1}{2} m v_B^2 = m g \sin\alpha AB - F_f AB$$

And

$$\frac{1}{2} m v_B^2 = m g \sin\alpha AB - \mu_d m g \cos\alpha AB$$

$$\text{so } v_B = \sqrt{g \cdot 2 \cdot AB \sin\alpha - 2\mu_d g \cos\alpha AB}$$

$$v_B = \sqrt{g \cdot 2 \cdot AB (\sin\alpha - \mu_d \cos\alpha)}$$

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4- At point B, the block touches a spring with stiffness constant  $k=100\text{N/m}$  at speed  $V_B$ .

Calculate the maximum compression ( $x$ ) of the spring?

(we give  $g = 10 \text{ m/s}^2$ ).

$$\Delta E_M = E_{M_C} - E_{M_B} = \Sigma W_{f_{NC}} \Rightarrow (E_{C_C} + E_{P_C}) - (E_{C_B} + E_{P_B}) = W_{F_f}$$

$$\Rightarrow \frac{1}{2} k x^2 - \frac{1}{2} m v_B^2 - m g h' = \frac{1}{2} k x^2 - \frac{1}{2} m v_B^2 - m g (AB \sin \alpha) = -\mu_d m g \cos \alpha AB$$

So;

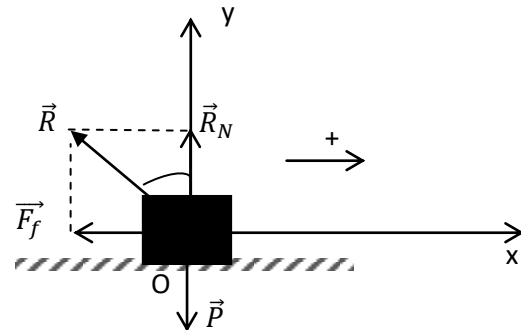
$$x = \sqrt{\frac{m(v_B^2 + g 2(AB \sin \alpha) - \mu_d 2 g \cos \alpha AB)}{k}} = 1.07 \text{ m}$$

### Exercise 6

1- Representation of forces on the figure:

2- Acceleration:

According to the fundamental principle of dynamics:



$$\Sigma \vec{F} = m \vec{a} \Rightarrow \vec{R}_N + \vec{F}_f + \vec{P} = m \vec{a}$$

Following (Ox):  $-F_f = ma$  (1)

Following (Oy):  $R_N - P = 0$

$$\mu = \frac{F_f}{R_N} \Rightarrow F_f = \mu R_N ; \text{ with } P = R_N \text{ so } F_f = \mu m g$$

$$(1) \Rightarrow -\mu m g = m a$$

so  $a = -\mu g = -1 \text{ m/s}^2$

$a = -1$  so we have uniformly decelerated rectilinear motion.

3- Il n'y a pas de frottement donc d'après le principe de conservation de l'énergie mécanique.

$$E_{M_i} = E_{M_f} \Rightarrow E_{C_i} + E_{P_i} = E_{C_f} + E_{P_f} \text{ with } v_i = 0 \text{ so } E_{C_i} = 0$$

$$E_{P_i} = m g l (1 - \cos \alpha), E_{C_f} = \frac{1}{2} m v_f^2 \text{ and } E_{P_f} = 0$$

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$$mgl(1 - \cos\alpha) = \frac{1}{2}mv_f^2 \Rightarrow \frac{1}{2}mv_B^2 = mgl(1 - \cos\alpha)$$

$$\text{So } v_B = \sqrt{2gl(1 - \cos\alpha)}$$

4- According to the principle of conservation of momentum:

$$P_i \text{ (before impact)} = P_f \text{ (after impact)}$$

$$m\vec{v}_B + \vec{0} = \vec{0} + m\vec{v}_A \Rightarrow v_A = v_B = \sqrt{2gl(1 - \cos\alpha)}$$

5- According to the principle of kinetic energy between O and C:

$$\Delta E_C = \Sigma W_{(F_{ext})} \Rightarrow E_{C_C} - E_{C_O} = W_P + W_{R_N} + W_{F_f}$$

$$\frac{1}{2}mv_C^2 - \frac{1}{2}mv_O^2 = -F_f(OC)$$

With  $F_f = \mu mg$ ,  $v_O^2 = 2gl(1 - \cos\alpha)$  and  $OC = d$

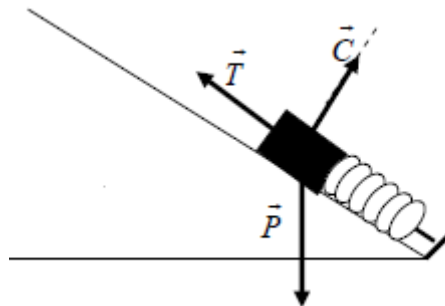
$$\text{So } v_C = \sqrt{2gl(1 - \cos\alpha) - 2\mu gd}$$

6- Find the angle  $\alpha$  so that  $v_C = 0$

$$v_C = 0 \Rightarrow \sqrt{2gl(1 - \cos\alpha) - 2\mu gd} = 0$$

$$\text{So } \cos\alpha = 1 - \frac{\mu d}{l} \Rightarrow \alpha_{min} = \frac{\pi}{2}$$

- 1- Representation of forces on the figure:



2- There is no friction, so according to the principle of conservation of mechanical energy between point C and D:

$$E_{MC} = E_{MD} \Rightarrow \Delta E_C = -\Delta E_C \Rightarrow E_{C_C} + E_{P_C} = E_{C_D} + E_{P_D}$$



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With  $v_D=0$  because body A changes direction at point D, so  $E_{CD} = 0$

The body approaches the CD section without any initial speed, so  $v_C=0$  and  $E_C=0$

$$E_{Pc} = mgh \text{ with } h = (L + x)\sin\beta \text{ so } E_{Pc} = mg(L + x)\sin\beta$$

$$E_{PD} = \frac{1}{2}kx^2$$

$$\text{So } \frac{1}{2}kx^2 = mg(L + x)\sin\beta \Rightarrow \frac{1}{2}kx^2 - mgL\sin\beta - mgx\sin\beta = 0$$

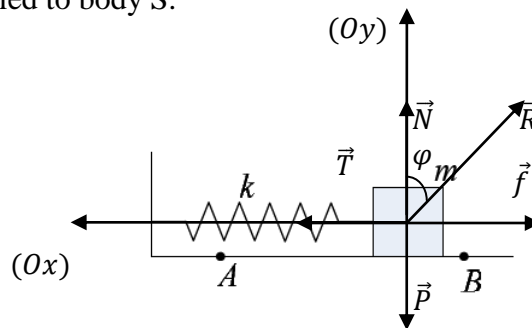
$$\Rightarrow 70x^2 - x - 1 = 0$$

$$\text{So } x=12.7\text{cm}$$

### Exercise 7

A solid body S of mass  $m$  is linked on one side to a spring of stiffness  $K$ , while the other side of the spring is fixed. The body is moved horizontally from its equilibrium position by a distance  $x$  and then released ( $\mu = \tan \varphi$  : coefficient of friction).

1. Show the forces applied to body S.



- 2- Calculate the speed  $V_B$  corresponding to the movement of S from its equilibrium position.

$$\Delta E_c = \Sigma W_{f_{ext}} \Rightarrow E_{c_B} - E_{c_A} = W_p + W_T + W_N + W_f$$

With,  $W_{\vec{p}} = W_{\vec{R}} = 0$  because  $\vec{R}$  and  $\vec{p} \perp \vec{Ox}$

So  $E_{c_B} = W_T + W_f$  (\*) ( $E_{c_A} = 0$  because the initial velocity is null)

We have,  $\vec{F} = -kx\vec{i}$  and  $d\vec{l} = dx\vec{i}$

$$\Rightarrow W_{\vec{F}} = \int dW_{\vec{F}} = \int \vec{F} \cdot d\vec{l} = -k \int_x^0 x dx = \frac{1}{2}kx^2$$

## Chapter VI: Work and Energy

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so  $W_T = \frac{1}{2}kx^2$  and  $W_f = -f x$

According to the fundamental principle of dynamics:

$$\Sigma \vec{F} = m\vec{a} \Rightarrow \vec{N} + \vec{p} + \vec{T} + \vec{f} = m\vec{a}$$

$$\mu = \operatorname{tg} \varphi = \frac{f}{N}$$

We choose a frame of reference composed of the (Ox) axis following the axis of motion, i.e. in the same direction as  $\vec{T}$ , and the (Oy) axis following  $\vec{N}$ .

By projecting onto the Oy axis

$$N - p = 0 \Rightarrow N = mg \text{ and } \mu = \operatorname{tg} \varphi = \frac{f}{mg}$$

So  $f = mg \operatorname{tg} \varphi$

$$\text{Then } (*) \Rightarrow \frac{1}{2}mv_B^2 = \frac{1}{2}kx^2 - mgx\mu$$

Hence,

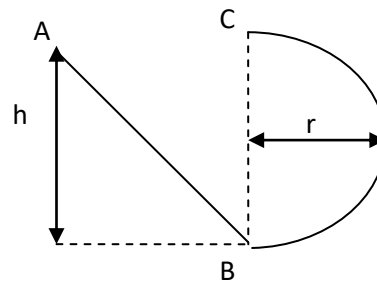
$$v_B^2 = \frac{2}{m} \left( \frac{1}{2}kx^2 - mgx\mu \right)$$

### Exercise 8

A ball slides without friction inside a gutter.

Find the smallest height  $h_{\min}$  from which the ball is

launched to reach point C, without leaving the gutter.



According to the principle of conservation of mechanical energy:

- Between two points A and B:

$$E_{M_A} = E_{M_B} \Rightarrow E_{C_A} + E_{P_A} = E_{C_B} + E_{P_B}$$

Then  $E_{P_A} = E_{C_B} (*)$

Because  $E_{C_A} = 0$

since  $v_A = 0$  because the ball is launched without initial velocity

and  $E_{P_B} = 0$  hence  $E_{P_B} = mgh$  and  $h = 0$

$$\text{So } (*) \Rightarrow mgh = \frac{1}{2}mv_B^2$$

## Chapter VI: Work and Energy

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- Between points B and C:

$$E_{M_B} = E_{M_C} \Rightarrow E_{C_B} + E_{P_B} = E_{C_C} + E_{P_C}$$

Then  $E_{C_B} = E_{C_C} + E_{P_C}$

So  $mgh = \frac{1}{2}mv_C^2 + 2mgr$  (\*)

The ball leaves the gutter at point C when  $N=0$ ,

According to the fundamental principle of dynamics:

$$\Sigma \vec{F} = m\vec{a} \Rightarrow \vec{N} + \vec{p} = m\vec{a}$$

We choose a reference frame consisting of the axis (OT) tangent to the half-sphere and the axis (ON) following the radius and in the direction of  $\vec{N}$  and  $\vec{p}$ .

Projecting onto the ON axis:

$$N+p = m.a_N \Rightarrow N + mg = m \frac{v^2}{r}$$

At point C for  $N=0$  the speed will be:

$$mg = m \frac{v_C^2}{r} \Rightarrow v_C^2 = r g$$

$$(*) \Rightarrow mgh_C = \frac{1}{2}mgr + 2mgr = \frac{5}{2}mgr$$

$$\text{So } h_C = \frac{5}{2}r$$

$h_C$  is the minimum value at which the ball reaches point C without leaving the gutter.

for  $h < h_C$  the ball does not reach point C.

for  $h > h_C$  the ball reaches point C and leaves the trough.

### Exercise 9

1. There are no non-conservative forces. We can therefore use the principle of conservation of mechanical energy between points A and B:

$$\Delta E_m = 0 \quad \text{so} \quad E_{TA} + E_{ppA} = E_{TB} + E_{ppB}$$

Taking the origin of the potential energies in B, and knowing that the velocity in A is equal to 0, we obtain:  $gh = \frac{1}{2}m.v_B^2$

$$\text{So: } v_B = \sqrt{2gh} \quad \text{with } h = l \sin \alpha$$

$$\text{Then } v_B = \sqrt{2g.l \sin \alpha}$$

## Chapter VI: Work and Energy

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2. Study of the system between points  $B$  and  $C$ :

If we choose the origin of abscissas at  $B$  ( $x_B = 0$ ) and therefore  $x_C = d$  (maximum compression distance), we find:

$$v_C^2 - v_B^2 = 2ax \quad (*)$$

To find the expression for acceleration, we use the fundamental principle of dynamics:

$$\Sigma \vec{F}_{ext} = m \vec{a}$$

The forces are weight, reaction  $R$  and spring return force  $F_e$ :

$$\vec{P} + \vec{R} + \vec{F}_e = m \vec{a}$$

Projection on axis ( $Ox$ ):

$$mg \sin \alpha - F_e = mg \sin \alpha - kx = ma$$

$$\text{Then: } a = g \sin \alpha - (k/m) x$$

Replace in (\*), to find:

$$\begin{aligned} -v_B^2 &= 2(g \sin \alpha - (k/m) x)x \Rightarrow (k/m) d^2 - 2g \sin \alpha d - v_B^2 = 0 \\ (k/2) d^2 - (g m \sin \alpha) d - (m g l \sin \alpha) &= 0 \end{aligned}$$

Solving this second-degree equation in  $d$  yields two solutions. One is negative; we retain only the positive solution (the physical solution):

$$d = \frac{mgsin\alpha + \sqrt{m^2g^2sin^2\alpha + 2kmlsin\alpha}}{k}$$

# **EXAMS**

## Continuous Assessment N°01

### Continuous assessment in mechanics

Exam duration: 01 h 30mn

#### Exercise 1: (05 Pts)

A. The momentum  $P$  ( $P=m\vartheta$  where  $m$  is a mass and  $\vartheta$  is a velocity) associated with a photon depends on its frequency  $f$  according to the following expression:

$$P = \sigma^\alpha f^\beta c^\gamma$$

Where  $c$  is the speed of light and  $\sigma$  has the following dimension  $[\sigma] = M \cdot L^2 \cdot T^{-1}$ .

Using dimensional analysis, find the exponents  $\alpha$ ,  $\beta$  and  $\gamma$ .

B. The average velocity of the molecules of a gas is written in the following formula:

$$\vartheta = \sqrt{\frac{PV}{m}}$$

$m$  being the mass of the molecule,  $V$  the volume, and  $p$  the pressure of the gas.

Calculate the relative uncertainty in  $\vartheta$  as a function of  $\Delta p$ ,  $\Delta m$  and  $\Delta V$ .

#### Exercise 2: (05 Pts)

A.  $\vec{i}$ ,  $\vec{j}$  and  $\vec{k}$  being the unit vectors of an orthonormal reference frame (Oxyz), consider the vectors.  $\vec{r}_1 = 2\vec{i} - 2\vec{j} + 3\vec{k}$ ,  $\vec{r}_2 = \vec{i} + \vec{j} + \vec{k}$

1- Calculate the vector product  $\vec{r}_1 \wedge \vec{r}_2$ .

2- Deduce the angle  $\theta$  formed by the two vectors  $\vec{r}_1$  and  $\vec{r}_2$ .

B. Let be a polar coordinate system with origin O and unit vectors  $\vec{u}_\rho$ ,  $\vec{u}_\theta$ .

M is a point with coordinates  $\begin{cases} \rho = 2t^3 + 1 \\ \theta = \omega t \end{cases}$  ( $\omega$  constant).

1- Using a detailed diagram, give the expression of the position vector  $\vec{OM}$  and calculate the velocity vector of point M in polar coordinates.

2- Write this velocity vector  $\vec{v}$  (M) in cartesian coordinates ( $\vec{i}$ ,  $\vec{j}$ ,  $\vec{k}$ ).

#### Exercise 3: (05 Pts)

A particle moves along a trajectory whose equation is  $x^2 + y^2 = 4$  such that  $\mathbf{x}(t) = 2 \sin(\omega t)$ .

Knowing that  $\omega$  is constant and at  $t=0$ , the mobile is at point M (0, R), Determine:

- 1) The component  $y(t)$ .
- 2) Velocity and acceleration vector components and their moduli.
- 3) Tangential and normal accelerations.
- 4) The nature of the motion.

Good luck

## Continuous Assessment N°02

### Continuous assessment in mechanics

Exam duration: 01 h 30mn

#### Exercise 1: (06 pts)

A. Check the homogeneity of this formula:

$$p = \rho g h_1 + h_2 F$$

Such as: P a pressure,  $\rho$  the density, g an acceleration of gravity,  $h_1$  and  $h_2$  are heights and F a force.

B. The period T of oscillation of a pendulum of length l in a gravitational field (gravitational acceleration) g is proposed in the following form:

$$T = k \cdot l^\alpha \cdot g^\beta$$

1. Find  $\alpha$  and  $\beta$  such that k is a dimensionless constant. Write the law of the period T.
2. Calculate the relative uncertainty on T as a function of  $\Delta l$  and  $\Delta g$ .

#### Exercise 2: (07 pts)

A. Let the points  $M_1 (+1,+1,+1)$ ,  $M_2 (+2,+2,+1)$  et  $M_3 (+2,+1,0)$  ;

1. Find the angle  $\widehat{M_1 M_2 M_3}$ .
2. Calculate the area of the triangle  $M_1 M_2 M_3$

B. A material point M is identified by its Cartesian coordinates (x, y, z).

1. Write down the relationship between Cartesian coordinates and cylindrical coordinates (using a diagram).
2. Write the position vector in cylindrical coordinates and deduce the velocity vector in the same coordinate system.

3. If the position of the point is marked in cylindrical coordinates by: 
$$\begin{cases} \rho = 4t^2 \\ \theta = \omega t \\ z = \sqrt{t} \end{cases}$$

Find the expression of the velocity vector  $\vec{v}$  in cylindrical coordinates.

#### Exercise 3: (06 pts)

The coordinates x and y of a moving point M in the plane (oxy) vary with time t according to the following relationships:  $x = 2t$  ,  $y = 2t^2$

Find:

1. The equation of the trajectory.
2. Velocity components and modulus v.
3. Components of acceleration and its modulus a.
4. Nature of motion.
5. Tangential acceleration  $a_T$  and normal acceleration  $a_N$ ; deduce the radius of curvature.

Good luck

## Continuous Assessment N°03

### Exam to replace the Continuous assessment in mechanics

Exam duration: 01 h 30mn

#### Exercise 01: (05 Pts)

The limiting **velocity**  $\vartheta$  of a ball of radius  $r$  of **mass**  $m$  and **density**  $\rho$ , falling into a viscous medium (a fluid) of **density**  $\rho_f$  and viscosity coefficient  $\eta$  معامل الميوعة , is given by :

$$\vartheta = \frac{2r^2(\rho - \rho_f)}{9\eta} g$$

$g$  is the acceleration of gravity.

1. Using the dimensional equations, determine the dimension of  $\eta$  and derive its unit in the MKSA system.
2. Determine the relative uncertainty on  $\eta$ , as a function of  $\Delta r$ ,  $\Delta \rho$ ,  $\Delta \rho_f$  and  $\Delta v$ .

#### Exercise 0 2 : (05 Pts)

A. Let the points A (+1,+1,+1), B (+2,+2,+1) et C (+2,+1,0)

Calculate the scalar product  $\overrightarrow{AB} \cdot \overrightarrow{AC}$  and the vector product  $\overrightarrow{AB} \wedge \overrightarrow{AC}$ . What do these two products represent? Deduce the angle between the vectors  $\overrightarrow{AB}$  and  $\overrightarrow{AC}$ .

B. Define cylindrical coordinates and give the transition relationships between Cartesian coordinates and cylindrical coordinates.

Write the position vector in cylindrical coordinates and calculate the velocity vector in this coordinate system.

#### Exercise 03: (05 Pts)

From the ground, a balloon rises with a constant initial velocity  $v_0$  (following  $y$ ). The wind gives the balloon a horizontal velocity  $\mathbf{V}_x = \gamma \cdot \mathbf{y}$  ( $\gamma$  constant).

- a- Determine the equations of motion  $x(t)$  and  $y(t)$ . Deduce the equation of the trajectory  $y=f(x)$ .
- b- Calculate the accelerations  $a$ ,  $a_N$  and  $a_T$ . Deduce the radius of curvature.

Good luck



**Continuous Assessment N°04**

**Continuous assessment in mechanics**

Exam duration: 01 h 30mn

**Exercise 1: (6 pts)**

The velocity limit reached by a weighted parachute is a function of its weight  $P$  and its surface area  $S$ , and is given by:

$$v = \sqrt{\frac{P}{k.S}}$$

- 1) Give the dimension of the constant  $k$ .
- 2) Calculate the limiting speed of a parachute with the following characteristics :  
 $M=90 \text{ kg}$ ,  $S=80 \text{ m}^2$ ,  $g=9,81 \text{ m/s}^2$ , and  $k=1,15 \text{ MKS}$ .
- 3) With the weight known to within 2% and the surface area to within 3%, calculate the relative uncertainty  $\Delta v/v$  on the velocity  $v$ , and the absolute uncertainty  $\Delta v$  and deduce the condensed form of this velocity.

**Exercise 2: (5 pts)**

**A.** In the vector space related to the orthonormal basis  $(\vec{i}, \vec{j}, \vec{k})$ , consider the vectors  $\vec{U}(0, 3, 1)$ ,  $\vec{V}(0, 1, 2)$ .

- 1) Calculate the scalar product  $\vec{U} \cdot \vec{V}$  and the angle  $\varphi$  acute between  $\vec{U}$  and  $\vec{V}$ .
- 2) Determine the components of the vector  $\vec{W} = \vec{U} \wedge \vec{V}$  then calculate  $\|\vec{W}\|$  by two methods. What does the latter represent.
- 3) Calculate the mixed product  $(\vec{U}, \vec{V}, \vec{W})$ , what does this product represent.

**B.** Does each of the following expressions have a meaning? If yes, specify whether it is a vector or a real. If not, say why (without calculation).

- |   |   |  |
|---|---|--|
| 1) $\vec{A} \cdot (\vec{B} \wedge \vec{C})$ | 2) $\vec{A} \wedge (\vec{B} \cdot \vec{C})$                   | 3) $\vec{A} \wedge (\vec{B} \wedge \vec{C})$                 |
| 4) $\vec{A} \cdot (\vec{B} \cdot \vec{C})$  | 4) $(\vec{A} \wedge \vec{B}) \wedge (\vec{C} \wedge \vec{B})$ | 5) $(\vec{A} \wedge \vec{B}) \cdot (\vec{C} \wedge \vec{C})$ |

**Exercise 3: (8 pts)**

Let be a cylindrical reference frame with origin  $O$ , unit vectors  $\vec{u}_\rho, \vec{u}_\theta, \vec{u}_z$ .  $M$  is any point with coordinates  $(\rho, \theta, z)$ .

Using a detailed diagram, give the expression for the position vector  $\vec{OM}$  as a function of the unit vectors  $\vec{u}_\rho, \vec{u}_\theta, \vec{u}_z$ .

1. Find the velocity vector in cylindrical coordinates.
2. Express the elementary displacement vector in cylindrical coordinates.
3. Write the expression for the elementary volume in this frame of reference and deduce the volume of a cylinder.

Good luck

## Continuous Assessment N°05

### Continuous assessment in mechanics

Exam duration: 01 h 30mn

#### Exercise 1: (6pts)

A) A particle of mass  $m$  enclosed in a cubic box of side  $L$ , has a kinetic energy  $E$  such that:

$$E = \frac{\pi^2 \sigma^2}{2mV^{\frac{2}{3}}} n^2$$

Where  $V$  the volume of the box and  $n$  a dimensionless number.

Using the dimensional equations, find the dimension of  $\sigma$ .

B) The average velocity of the molecules in a gas can be written as:

$$\vartheta = \sqrt{\frac{pV}{m}}$$

$m$  being the mass of the molecule,  $V$  the volume, and  $p$  the pressure of the gas.

Calculate the relative uncertainty in  $\vartheta$  as a function of  $\Delta p$ ,  $\Delta m$  and  $\Delta V$ .

#### Exercise 2 : (8pts)

A material point  $M$  is identified by its Cartesian coordinates  $(x, y)$ .

1. Write the relationship between Cartesian coordinates and polar coordinates.
2. Give the expression of the unit vectors  $\vec{U}_r$  and  $\vec{U}_\theta$  as a function of the unit vectors  $\vec{i}$  and  $\vec{j}$ .
3. Find the expression of the velocity vector  $\vec{v}$  of point  $M$  in polar coordinates.
4. Give the expression of the vector  $\vec{A} = 2x\vec{i} - y\vec{j}$  in polar coordinates.

#### Exercise 3: (6pts)

The coordinates  $x$  and  $y$  of a moving point  $M$  in the plane  $(oxy)$  vary with time  $t$  according to

the following relationships:  $x = t + 1$ ,  $y = \frac{t^2}{2} + 2t$

Find :

1. The equation of the trajectory.
2. Components of velocity and its modulus  $v$ .
3. Components of acceleration and its modulus.
4. Nature of motion.
5. Tangential and normal accelerations.
6. Radius of curvature  $R$ .

Good luck

## Continuous Assessment N°06

### Exam to replace the Continuous assessment in mechanics

Exam duration: 01 h 30mn

#### Exercise 1: (6pts)

The energy of a photon is given by the expression  $E=h\nu$ , where  $h$  is the Planck constant and  $\nu$  the photon frequency.

- 1- Give the dimension of  $h$ .
- 2- Find the expression for the wavelength  $\lambda$ , assuming it to be of the form  $\lambda=k.h^x m^y V^z$ . Where  $k$  is a dimensionless constant,  $m$  and  $V$  represent the photon's mass and velocity respectively.
- 3- Determine the relative uncertainty on  $\lambda$  as a function of  $\Delta m$ ,  $\Delta h$  and  $\Delta V$ .

#### Exercise 2: (8pts)

1. Define cylindrical coordinates.
2. Write the unit vectors  $\vec{u}_\rho$ ,  $\vec{u}_\theta$  and  $\vec{u}_z$  in the cylindrical coordinate system in terms of the unit vectors  $(\vec{i}, \vec{j}$  and  $\vec{k})$ .
3. Write the elementary displacement vector in cylindrical coordinates.
4. Deduce the volume of a cylinder.
5. Write the vector  $\vec{A} = x\vec{i} + 2y\vec{j} - z\vec{k}$  in cylindrical coordinates.

#### Exercise 3: (6pts)

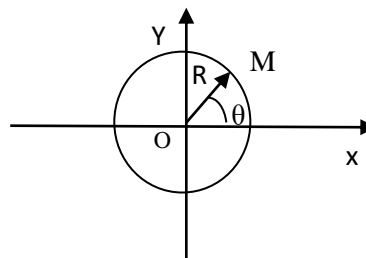
A. Let a moving point  $M$  describe a circle of radius  $R$  and center  $O$  with angular velocity  $\omega=d\theta/dt$ . At time  $t=0$  the point  $M$  is at  $A$ .

1. Write the coordinates of  $M$  as a function of  $R$  and  $\theta$ .
2. Calculate the modulus of the velocity of point  $M$ .
3. Determine the components of the acceleration on the axes  $Ox$  and  $Oy$  (Cartesian coordinates) on the one hand, and on the axes parallel and perpendicular to  $OM$  on the other (polar coordinates).

B. Assume  $\alpha=d\omega/dt$  ( $\alpha$  is a non-zero constant). Give the expressions for  $\omega$  and  $\theta$  as a function of time.

Recall that at  $t=0$ ,  $\theta_0=0$  and  $\omega=\omega_0$ .

What relationship exists between  $\omega$  and  $\theta$ .



Good luck

## Continuous Assessment N°07

### Continuous assessment in mechanics

Exam duration: 01 h 30mn

#### Exercise 1: (06pts)

A) Experience has shown that the velocity  $v$  of sound in a gas is a function only of the gas density  $\rho$  and its coefficient of compressibility  $\chi$ . It is given by :

$$v = k \rho^x \chi^y$$

Recall that  $\chi$  is homogeneous to the inverse of a pressure, with  $k$  a dimensionless constant. Determine the velocity of sound relationship.

B) The focal length  $f$  of a lens is determined from the formula:

$$f = \frac{D^2 - a^2}{4D}$$

Calculate the absolute uncertainty  $\Delta f$  as a function of  $\Delta D$  and  $\Delta a$ .

#### Exercise 2: (08pts)

Let two vectors  $\vec{A}$  and  $\vec{B}$  in the orthonormal reference frame (Oxyz), be defined by :

$$\vec{A} = 2 \cdot \vec{i} + 4 \cdot \vec{j} - 5 \cdot \vec{k} , \vec{B} = -\vec{i} + \vec{j} - 2 \cdot \vec{k}$$

1. Calculate and represent the two vectors ( $\vec{A} + \vec{B}$ ) and ( $\vec{A} - \vec{B}$ ).
2. Calculate their moduli.
3. Calculate the scalar product ( $\vec{A} \cdot \vec{B}$ ). Determine the angle  $\theta = (\vec{A}, \vec{B})$ .
4. Calculate the vector product ( $\vec{A} \wedge \vec{B}$ ), what present  $|\vec{A} \wedge \vec{B}|$  ?
5. Write the vector  $\vec{C} = y \cdot \vec{i} - 2 \cdot x \cdot \vec{j} + z \cdot \vec{k}$  in cylindrical coordinates i.e. as a function of  $\rho, \theta, z$  and  $\vec{u}_\rho, \vec{u}_\theta, \vec{u}_z$ . (**indication:** use the relations for passing coordinates (x,y,z) and unit vectors  $\vec{i}, \vec{j}, \vec{k}$  as a function of unit vectors  $\vec{u}_\rho, \vec{u}_\theta, \vec{u}_z$ ).

#### Exercise 2: (06pts)

The coordinates  $x$  and  $y$  of a moving point  $M$  in the plane (oxy) vary with time  $t$  according to the following relationships:  $x = t, y = t^2 - t$

Find:

1. The equation of the trajectory.
2. Velocity components and modulus  $v$ .
3. Components of acceleration and its modulus.
4. Tangential and normal accelerations as a function of the modulus of  $v$ .
5. Radius of curvature  $R$ .

Good luck

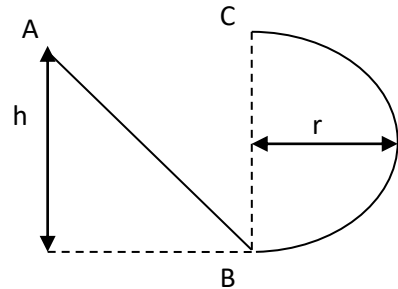
**Final Exam N°01**

**Final Exam of Mechanics**

Exam duration: 01 h 30mn

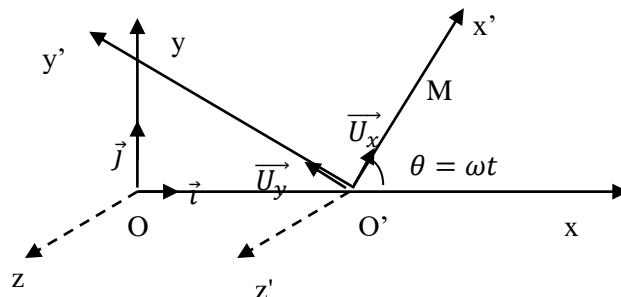
**Course questions: (6pts)**

- 1- State and demonstrate the kinetic energy theorem.
- 2- A ball slides without friction inside a gutter. Find the smallest height  $h_{\min}$  from which the ball is launched to reach point C, without leaving the gutter.



**Exercise 1 : (7pts)**

In the (Oxy) plane, a point O' (the origin of the moving reference frame) moves along the (Ox) axis such that  $|OO'|=t$ . The reference frame (O'X'Y') rotates around Oz with a constant angular velocity  $\omega$ . A moving point M ( $O'M=r$ ) moves along the axis (O'X') according to the law  $r = r_0 (\cos \omega t + \sin \omega t)$  with  $r_0 = \text{constant}$ . Determine at time t as a function of  $r_0$  and  $\omega$ .

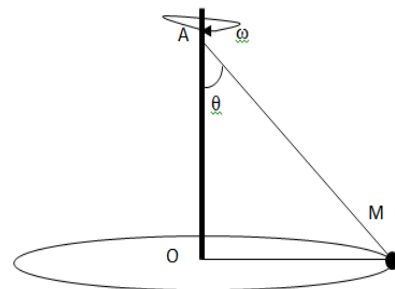


- 1- The velocity  $\vec{v}_r$  and the entrainment velocity  $\vec{v}_e$ , in the moving reference frame (OX'Y'), deduce the absolute velocity  $\vec{v}_a$  in the same reference frame.
- 2- Relative acceleration  $\vec{a}_r$ , entrainment acceleration  $\vec{a}_e$  and Coriolis acceleration  $\vec{a}_c$  in the moving frame of reference, deduce the absolute acceleration  $\vec{a}_a$  in this frame of reference.

**Exercise 2 : (7pts)**

A ball of mass m is attached by two wires (Am and Om) to a vertical pole. The whole system rotates with a constant angular velocity  $\omega$  around the axis of the pole (we know g the acceleration of gravity,  $\theta$  and  $L = |\overline{OM}|$ )

1. Assuming  $\omega$  is large enough to keep both wires taut, find the force (wire tension) each wire exerts on the ball.
2. What is the minimum angular velocity  $\omega_{\min}$  for which the bottom wire remains taut.



Good luck

**Final Exam N°02**

**Final Exam of Mechanics**

Exam duration: 01 h 30mn

**Course questions: (5pts)**

- 1- Why use dimensional analysis?
- 2- What can it say about the total mechanical energy of a system in the presence of frictional forces?
- 3- What is the difference between a conservative force (قوة منحفضة) and a non-conservative force? Give an example for each one.
- 4- Calculate the work of a force  $F=1.5 \cdot 10^4 \text{N}$  supplied to move a body a height (AB) of 3 meters (vertically).
- 5- Calculate the work of the spring return force with stiffness constant  $k$  ( $d\vec{l} = dx \cdot \vec{i}$ ).

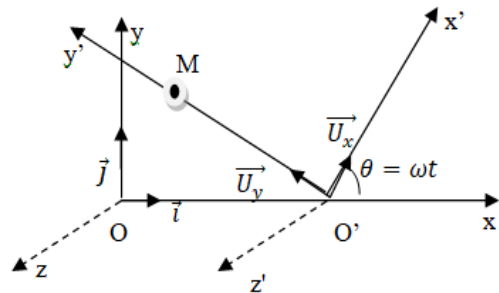
**Exercise 1: (7pts)**

Consider the fixed reference frame  $R(Oxyz)$  where point  $O'$  moves **along axis (Ox)** with **constant velocity**  $v_0$ . Linked to  $O'$  is the moving reference frame  $(O'x'y'z')$  which rotates **around (Oz)** with **constant** angular velocity  $\omega$ . A moving point  $M$  moves **along the (O'y')** axis with constant **acceleration**  $\gamma$ .

At time  $t=0$ , the axes  $(Ox)$  and  $(O'x')$  are coincident and  $M$  is at  $O$ .

Calculate in the moving frame:

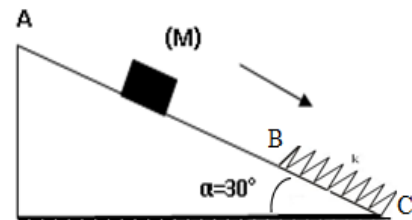
- 1- The relative velocity  $\vec{v}_r$  and the entrainment velocity  $\vec{v}_e$ , deduce the absolute velocity  $\vec{v}_a$ .
- 2- The relative acceleration  $\vec{a}_r$ , the entrainment acceleration  $\vec{a}_e$  and Coriolis acceleration  $\vec{a}_c$ , deduce the absolute acceleration  $\vec{a}_a$ .



**Exercise 2: (8pts)**

Consider a small block of mass  $m = 2\text{kg}$  dropped without initial velocity at point A of an inclined plane at an angle  $\alpha=30^\circ$  to the horizontal. Point A is at a height  $h_A=5\text{m}$  from the horizontal.

- 1- Knowing that the coefficient of dynamic friction on **plane AB is  $\mu_d=0.2$** , applying the fundamental principle of dynamics, what is the acceleration of the block on plane  $AB=8\text{m}$ ?
- 2- Calculate the speed of the block when it reaches point B.
- 3- Using the kinetic energy theorem, find the speed of the block at point B.
- 4- At point B, the block hits a spring with stiffness constant  $k=100\text{N/m}$  at speed  $V_B$ . Calculate the maximum compression ( $x$ ) of the spring (given  $g = 10 \text{ m/s}^2$ ).



Good luck

**Final Exam N°03**

**Final Exam of Mechanics**

Exam duration: 01 h 30mn

**Exercise 1:** (6 pts)

A) The position of a moving body, subjected to a force  $F$ , is marked by its abscissa  $x$ , at any instant  $t$  according to the relation:  $x = a.(bt) + F.c$

Give the dimensions of the different quantities  $a$ ,  $b$  and  $c$ . What might the quantity  $c$  represent?

B) The average velocity of the molecules of a gas is written in the following formula:

$$\vartheta = \sqrt{\frac{pV}{m}}$$

$m$  being the mass of the molecule,  $V$  the volume, and  $p$  the pressure of the gas.

Calculate the relative uncertainty on  $\vartheta$  as a function of  $\Delta p$ ,  $\Delta m$  and  $\Delta V$ .

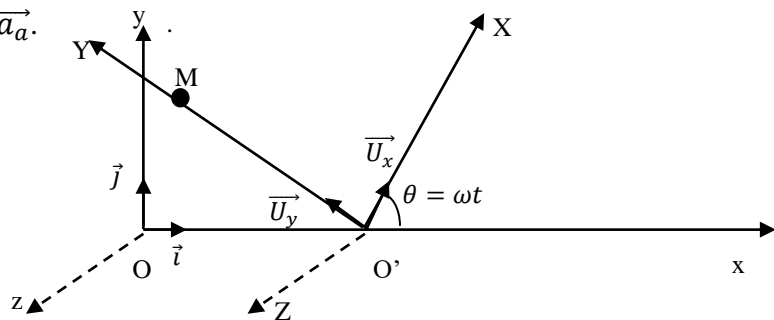
**Exercise 2:** (8 pts)

Let the reference frame be  $R(Oxyz)$  and the point  $O'$  moves on the axis  $(Ox)$  with a constant velocity  $v_0$ . We link to  $O'$  the reference frame  $(O'XYZ)$  which rotates around  $(Oz)$  with a constant angular velocity  $\omega$ . Mobile  $M$  moves along axis  $(O'Y)$  such that  $|O'M| = (t^2 + 2)$  (without initial velocity).

At time  $t=0$ , axis  $(O'X)$  is coincident with  $(Ox)$  and point  $M$  is at  $O'$ .

Calculate in the **moving reference frame** :

- 1- The relative velocity  $\vec{v}_r$  and the entrainment velocity  $\vec{v}_e$ , deduce the absolute velocity  $\vec{v}_a$ .
- 2- The relative acceleration  $\vec{a}_r$ , the entrainment acceleration  $\vec{a}_e$  and Coriolis acceleration  $\vec{a}_c$ , deduce the absolute acceleration  $\vec{a}_a$ .



**Exercise 3:** (6 pts)

A particle moves along a trajectory whose y-coordinate equation is given by:  $y(t) = t^2 + 1$  such that at each instant  $v_x = v_0 = \text{cste}$ .

If at  $t=0$ ,  $x_0=0$ ; determine :

- 1- The equation of the particle's trajectory.
- 2- The particle's velocity and acceleration.
- 3- Normal and tangential accelerations and radius of curvature.

Good luck

**Final Exam N°04**

**Final Exam of Mechanics**

Exam duration: 01 h 30mn

**Cours Questions: (5pts)**

**I.** What is the work done by a force  $\vec{F} = 2t \vec{i}$  acting on a particle of mass  $m=2\text{kg}$  for a displacement along the horizontal ( $dl = dx$ ).

1. Is this force conservative?
2. Name three conservative forces.
3. Deduce the power of this force.

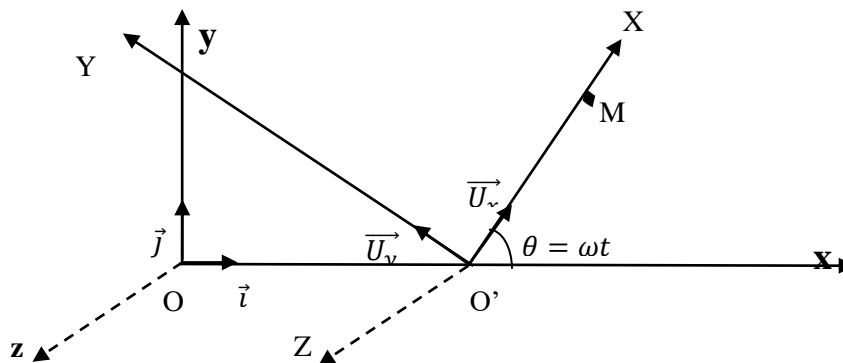
**II.** Define the dynamic and static coefficients of friction. What is the relationship between these two coefficients? what is the dimension of the friction coefficient?

**Exercise 1 : (8 pts)**

In the Oxy plane, consider a system of moving axes (O'XY), such that (Ox) makes a variable angle  $\theta$  with (O'X). Point O' moves along axis (Ox) with constant acceleration  $a$ . A point M moving along the O'X axis is marked by  $\vec{O'M} = \vec{r}$ . We call relative motion of M its motion with respect to (O'XYZ) and absolute motion with respect to (Oxyz). (see figure 1)

Calculate in the moving reference frame :

- 1- The relative velocity  $\vec{v}_r$  and the entrainment velocity  $\vec{v}_e$ , deduce the absolute velocity  $\vec{v}_a$ .
- 2- The relative acceleration  $\vec{a}_r$ , the entrainment acceleration  $\vec{a}_e$  and Coriolis acceleration  $\vec{a}_c$ , deduce the absolute acceleration  $\vec{a}_a$ .



**Figure 1**

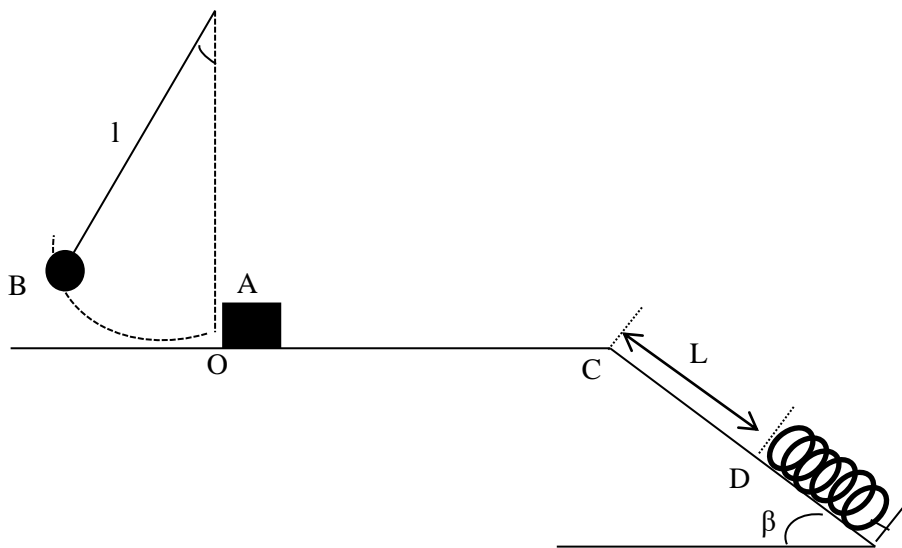
**Exercise 2 : (7 pts)**

- A ball B of mass  $m$ , attached to an inextensible wire of length  $l$ , is moved away from its equilibrium position by an angle  $\alpha$ . It is dropped without initial velocity. Passing through the vertical position, the ball strikes (touches) a body A of the same mass and stops, body A passes from point O to point C ( $\vec{OC}=\vec{d}$ ) on a rough horizontal plane of friction coefficient  $\mu$ .

- 1- Show the forces exerted on body A.
- 2- What is the nature of the motion on the horizontal plane?
3. Express the velocity of ball B just before it touches body A.



4. Using the principle of conservation of momentum of the system, determine the velocity of body A after the interaction.
  5. If  $v_A = v_B$  at point O, give the velocity of body A at point C as a function of  $g$ ,  $l$ ,  $d$ ,  $\alpha$  and  $\mu$ .
  6. By what angle must ball B be moved away for body A to arrive at point C with zero velocity.
- From point C, body A approaches the perfectly smooth (no friction) path  $CD=L$ , inclined at an angle  $\beta$  to the horizontal (Fig. 2). It arrives, without initial velocity, on a perfect spring of length  $l_0$  and stiffness constant  $k$ .
    1. Show the forces exerted on A as the spring compresses.
    2. What is the value of the spring's maximum compression?
- We give  $m=200g$ ,  $d=OC=1m$ ,  $l=10\text{ cm}$ ,  $L=1m$ ,  $\mu=0.1$ ,  $g=10m/s^2$ ,  $k=140N/m$ ,  $\beta=30^\circ$ .



**Figure 2**

**Final Exam N°05**

**Final Exam of Mechanics**

Exam duration: 01 h 30mn

**Exercise 1 : (6pts)**

KEPLER's third law relates the period  $T$  to the semi-major axis  $a$  of a planet's orbit around the

sun as follows: 
$$\frac{T^2}{a^3} = \frac{4\pi^2}{GM_S}$$

With  $G$  the universal gravitational constant and  $M_S$  the mass of the sun.

We give:  $G = (6.668 \pm 0.005) \cdot 10^{-11} \text{SI}$

For the earth:  $T = (365.25636567 \pm 0.00000001)$  days and  $a = (1.4960 \pm 0.0003) \cdot 10^{11} \text{m}$

- 1- Determine the dimension and unit of  $G$ .
- 2- Determine the mass of the sun  $M_S$  and the absolute uncertainty  $\Delta M_S$  on this mass.

**Exercise 2: (8pts)**

Let the reference frame be  $R(Oxyz)$  and the point  $O'$  moves along the **axis (Ox)** with a constant **acceleration**  $\gamma$ ; and a **positive initial velocity**  $V_0$ . We link to  $O'$  the reference frame  $(O'XYZ)$  which rotates around  $(Oz)$  with a constant angular velocity  $\omega$ . The coordinates of a moving body  $M$  in the moving frame of reference are:  $X' = t^2 + 2$  and  $Y' = 2t$ .

At time  $t=0$ , the axis  $(O'X)$  coincides with  $(Ox)$ . (Fig 1)

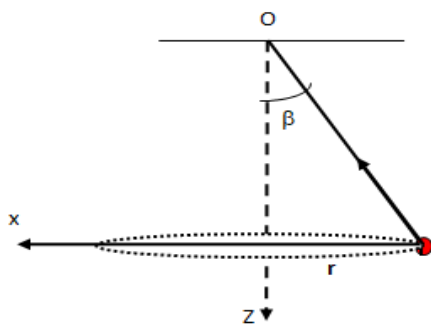
Calculate in the **moving frame** of reference  $(O'XYZ)$  :

- 1- The relative velocity  $\vec{v}_r$  and the entrainment velocity  $\vec{v}_e$ , deduce the absolute velocity  $\vec{v}_a$ .
- 2- The relative acceleration  $\vec{a}_r$ , the entrainment acceleration  $\vec{a}_e$  and Coriolis acceleration  $\vec{a}_c$ , deduce the absolute acceleration  $\vec{a}_a$ .

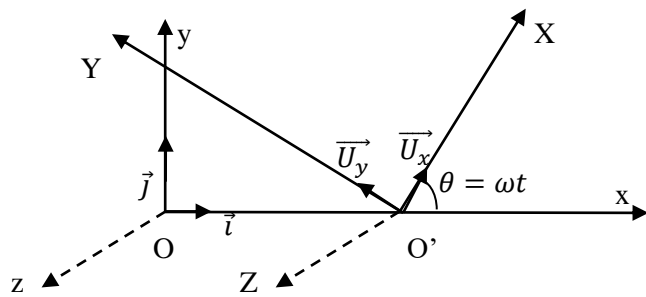
**Exercise 3: (6pts)**

A material point, of mass  $m$ , is suspended at a fixed point  $O$  by a wire of length  $l$  inextensible and negligible mass. By rotating it around axis  $Oz$ , it acquires a constant angular velocity  $\omega$ . It describes a horizontal circle of radius  $r$ . (Fig 2)

1. Find the expression for the wire tension.
2. Find the expression for the inclination  $\beta$  of the wire with respect to the vertical.



(Fig 2)



(Fig 1)

Good luck

**Final Exam N°06**

**Final Exam of Mechanics**

Exam duration: 01 h 30mn

**Exercise 1 (6pts)**

1. A moving body whose motion is uniformly circular is subject to an acceleration:
- |             |  |
|-------------|--|
| a- constant | c- whose modulus is constant                     |
| b- zero     | d- directed towards the center of the trajectory |

Give the correct answers.

2. A particle M moves along a parabolic trajectory with equation :

$$y = x^2 \text{ with } x(t) = 2t$$

- a- Determine the components of velocity and acceleration, and calculate their moduli.  
 b- Determine the tangential and normal accelerations, and deduce the radius of curvature R.

**Exercise 2 (8pts)**

A point M moves with constant velocity  $V_0$  on the axis (OX) of a reference frame (OXYZ) which rotates with constant angular velocity  $\omega$  around (Oz) in the plane (Oxy) ( $\overrightarrow{OO'} = \vec{0}$ ).

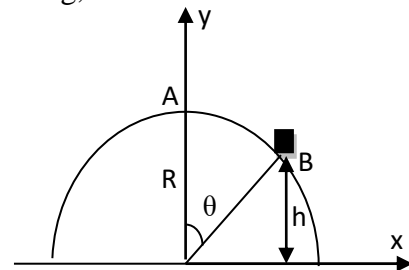
- 1- What is the expression of  $\overrightarrow{OM}$  in the fixed reference frame. Calculate absolute velocity and absolute acceleration.
- 2- Calculate relative velocity and entrainment velocity, check that  $\vec{v}_a = \vec{v}_r + \vec{v}_e$ .
- 3- Calculate the relative acceleration  $\vec{a}_r$ , entrainment acceleration  $\vec{a}_e$  and Coriolis acceleration  $\vec{a}_c$ , check that  $\vec{a}_a = \vec{a}_r + \vec{a}_e + \vec{a}_c$ .

**Exercise 3 (6pts)**

A piece of ice M of mass m slides frictionlessly over the outer surface of an igloo, which is a half-sphere of radius R with a horizontal base.

At  $t=0$ , it is released from point A without any initial velocity.

1. Find the expression for the velocity at point B, as a function of g, R and  $\theta$ .
2. Using the fundamental relation of dynamics, determine the expression of  $|\vec{N}|$  the reaction of the igloo on M at point B as a function of velocity  $v_B$ .
3. At what height does M leave the sphere?
4. At what speed does M arrive at the axis (Ox)?



Good luck

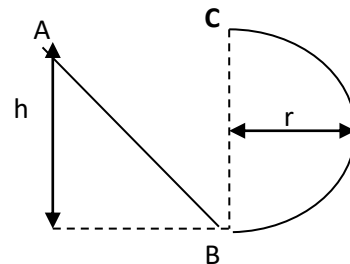
**Final Exam N°07**

**Final Exam of Mechanics**

Exam duration: 01 h 30mn

**Cours Questions: (6pts)**

- 1- Give Newton's three laws.
  - 2- What is the difference between a conservative force and a non-conservative force, with examples for each case?
  - 3- In which case do we have conservation of mechanical energy, and what do we have in the opposite case?
  - 4- State the kinetic energy theorem and demonstrate it.
  - 5- A ball is thrown without initial velocity and without friction inside a gutter.
- Find the height at which the ball reaches point C and changes direction.



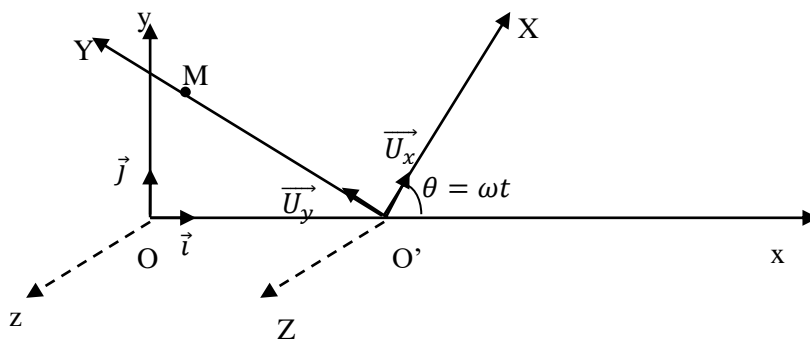
**Exercise 1 (7 pts) :**

Let the reference frame be  $R(Oxyz)$  and the point  $O'$  moves on **the axis (Ox)** with a constant **velocity  $v_0$** . We link to  $O'$  the reference frame  $(O'XYZ)$  which rotates around  $(Oz)$  with a constant angular velocity  $\omega$ . Mobile  $M$  moves along **axis (O'Y)** with constant **acceleration  $\gamma$**  (no initial velocity).

At time  $t=0$ , axis  $(O'X)$  is coincident with  $(Ox)$  and point  $M$  is at  $O'$ .

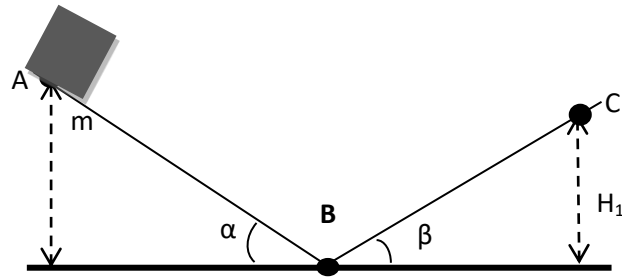
Calculate in the **moving reference** frame:

- 1- The relative velocity  $\vec{v}_r$  and the entrainment velocity  $\vec{v}_e$ , deduce the absolute velocity  $\vec{v}_a$ .
- 2- The relative acceleration  $\vec{a}_r$ , the entrainment acceleration  $\vec{a}_e$  and Coriolis acceleration  $\vec{a}_c$ , deduce the absolute acceleration  $\vec{a}_a$ .



**Exercise 2 (7 pts) :**

Consider a small block of mass  $m = 5\text{kg}$  dropped without initial velocity at point A of an inclined plane at an angle  $\alpha = 30^\circ$  to the horizontal (see figure 2). Point A is at a height  $h_0 = 5\text{m}$ .



1- What is the value of the coefficient of static friction  $\mu_s$  that keeps the mass in equilibrium at point A?

2- Knowing that the coefficient of dynamic friction on plane  $AB$  is  $\mu_d=0.2$ , apply the fundamental principle of dynamics:

- What is the nature of the motion on  $AB$ ?
- Calculate the speed of the block when it reaches point  $B$ .
- What can be said about the total mechanical energy of the mass  $m$ ?

3- After passing through point  $B$  at speed  $V_B$ , the mass moves up the inclined plane  $BC$  (angle= $20^\circ$ ), and stops at point  $C$ . Knowing that the coefficient of friction remains the same, determine the height  $h_1$  of point  $C$ ?

Good luck

**Final Exam N°08**

**Final Exam of Mechanics**

Exam duration: 01 h 30mn

**Course question (5pts)**

- 1- Which physical quantity has a unit and no dimension?
- 2- What are the different coordinate systems? In which system are the components dependent?
- 3- In the case where we have only conservative forces, give the theorems we can use?
- 4- Define the dynamic and static coefficients of friction. Which is the most important?

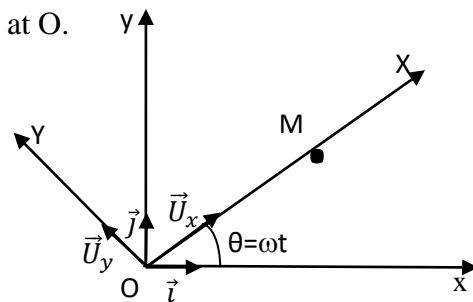
**Exercise 1 (8pts)**

In the (Oxy) plane, consider a system of moving axes (OXY) with the same origin O, rotating with a constant angular velocity  $\omega$  around (OZ). A moving point M moves along axis (OX) with constant **acceleration**  $\gamma$  and no initial velocity. We call relative motion of M its motion with respect to (OXY), and absolute motion with respect to (Oxy).

At time  $t=0$ , axes (Ox) and (OX) are coincident and M is at O.

Calculate in the moving reference frame :

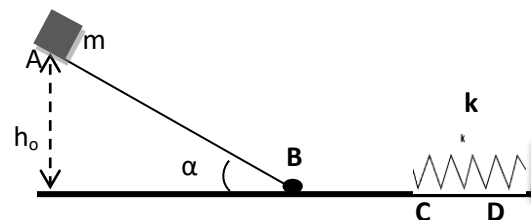
- 1- The velocity and relative acceleration of M.
- 2- Entrainment velocity of M and acceleration.
- 3- Coriolis acceleration.
- 4- Deduce its absolute velocity and acceleration.



**Exercise 2 (7 pts)**

Consider a small block of mass  $m = 5\text{kg}$  dropped without initial velocity at point A of an inclined plane at an angle  $\alpha = 30^\circ$  to the horizontal. Point A is at a height  $h_0 = 5\text{m}$  from the horizontal.

. What is the value of the coefficient static friction coefficient  $\mu_s$  that keeps the mass the mass in equilibrium at point A.



2. Knowing that the coefficient of dynamic friction on plane AB is  $\mu_d = 0.2$ , apply the fundamental principle of dynamics:

- What is the nature of the motion on plane AB?
- Calculate the speed of the block when it reaches point B.

3. After passing through point B at speed  $V_B$ , the mass arrives at point C. Knowing that the coefficient of friction is negligible on the plane BC :

- Deduce the velocity at point C?
- Calculate the maximum compression of the spring, given a stiffness constant equal to  $k = 100\text{N/m}$ ? (we given  $g = 10\text{ m/s}^2$ )

Good luck

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